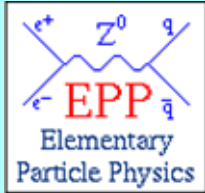


Quantum Numbers

In this section we will cover the following topics:

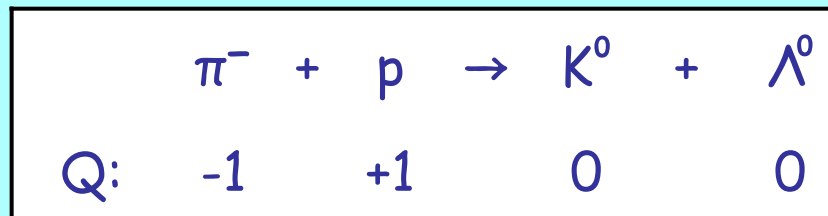
- Electric Charge
- Baryon Number
- Lepton Number
- Spin
- Parity
- Isospin
- Strangeness
- Charge Conjugation



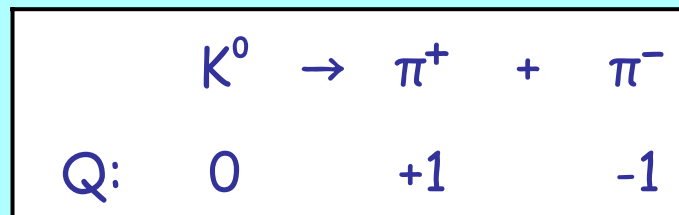
Electric Charge Q

Quantum Numbers are quantised properties of particles that are subject to constraints. They are often related to **symmetries**.

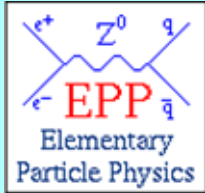
Electric Charge Q is conserved in **all** interactions.



Strong
Interaction



Weak
Interaction



Baryon Number B

Baryons have $B = +1$
Antibaryons have $B = -1$
Everything else has $B = 0$

or

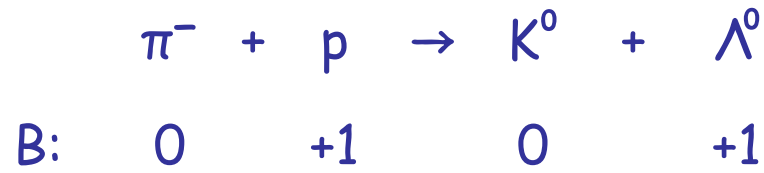
Quarks have $B = +\frac{1}{3}$
Antiquarks have $B = -\frac{1}{3}$
Everything else has $B = 0$

$$\text{Baryons} = qqq = \frac{1}{3} + \frac{1}{3} + \frac{1}{3} = 1 \quad \text{Mesons} = q\bar{q} = \frac{1}{3} + -\frac{1}{3} = 0$$

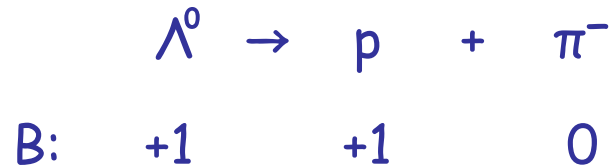
Baryon Number B is conserved in Strong, EM and Weak interactions.

Total (quarks - antiquarks) is constant.

Baryon Number

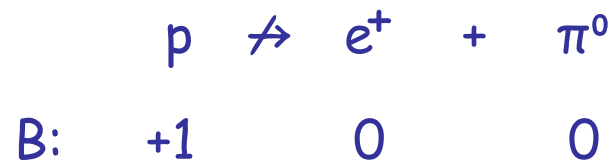


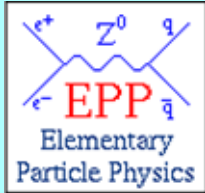
Strong
Interaction



Weak
Interaction

Since the **proton** is the **lightest baryon** it cannot decay if **B** is conserved e.g:





Lepton Number L

Leptons have $L = +1$
Antileptons have $L = -1$
Everything else has $L = 0$

Lepton Number L is conserved in **Strong**, **EM** and **Weak** interactions but is also **separately** conserved within lepton families:

e^- and ν_e have $L_e = 1$ e^+ and $\bar{\nu}_e$ have $L_e = -1$
 μ^- and ν_μ have $L_\mu = 1$ μ^+ and $\bar{\nu}_\mu$ have $L_\mu = -1$
 τ^- and ν_τ have $L_\tau = 1$ τ^+ and $\bar{\nu}_\tau$ have $L_\tau = -1$

L_e , L_μ and L_τ are separately conserved.

Lepton Number



$$\gamma + N \rightarrow e^+ + e^- + N$$

$L_e:$	0	0	-1	+1	0
--------	---	---	----	----	---

Pair
Production



$$\pi^+ \rightarrow \mu^+ + \nu_\mu$$

$L_\mu:$	0	-1	+1
----------	---	----	----

Pion Decay



$$\mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_\mu$$

$L_e:$	0	-1	+1	0
$L_\mu:$	-1	0	0	-1
$L:$	-1	-1	+1	-1

Muon Decay

Lepton Number

✗

	μ^+	\nrightarrow	e^+	+	γ
$L_e :$	0		-1		0
$L_\mu :$	-1		0		0
$L :$	-1		-1		0

Forbidden

OK

L is conserved but neither L_e or L_μ separately.

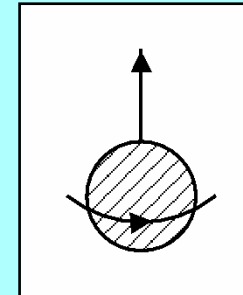
The decay $\mu^+ \rightarrow e^+ + \gamma$ has not been observed and has a "Branching Ratio" $< 10^{-9}$.

Spin S

Spin is an intrinsic property of all particles:

$0\hbar, 1\hbar, 2\hbar, 3\hbar, \dots$ **Bosons**

$\frac{1}{2}\hbar, \frac{3}{2}\hbar, \frac{5}{2}\hbar, \frac{7}{2}\hbar, \dots$ **Fermions**



Spin is like **angular momentum** but a Quantum Mechanical effect.

For spin S there are $2S+1$ states of different S_z (like $2J+1$ in Angular Momentum) e.g.

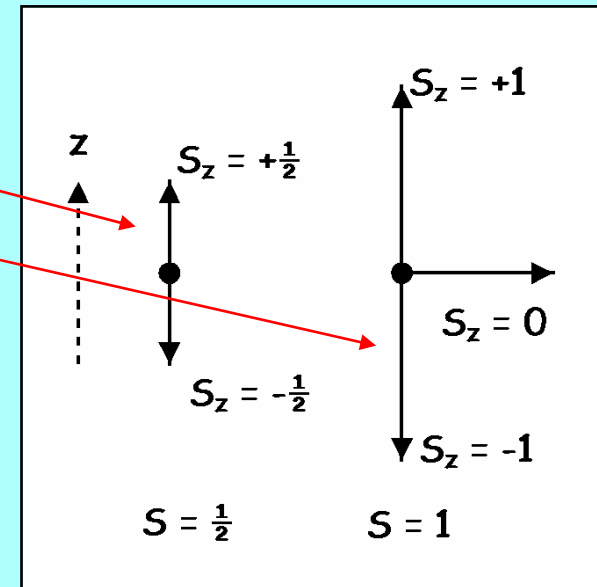
For Spin $S = \frac{1}{2}$, S_z can be $+\frac{1}{2}$ or $-\frac{1}{2}$ (2 states).

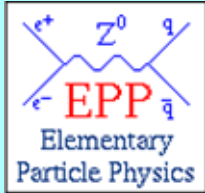
For Spin $S = 1$, S_z can be $+1, 0, -1$ (3 states).

For a process $a + b \rightarrow c + d$ the cross section

$$\sigma \sim (2S_c + 1)(2S_d + 1) \times \text{Other Factors}$$

This can be used to determine the spin of unknown particles.





Parity P

Parity is a Quantum Mechanical concept.

For a wavefunction $\Psi(r)$ and Parity operator P , the **Parity** Operator reverses the coordinates r to $-r$.

$$P \Psi(r) \rightarrow \Psi(-r) = \lambda \Psi(r)$$

$$\text{But } P^2 \Psi(r) \rightarrow \Psi(r) = \lambda^2 \Psi(r)$$

$$\text{i.e. } \lambda^2 = 1 \text{ so that } \lambda = \pm 1$$

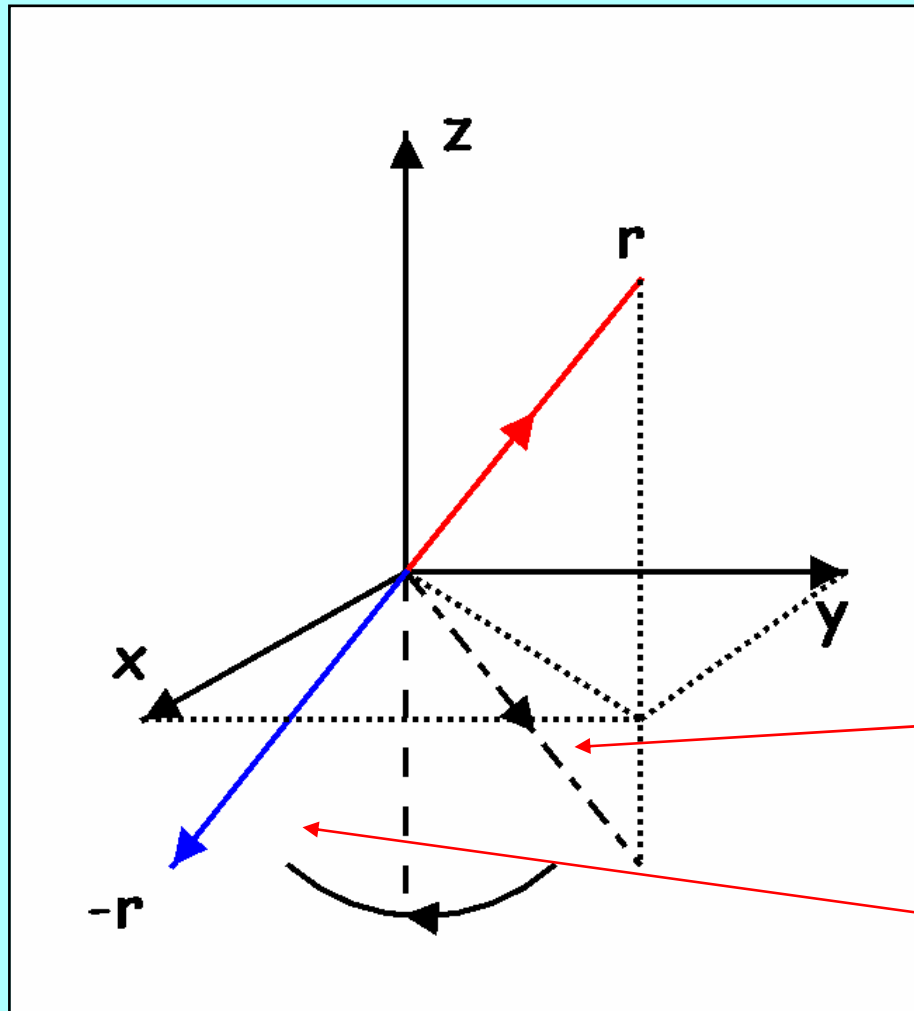
Hence the eigenvalues of Parity are $+1$ (**even**) and -1 (**odd**).

If an operator O acts on a wavefunction Ψ such that Ψ is unchanged

$$O \Psi = \lambda \Psi$$

*Ψ is an **Eigenfunction** of O and λ is the **Eigenvalue**.*

Parity

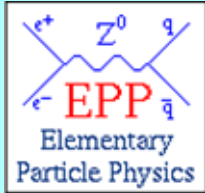


The **Parity** Operator reverses the coordinates r to $-r$.

Equivalent to a **reflection** in the **x-y** plane followed by a **rotation** about the **z** axis

Reflection in x-y plane

Rotation about z axis



Parity

Parity is a **multiplicative** quantum number. The parity of a composite system is equal to the product of the parities of the parts:

$$\Psi = \Phi_A \Phi_B \dots \quad P_\Psi = P_A \times P_B$$

One can show that a state with angular momentum ℓ has parity

$$P = (-1)^\ell$$

For a system of particles:

$$P (\text{overall}) = P (\text{relative motion}) \times P (\text{intrinsic})$$

For **Fermions** $P (\text{antiparticle}) = (-1) \times P (\text{particle})$

For **Bosons** $P (\text{antiparticle}) = P (\text{particle})$

Arbitrarily assign $n, p \rightarrow P = +1$ $\bar{p}, \bar{n} \rightarrow P = -1$

Others determined from experiment (angular distributions)

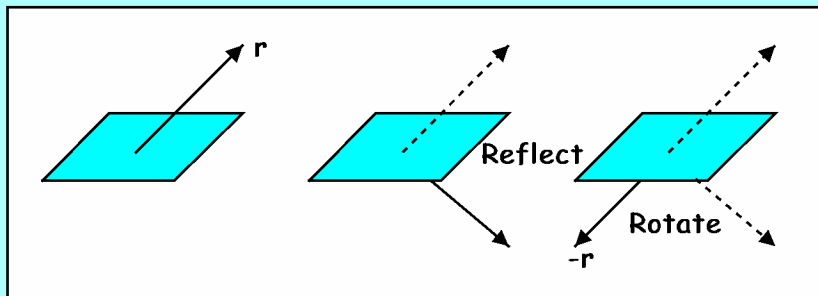
$$\text{Parity of } \pi^+, \pi^-, \pi^0 \rightarrow P = -1$$

Parity

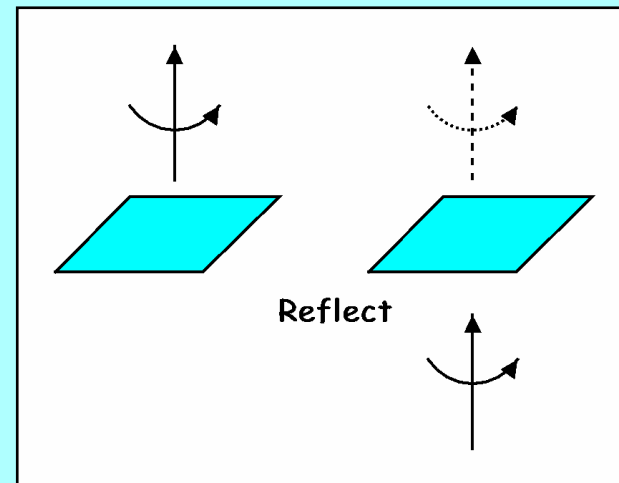
We label mesons by J^P - $\text{Spin}^{\text{Parity}}$ corresponding to how their wavefunctions behave:

- $J^P = 0^-$ Pseudoscalar (Pressure...)
- 0^+ Scalar (Mass, time, wavelength...)
- 1^- Vector (Momentum, position...)
- 1^+ Axial Vector (Spin, angular momentum...)
- 2^+ Tensor

Examples of things that have these properties



Vector $r \rightarrow -r \therefore P = -1$



Axial Vector $r \rightarrow r \therefore P = +1$

Parity is conserved Strong and EM Interactions but NOT Weak.

Isospin I

Used mostly in Nuclear Physics from charge independence of nuclear force $p \leftrightarrow p = n \leftrightarrow n = p \leftrightarrow n$ sometimes called **Isobaric Spin/Isotopic Spin T** (or **t!**).

Isospin is represented by a 'spin' vector **I** with component I_3 along some axis.

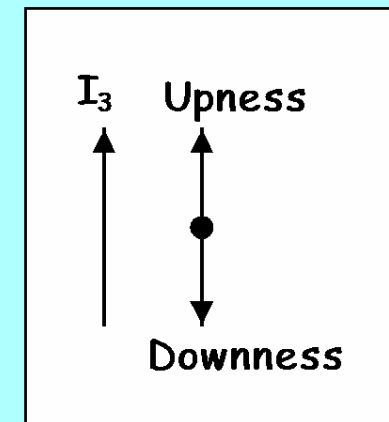
$I = \frac{1}{2}$: p has $I_3 = +\frac{1}{2}$ (\uparrow), n has $I_3 = -\frac{1}{2}$ (\downarrow)

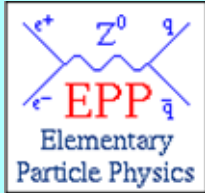
$I = 1$: π^- has $I_3 = -1$, π^0 has $I_3 = 0$, π^+ has $I_3 = +1$

I_3 really only counts the number of **u** and **d** quarks

u has $I_3 = +\frac{1}{2}$ \bar{u} has $I_3 = -\frac{1}{2}$ d has $I_3 = -\frac{1}{2}$ \bar{d} has $I_3 = +\frac{1}{2}$
--

$$p = uud = \frac{1}{2} + \frac{1}{2} + -\frac{1}{2} = \frac{1}{2} \quad \pi^- = \bar{u}d = -\frac{1}{2} + -\frac{1}{2} = -1$$





Isospin

I_3 can be related to charge Q and baryon number B :

$$Q = B/2 + I_3$$

For a proton $B = 1$, $I_3 = +\frac{1}{2}$ and hence $Q = 1$

Since the Strong Interaction doesn't distinguish p from n or u from d , I and I_3 are conserved in Strong Interactions.

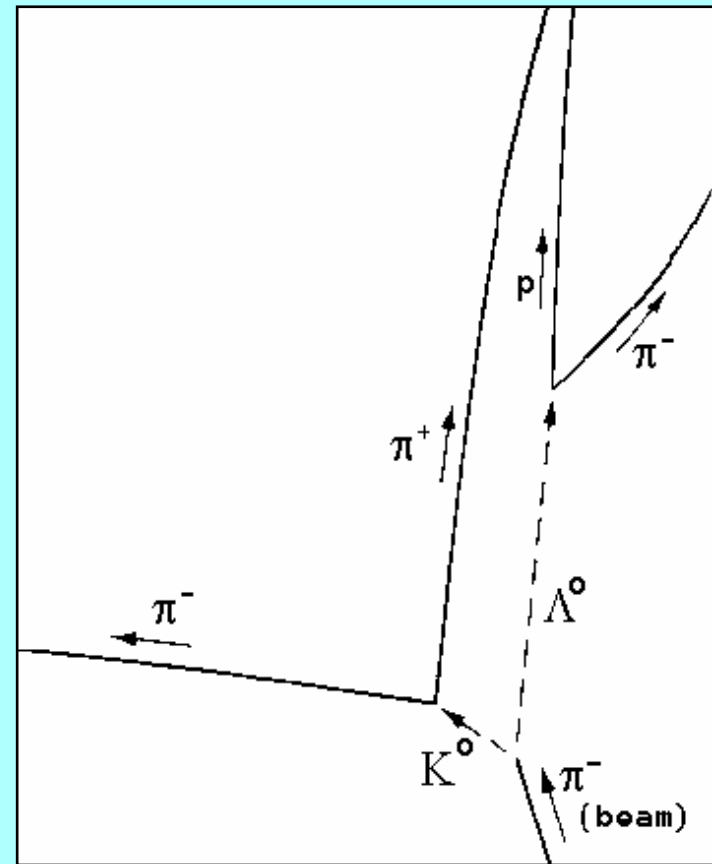
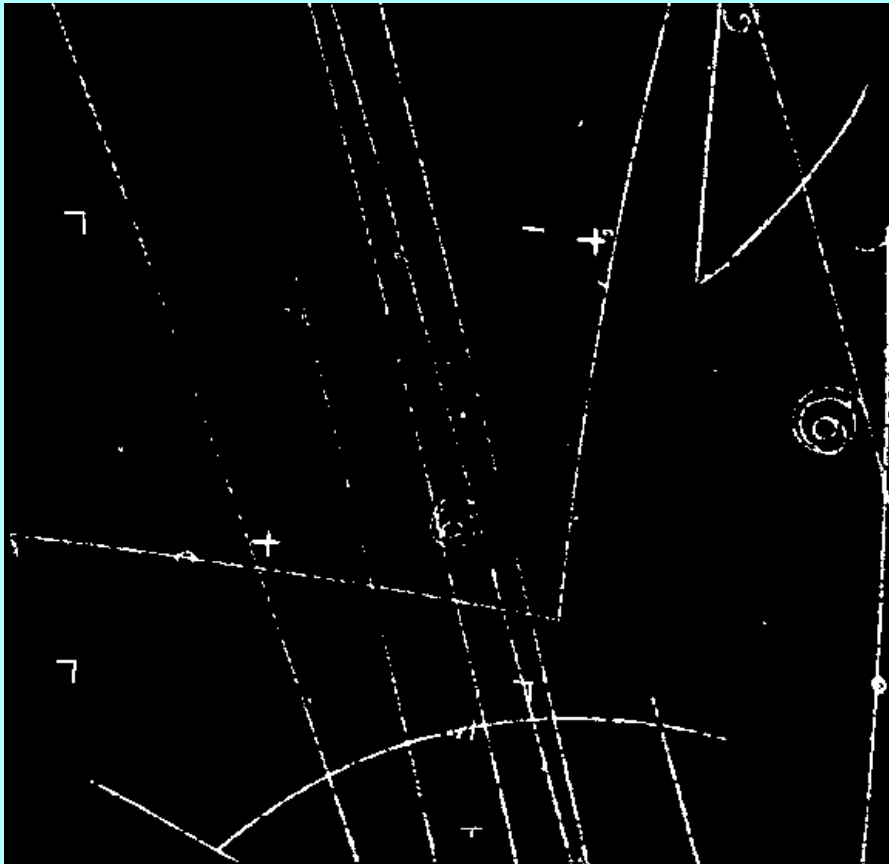
This is equivalent to saying that the number of
 $(u - \bar{u}) - (d - \bar{d}) = \text{constant}$

In Weak Interactions where $u \rightleftharpoons d$, I and I_3 are NOT conserved.

In EM Interactions u and d are not changed but because of the different charges u and d can be distinguished. Hence I_3 is conserved but I is NOT conserved.

Strangeness S

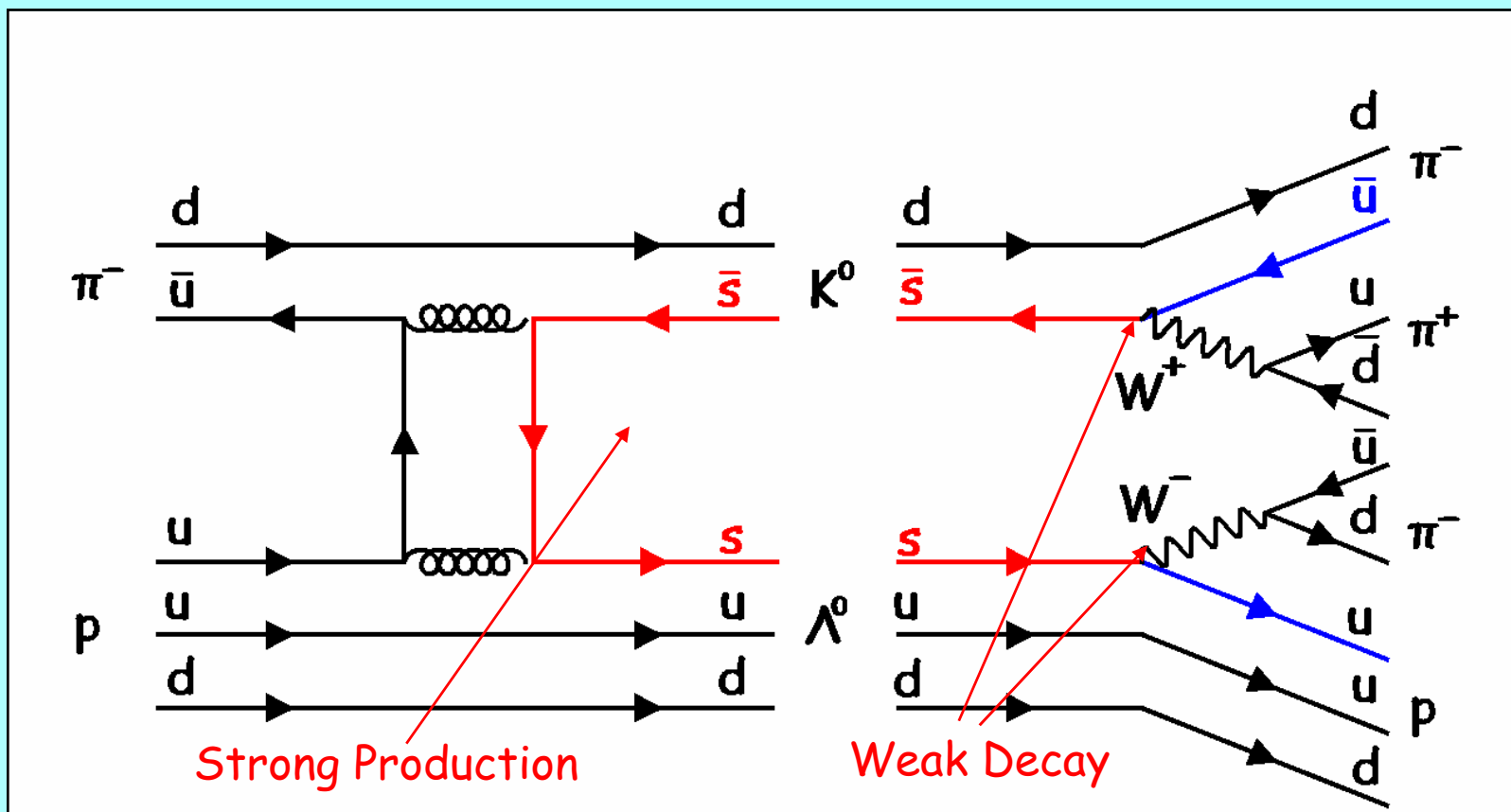
Associated production (via SI) of 'strange' particles $\pi^- + p \rightarrow K^0 + \Lambda^0$



K^0 and Λ^0 'Strange' - decay **weakly** not strongly.

Strangeness

Assume pair production of new quark s and antiquark \bar{s} by Strong Interaction but once produced s and \bar{s} can only decay *weakly*.



Strangeness

Λ^0 is uds

$\bar{\Lambda}^0$ is $\bar{u}\bar{d}\bar{s}$

K^0 is $d\bar{s}$

\bar{K}^0 is $\bar{d}s$

K^+ is $u\bar{s}$

K^- is $\bar{u}s$

The K^0 has an antiparticle the \bar{K}^0 although it is neutral, unlike the π^0 which is its own antiparticle.

Strangeness can be combined with Isospin if

$$Q = I_3 + (B + S)/2 \quad \text{Gell Mann - Nishijima relation}$$

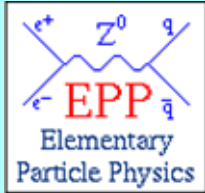
The s quark has strangeness $S = -1$

Strangeness is conserved in Strong and EM Interactions but NOT in Weak Interactions.

Likewise charm, bottom, top quantum numbers.

Strong and EM Interactions do not change quark flavours. Number of $(u - \bar{u})$, $(d - \bar{d})$, $(s - \bar{s})$, $(c - \bar{c})$, $(b - \bar{b})$, $(t - \bar{t})$ constant.

Weak Interaction changes one quark type to another.



Charge Conjugation C

The **Charge Conjugation** operator reverses the sign of **electric charge** and **magnetic moment** (μ).

This implies $\text{particle} \rightleftharpoons \text{antiparticle}$.

$$\begin{array}{ccc}
 \text{proton} & \rightleftharpoons & \text{antiproton} \\
 Q = +e & C & Q = -e \\
 B = +1 & & B = -1 \\
 \mu & & -\mu
 \end{array}$$

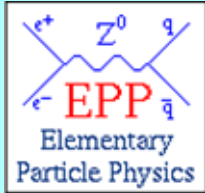
$|X\rangle$ is (*Dirac*) *bra/ket* notation for Ψ_X i.e.
 $|\pi^+\rangle \equiv \Psi_{\pi^+}$

$$C |\pi^+\rangle \rightarrow |\pi^-\rangle \quad \text{Hence } \pi^\pm \text{ not eigenstates of } C.$$

C only has definite eigenvalues for neutral systems such as the π^0 .

$$C |\pi^0\rangle = \lambda |\pi^0\rangle \quad C^2 |\pi^0\rangle = \lambda^2 |\pi^0\rangle = |\pi^0\rangle$$

$$\therefore \lambda = \pm 1$$



Charge Conjugation

EM fields come from moving charges which change sign under Charge Conjugation $\therefore C_\gamma = -1$.

$\therefore n$ photons have $C = (-1)^n$

Since $\pi^0 \rightarrow \gamma\gamma$ this implies $C_{\pi^0} = +1$ (assuming C invariance in EM decays).

Note $\pi^0 \rightarrow \gamma\gamma\gamma$ is then forbidden.

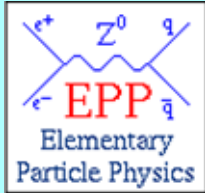
The η (eta) meson (mass $550 \text{ MeV}/c^2$)

$$\eta \rightarrow \gamma\gamma$$

$$\eta \not\rightarrow \gamma\gamma\gamma$$

$$\text{i.e. } C_\eta = +1$$

C is conserved Strong and EM Interactions but NOT Weak.



Summary

Conserved Quantum Numbers

Quantity		Strong	EM	Weak
Charge	Q	✓	✓	✓
Baryon Number	B	✓	✓	✓
Lepton Number	L	✓	✓	✓
Isospin	I	✓	✗	✗
	I_3	✓	✓	✗
Strangeness	S	✓	✓	✗
Parity	P	✓	✓	✗
Charge Conjugation	C	✓	✓	✗