

The Solomon equations

$$\frac{d\langle I_z \rangle}{dt} = -\rho_I (\langle I_z \rangle - \langle I_0 \rangle) - \sigma_{IS} (\langle S_z \rangle - \langle S_0 \rangle) \quad \dots [1]$$

$$\frac{d\langle S_z \rangle}{dt} = -\rho_S (\langle S_z \rangle - \langle S_0 \rangle) - \sigma_{SI} (\langle I_z \rangle - \langle I_0 \rangle) \quad \dots [2]$$

Steady state NOE

If spin S were saturated ($\langle S_z \rangle = 0$), in the steady state we have:

$$\begin{aligned} \frac{d\langle I_z \rangle}{dt} &= 0 = -\rho_I (\langle I_z \rangle - \langle I_0 \rangle) - \sigma_{IS} (\langle S_z \rangle - \langle S_0 \rangle) \\ &= -\rho_I (\langle I_z \rangle - \langle I_0 \rangle) + \sigma_{IS} \langle S_0 \rangle \\ \Rightarrow \frac{(\langle I_z \rangle - \langle I_0 \rangle)}{\langle S_0 \rangle} &= \frac{\sigma_{IS}}{\rho_I} \\ \therefore \frac{(\langle I_z \rangle - \langle I_0 \rangle)}{\langle I_0 \rangle} \frac{\langle I_0 \rangle}{\langle S_0 \rangle} &= \frac{\sigma_{IS}}{\rho_I} \Rightarrow \frac{(\langle I_z \rangle - \langle I_0 \rangle)}{\langle I_0 \rangle} = \frac{\langle S_0 \rangle}{\langle I_0 \rangle} \frac{\sigma_{IS}}{\rho_I} = \frac{\gamma_S}{\gamma_I} \frac{\sigma_{IS}}{\rho_I} \quad \dots [3] \end{aligned}$$

Note that under these conditions, Equation [2] above does not apply, since relaxation processes of the S spin are rendered ineffective ('short circuited') by saturation.

Equation [3] describes the steady state Overhauser effect, which leads to an enhancement of the signal of spin I when spin S is saturated, provided these two spins relax each other by mutual dipolar interactions as given by the Solomon equations.

Under motional narrowing ('extreme narrowing') conditions, the ratio of cross-relaxation rate σ and auto-relaxation rate ρ is $\frac{1}{2}$. For ^{13}C NMR therefore the nuclear Overhauser effect could result in a relative enhancement by a factor of 2 when protons are saturated. This implies that the observed ^{13}C signal under proton saturation conditions could be upto 3 times the normal ^{13}C signal.

Transient NOE

If we assume $\rho_I = \rho_S$, represent the cross-relaxation rate by σ ($= \sigma_{IS} = \sigma_{SI}$) and invert the S spins at time 0, we have from Eqs. [1] and [2]:

$$\frac{d(\langle I_z \rangle + \langle S_z \rangle - \langle I_0 \rangle - \langle S_0 \rangle)}{dt} = -(\rho + \sigma)(\langle I_z \rangle + \langle S_z \rangle - \langle I_0 \rangle - \langle S_0 \rangle)$$

$$\frac{d(\langle I_z \rangle - \langle S_z \rangle - \langle I_0 \rangle + \langle S_0 \rangle)}{dt} = -(\rho - \sigma)(\langle I_z \rangle - \langle S_z \rangle - \langle I_0 \rangle + \langle S_0 \rangle)$$

Solving these equations with the initial condition as stated gives rise to the following:

$$(\langle I_z \rangle + \langle S_z \rangle - \langle I_0 \rangle - \langle S_0 \rangle) = -2\langle S_0 \rangle e^{-(\rho+\sigma)t} \quad \dots[4]$$

$$(\langle I_z \rangle - \langle S_z \rangle - \langle I_0 \rangle + \langle S_0 \rangle) = 2\langle S_0 \rangle e^{-(\rho-\sigma)t} \quad \dots[5]$$

Adding Eq. [4] and Eq. [5], we have:

$$\langle I_z \rangle - \langle I_0 \rangle = \langle S_0 \rangle \left(e^{-(\rho-\sigma)t} - e^{-(\rho+\sigma)t} \right) \quad \dots[6]$$

On the other hand, subtracting Eq. [5] from Eq. [4], we have:

$$\langle S_z \rangle - \langle S_0 \rangle = -\langle S_0 \rangle \left(e^{-(\rho-\sigma)t} + e^{-(\rho+\sigma)t} \right) \dots[7]$$

Eq. [6] indicates a build-up followed by decay (or the opposite), while Eq. [7] describes a monotonic bi-exponential decay.

It may be noted that cross-relaxation rates σ as well as auto-relaxation rates ρ have an inverse sixth power dependence on the distance between the spins I and S , as well as a square dependence on the product of magnetogyric ratios of the two spins. In the two-spin model treated above, the steady state NOE loses the distance information because the ratio (σ/ρ) is involved. However, the transient experiment retains the distance information because the algebraic sum of the two rates is involved in the exponents. In particular, the initial rate of build-up of the I spin transient NOE upon inverting spin S corresponds clearly to 2σ .