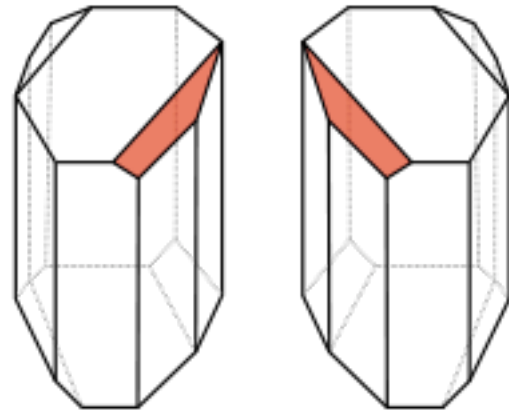# Chirality, particle physics, and "theory space"

Erich Poppitz   U of Toronto

Pasteur (1848):

different rotation of light polarization due
to "chiral" crystals

("sodium ammonium tartrate")

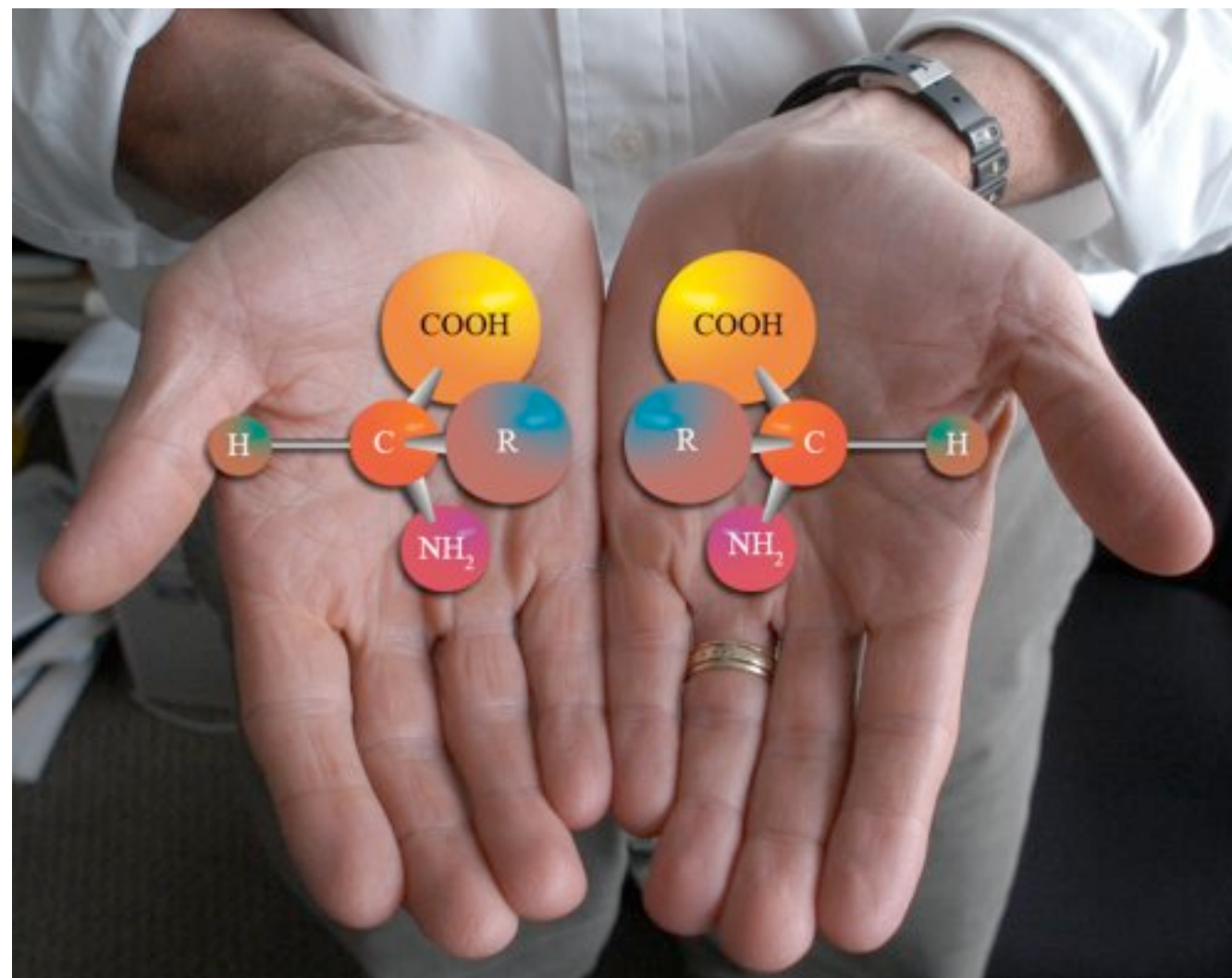- inferred existence of "L-/R- handed" molecules

Lord Kelvin's definition (1904):

*"I call any geometrical figure, or group of points, **chiral**, and say it has **chirality**, if its image in a plane mirror, ideally realized, cannot be brought to coincide by itself."*

"chiral":   an object that is non-superposable
          on its mirror image

χειρ = hand

"handedness"

"*I call any geometrical figure, or group of points, **chiral**, and say it has **chirality**, if its image in a plane mirror, ideally realized, cannot be brought to coincide by itself.*"

Objects with left- or right-handedness commonly occur in the **macroscopic** world.

Some substances only occur naturally in left-handed form, while others occur only in right-handed form - aka "homochirality" (left-handed amino acids and right-handed sugars), already observed by Pasteur

for example, we metabolize only right-handed glucose

Alice to Kitty:

"…Perhaps looking-glass milk isn't good to drink…."

In the **microscopic** world of elementary particles and the fundamental forces between them, however, the symmetry between L-(eft) and R-(ight) holds almost universally.

the four fundamental forces:

- gravity
- electromagnetism ⟶ the one really relevant for chemistry/biology
- weak interactions
- strong interactions

are blind to "chirality" - except the weak interactions

Lee, Yang (theory, 1956)   Wu (experiment, 1957)
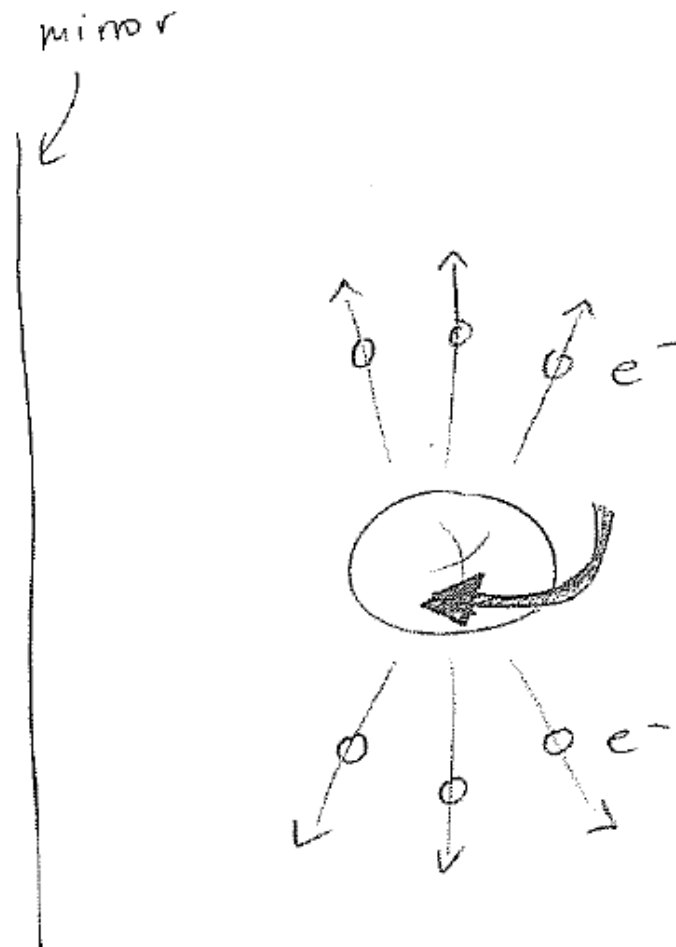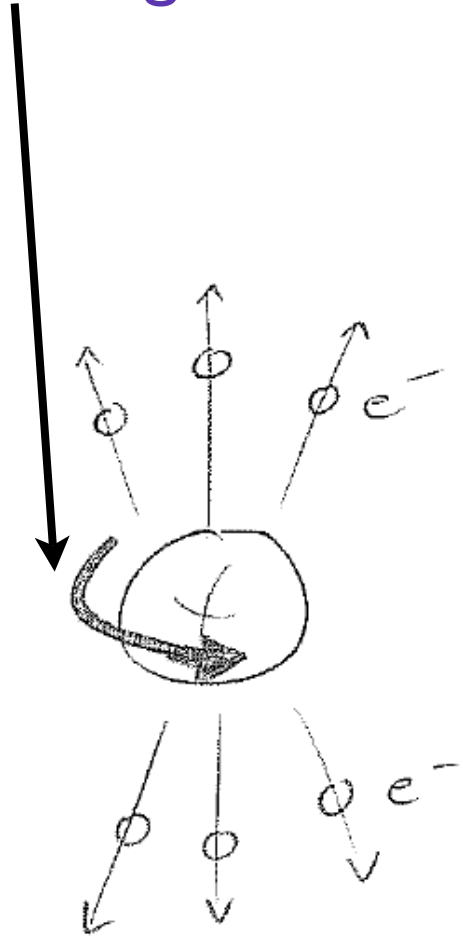
Cobalt ⟶ Nickel + electron + anti-neutrino

27 protons
33 neutrons

28 protons
32 neutrons

spinning nucleus

mirror

in the mirror the nuclei rotate in the opposite direction

symmetry with mirror image would require equal number of electrons going up and down

spin-polarized

Cobalt $\longrightarrow$ Nickel + electron + anti-neutrino

27 protons
33 neutrons

28 protons
32 neutrons

spinning nucleus

mirror

in the mirror the nuclei rotate in the opposite direction

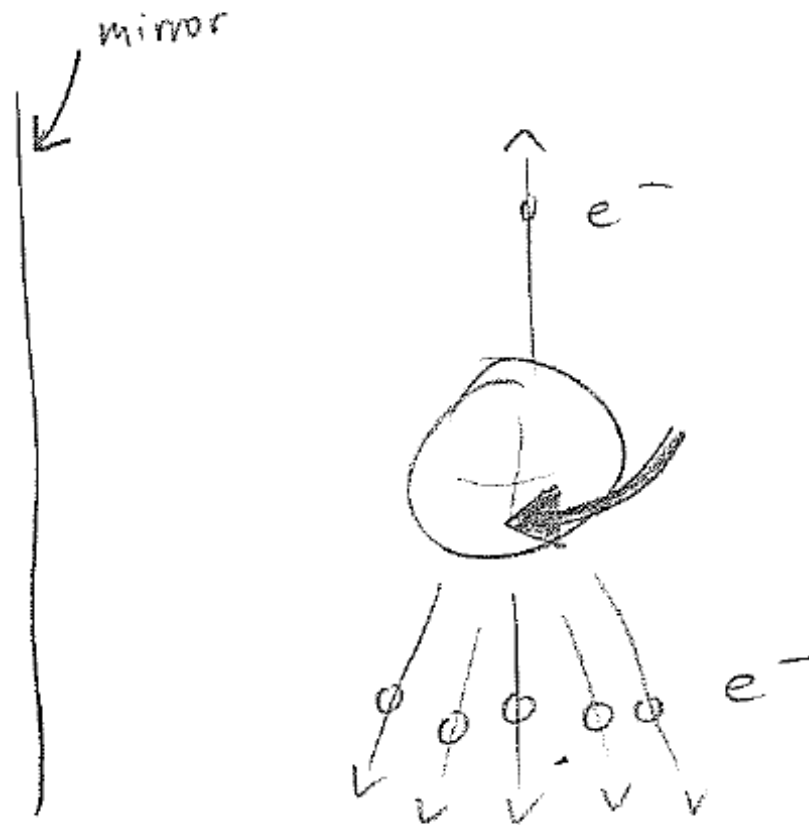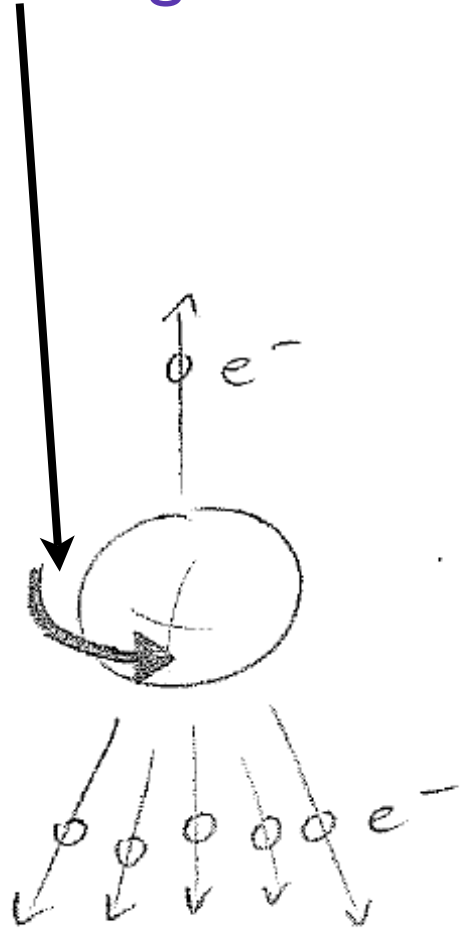symmetry with mirror image would require equal number of electrons going up and down

but this is not what was found!

Wu
(experiment, 1957)

thus, weak interactions are "chiral"

Wu's discovery led Gell-Mann (1958) to postulate the chiral ("V-A") structure of weak interactions, and ultimately, led to the establishment of the Standard Model of particle physics of Glashow, Salam, Weinberg (~1970).

It is not my purpose to review it in great detail today - but explain that it is a "chiral" (L-/R- asymmetric) model.

What, more precisely, do we mean by "chirality" in the particle physics world?

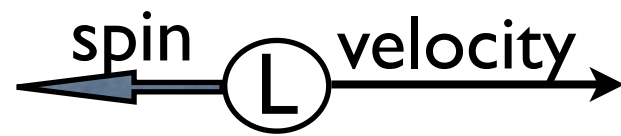What are the "building blocks" of the Standard Model?

gauge bosons of spin-1: photon, gluons, W, Z

fermions of spin-1/2: electron, muon, tau, neutrinos, quarks

electron, muon, tau, quarks, (neutrinos) are all "made of"

"L-handed" fermions

spin ⟵ (L) velocity ⟶

and

"R-handed" fermions

(R) velocity ⟶
spin ⟶
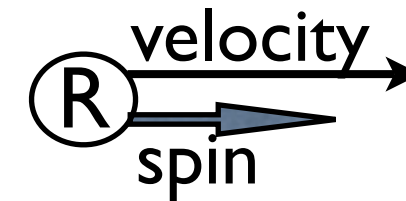
- which are mirror images of each other

massless fermions move with speed of light - projection of spin on direction of motion is a characteristic independent of observer - called the particle's "chirality"

w ⟶

v < c

(L)

v' = v - w

(R)

(picture after Y. Nambu's "Quarks")

massive fermions, on the other hand, do not have definite chirality, which can look different to different observers

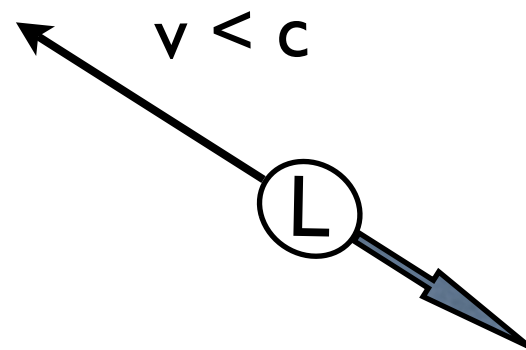e.g., at the fundamental level, the electron is "made of"

left-handed electron
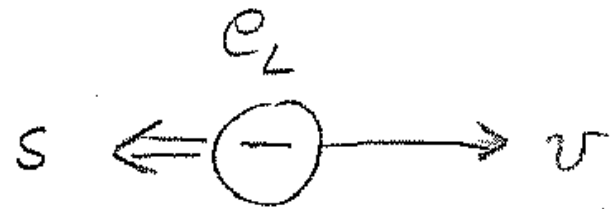
its anti-particle:
right-handed positron
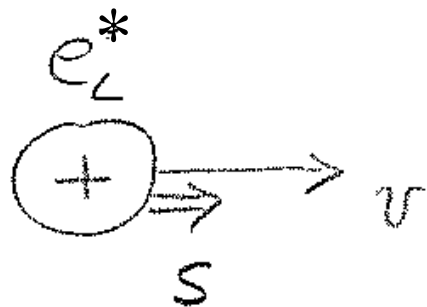
right-handed electron
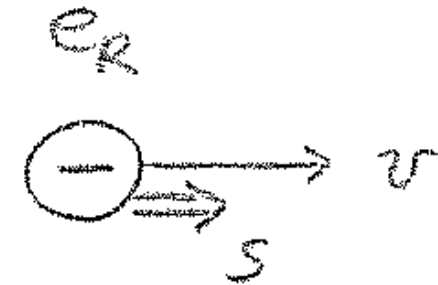
its anti-particle:
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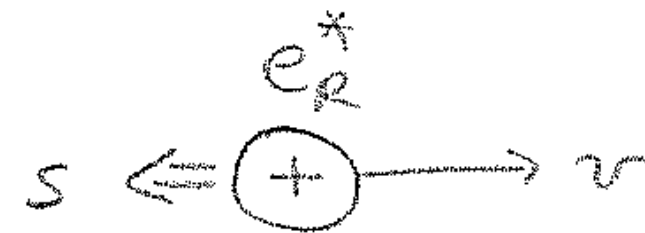
for neutrinos, we do not know yet if R-handed ones exist

for all the rest: muon, tau, quarks (u,d,c,s,b,t) simply put the appropriate value of charge in place of +/-

thus, apart from the neutrinos, the particle content of the Standard Model is L-R symmetric - we call it "vectorlike" or "Dirac"

but, electron not massless - recall: chirality $\longrightarrow$ v = c (massless)

(massive) v < c $\longrightarrow$ ~~chirality~~

there are two ways that a fermion can be "slowed down"
clearly, both must involve chirality violation

"Dirac" fermion mass - leptons, quarks in Standard Model

interaction with the "Higgs field" condensate "flips" chirality

"Majorana" fermion mass

e.g., if R-handed electron didn't exist - charge would
not be conserved - but a possibility for neutrinos

While the spectrum of the SM fermions is L-R symmetric ("vectorlike" or "Dirac"), some of the interactions - the ones responsible for the nuclear beta decay - are L-R asymmetric ("chiral")

W couples only to L!

neutron

proton
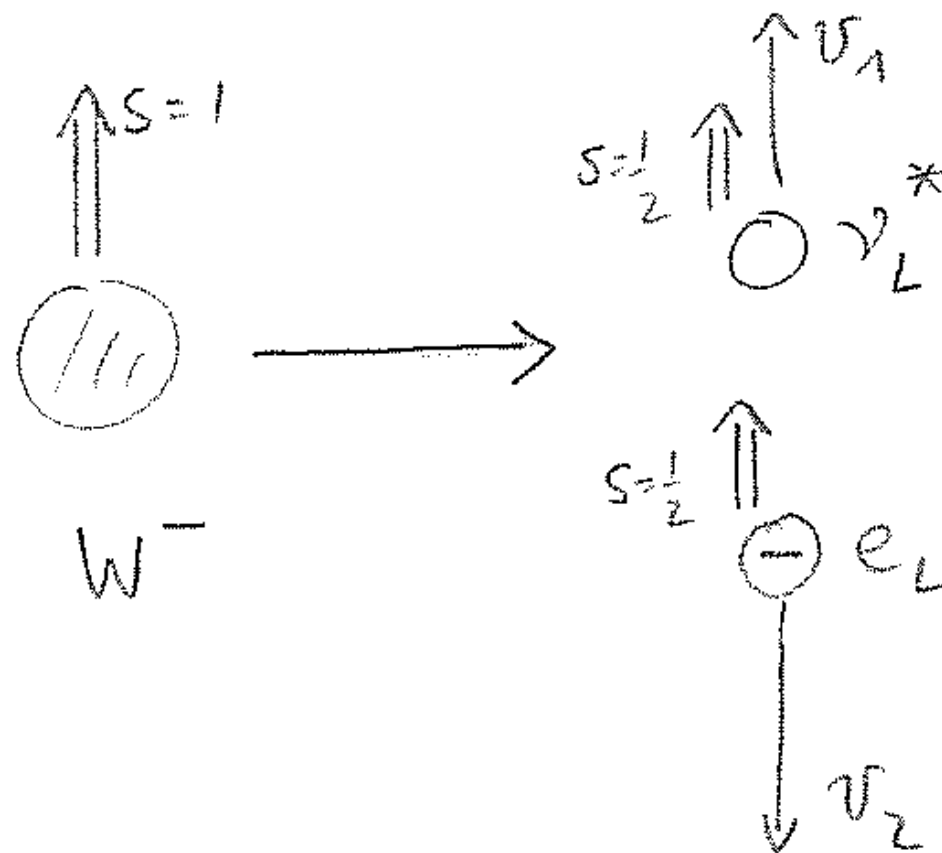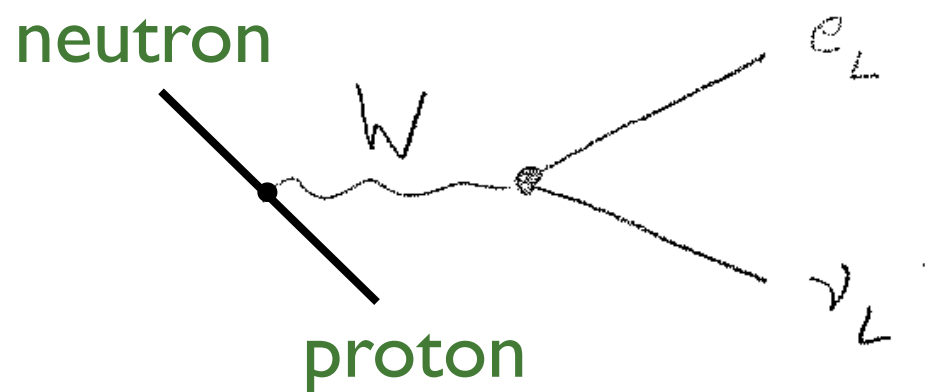
Thus, in Wu's experiment of Cobalt beta decay, only L-electron is produced, so its velocity correlates with spin of W, which in turn "remembers" spin of polarized Cobalt nucleus

(since the nuclear beta decay interactions are very weak, it took many years before parity violation was seen)

In summary, the Standard Model incorporates three of the fundamental forces

- electromagnetism
  "vectorlike"
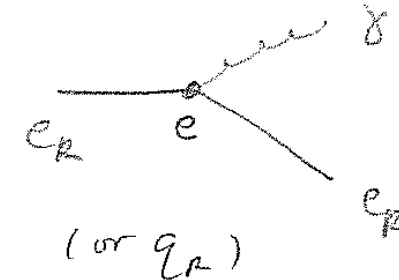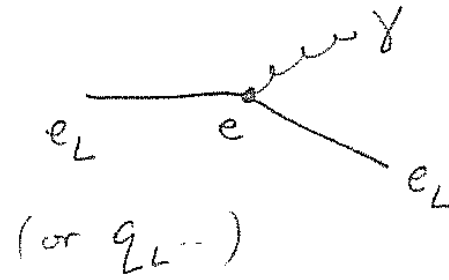
- weak interactions
  "chiral"
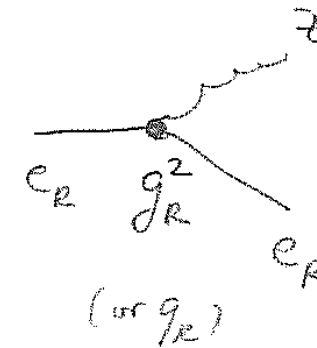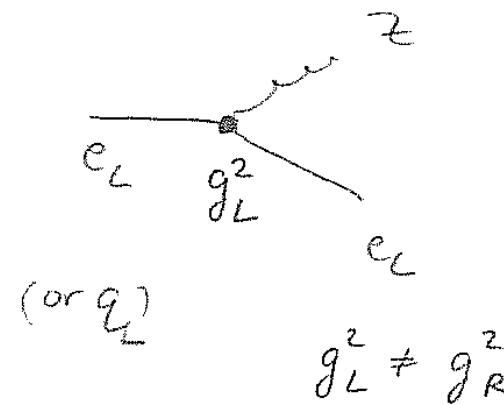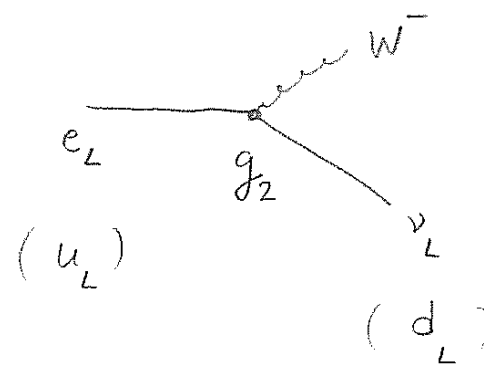
- strong interactions
  "vectorlike"
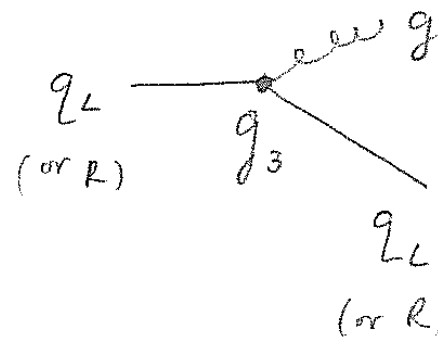
+ gluon self-interactions

as the names may suggest - all but the strong ones are "weak", in the sense that we can study most of their relevant aspects in "perturbation theory" in some small g:

$$H = H_0 + g H_I$$

$$H = H_0 + g\, H_I$$

H_0 — represents something we can solve (free fields)

H_I — represents a "small" perturbation
(use expansion in $g < 1$ to approximate true answer)

"perturbation theory" approach familiar from mechanics/ quantum mechanics - e.g. anharmonic oscillator

V(x)

"perturbation" $\sim g\, x^4$

"unperturbed" $\sim x^2$

x
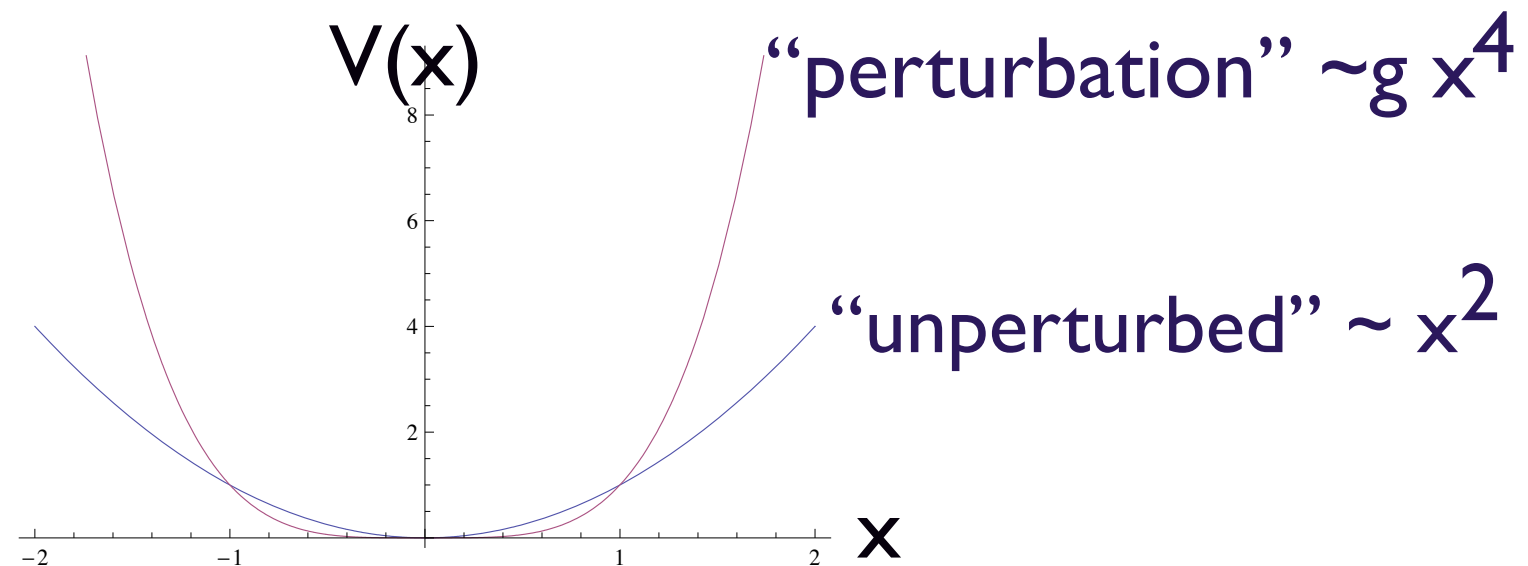
works OK for low-lying states
breaks down for highly-excited ones, even at weak g

For a theorist, the Standard Model is just a collection of **"gauge theories":**

- electromagnetism — "vectorlike", weak

  perturbative
  (largely; e.g., high T...)

- weak interactions — "chiral", weak

- strong interactions — "vectorlike", strong "in the IR"

an arena for a host of non-perturbative techniques:
  - models (quark models, instanton "liquid", SD eqns...)
  - first principles - large-N ideas, lattice
    e.g., solve H without expanding in g: $H = H_0 + g H_I$

The need for such non-perturbative studies of the strong interactions arises because of a peculiarity of their dynamics - "asymptotic freedom" - interactions are weak at short distances but become strong at long-distances (or, as we say, in the "IR" infrared):
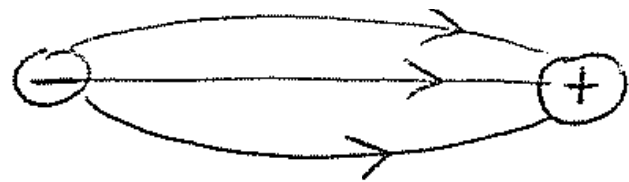
- a flux tube of "glue" ("QCD string") stretches between quarks and gives rise to a linearly-rising "confining" potential between quarks

"quarks" and "gluons" are useful to describe the short-distance behavior of the theory - but fail to capture its large-distance properties -
new "emergent" degrees of freedom become relevant.

Much progress in understanding QCD - the theory of the strong interactions - has been achieved by a variety of techniques - e.g., symmetries, effective field theory, lattice, large-N...

no doubt also aided by the nature's "analogue computer" power (experiment).

"Asymptotic freedom" is a generic property of nonabelian gauge theories. It can lead to a variety of "emergent" IR behavior, which can be quite distinct from that of QCD.

My focus here is on asymptotically free theories different from those describing the strong interactions:

- why is this interesting?
- what is "theory space"?
- how do we study the dynamics?

The reason is that, despite the spectacular agreement of Standard Model predictions with experiment, we're somewhat at a loss...

			experiment	theory	
M_Z	[GeV]	LEP	91.1876 ± 0.0021	91.1874 ± 0.0021	0.1
Γ_Z	[GeV]	LEP	2.4952 ± 0.0023	2.4972 ± 0.0011	-0.9
$\Gamma(\text{inv})$	[MeV]	LEP	499.0 ± 1.5	501.74 ± 0.15	—
σ_{had}	[nb]	LEP	41.541 ± 0.037	41.470 ± 0.010	1.9
R_e		LEP	20.804 ± 0.050	20.753 ± 0.012	1.0
R_μ		LEP	20.785 ± 0.033	20.753 ± 0.012	1.0
R_τ		LEP	20.764 ± 0.045	20.799 ± 0.012	-0.8
R_b		LEP + SLD	0.21644 ± 0.00065	0.21572 ± 0.00015	1.1
R_c		LEP + SLD	0.1718 ± 0.0031	0.17231 ± 0.00006	-0.2
$A_{FB}(b)$		LEP	0.0995 ± 0.0017	0.1036 ± 0.0008	-2.4
$A_{FB}(c)$		LEP	0.0713 ± 0.0036	0.0741 ± 0.0007	-0.8
A_b		SLD	0.922 ± 0.020	0.93477 ± 0.00012	-0.6
A_c		SLD	0.670 ± 0.026	0.6681 ± 0.0005	0.1
A_{LR} (hadrons)		SLD	0.15138 ± 0.00216	0.1478 ± 0.0012	1.6

...the list goes on...

(fairly old transparency, agreement has only improved since; point is to show %-level agreement)

"Dirac" fermion mass - leptons, quarks in Standard Model

e_L e_R e_L
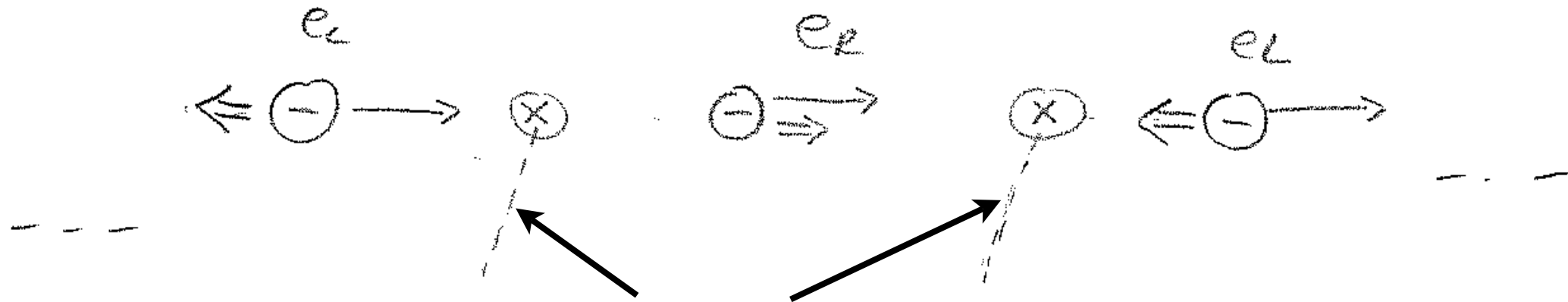
interaction with the "Higgs field" condensate "flips" chirality

the nature of this "thing" is not known

Precisely this unknown Higgs "thing," together with the L-R asymmetry of the electroweak theory links the "generation" of W, Z-boson masses to the fermion (quark and lepton) masses (these span some 12 orders of magnitude below W,Z - why?)

e.g., is "the Higgs" a fundamental scalar field? is it a composite object? is there strongly-coupled dynamics involved? is there supersymmetry? is the theory "natural"? ...

What will the LHC discover?

Blogger Jester says:

Higgs boson.

...

Non-SM Higgs boson.

...

New Beyond SM Particles.

...

Strong Interactions.

...

Dark matter.

...

Little Higgs and friends.

...

Supersymmetry.

...

Dragons.

...

Black Holes.

Blogger Jester says:

Here are my expectations. The probabilities were computed using all currently available data and elaborated Bayesian statistics.

Higgs boson. Probability 80%

...

Non-SM Higgs boson. Probability 50%

...

New Beyond SM Particles. Probability 50%

...

Strong Interactions. Probability 20%

...

Dark matter. Probability 5%

...

Little Higgs and friends. Probability 1%

...

Supersymmetry. Probability 0.1%

...

Dragons. Probability exp(-S(dragon))

...

*Black Holes. Probability 0.1*exp(-S(dragon))...*

What will the LHC discover?

Higgs boson.

...

Non-SM Higgs boson.

...

New Beyond SM Particles.

...

Strong Interactions.

...

Dark matter.

...

Little Higgs and friends.

...

Supersymmetry.

...

In all "scenarios" for "Beyond the Standard Model" physics, new gauge dynamics is invoked, at some scale.

(Unless "supersplit supersymmetry" turns out to be nature's choice.)

In many cases this gauge dynamics is of the "weak interactions" type, e.g., is chiral or L-R asymmetric.

My purpose here is not to discuss and advertise the various "scenarios" - also known as "model building", a separate and very long subject - but to focus on tools to study the "theory space" involved ...

"gauge theory space"   conventional wisdom:

SUSY

- very "friendly" to theorists
 beautiful - exact results

pure YM

- formal `but see www.claymath.org/millennium/`

**QCD-like
(vectorlike)**

- hard, leave it to lattice folks
 limitations: (m, V, $)

**non-SUSY chiral
gauge theories**

- poorly understood strong dynamics
 ...(almost) nobody talks about them anymore

"gauge theory space"   applications:

SUSY
```
superpartner
masses;
SUSY breaking
in chiral
SUSY theories
```

In addition, I believe that one of the most important "applications" of supersymmetry is to teach us about the many "unusual" things gauge field theories could do (or about math! - separate story) - often very much unlike QCD:

QCD-like (vectorlike)
```
W, Z-masses:
"walking" or
"conformal"
technicolor
```

- massless monopole/dyon condensation causing confinement and chiral symmetry breaking
- "magnetic free phases" - dynamically generated gauge fields and fermions
- chiral-nonchiral dualities
- last but not least: gauge-gravity dualities

non-SUSY chiral gauge theories
```
extended technicolor -
fermion mass generation;
quark and lepton compositeness;
& recent very speculative ideas of W, Z, t masses by monopole
condensation
```

"gauge theory space"

nonperturbative tools:

SUSY

- 't Hooft anomaly matching
- "power of holomorphy"
- mass and flat direction "deformations"
- semiclassical expansions
- strings/branes and gauge-gravity dualities
  (much of my older work, as well as recent/
  upcoming Unsal-EP 2009/10)

tools you don't really know whether to trust unless confirmed by other means - experiment or the tools on the left -
the equivalent of
"voodoo QCD"
[Intriligator]

QCD-like (vectorlike)

- most powerful: lattice
- the others mentioned already for QCD

"MAC"
(most attractive channel)
truncated Schwinger-Dyson equations

non-SUSY chiral gauge theories

- 't Hooft anomaly matching
- semiclassical expansions
  (e.g., Unsal-Shifman, Unsal-EP 2008/9)

tools one trusts*

* in some cases, must assume that nonperturbatively string theory exists

"gauge theory space"

nonperturbative tools:

SUSY

- 't Hooft anomaly matching
- "power of holomorphy"
- mass and flat direction "deformations"
- semiclassical expansions
- strings/branes and gauge-gravity dualities
  (much of my older work, as well as recent/upcoming Unsal-EP 2009/10)

QCD-like (vectorlike)

- **most powerful: lattice**
- the others mentioned already for QCD

non-SUSY chiral gauge theories

- 't Hooft anomaly matching
- semiclassical expansions
  (e.g., Unsal-Shifman, Unsal-EP 2008/9)

NOTE: "**most powerful: lattice**" is absent from "**SUSY**" and "**chiral**" parts of theory space

- the lattice is the only nonperturbative definition of general field theories we know of (although some vehemently disagree...)
- can it be useful in cases other than QCD and QCD-like theories?

- the lattice is the only nonperturbative definition of general field theories we know of (although some vehemently disagree...)
- can it be useful in cases other than QCD and QCD-like theories?

•••

In the remaining time, will briefly tell you one (...or two?) stories that I've worked on, unified by the motivational role that string theory has played.

i.) **lattice SUSY and D-brane orbifolds**

- 't Hooft anomaly matching
- "power of holomorphy"
- mass and flat direction "deformations"
- semiclassical expansions
- strings/branes and gauge-gravity dualities

- these are powerful tools - but not "tell all"
- many quantities of interest are not calculable
  - e.g., for dynamical susy breaking model-building of, say, "gauge mediation"
  - or, needless to say, would be nice to know what Seiberg-Witten's massless monopoles/dyons are "made of"

i.) lattice SUSY and D-brane orbifolds

Ex.: 1-d lattice SUSY: D-branes on a Z_N orbifold

open strings

"bulk" - $(9-p)$-dim space
closed strings (gravity)

$(p+1)$-dim
field theory

Dp-brane

(p+1)-dim gauge theory with fermions (and scalars) localized on the brane's world volume

Dp-brane — "point" in $(9-p)$-dim

p-brane

"winding open strings"

✳ 2 of $(9-p)$-dim shown
✳ a "Z_N orbifold"

$\frac{2\pi}{N}$
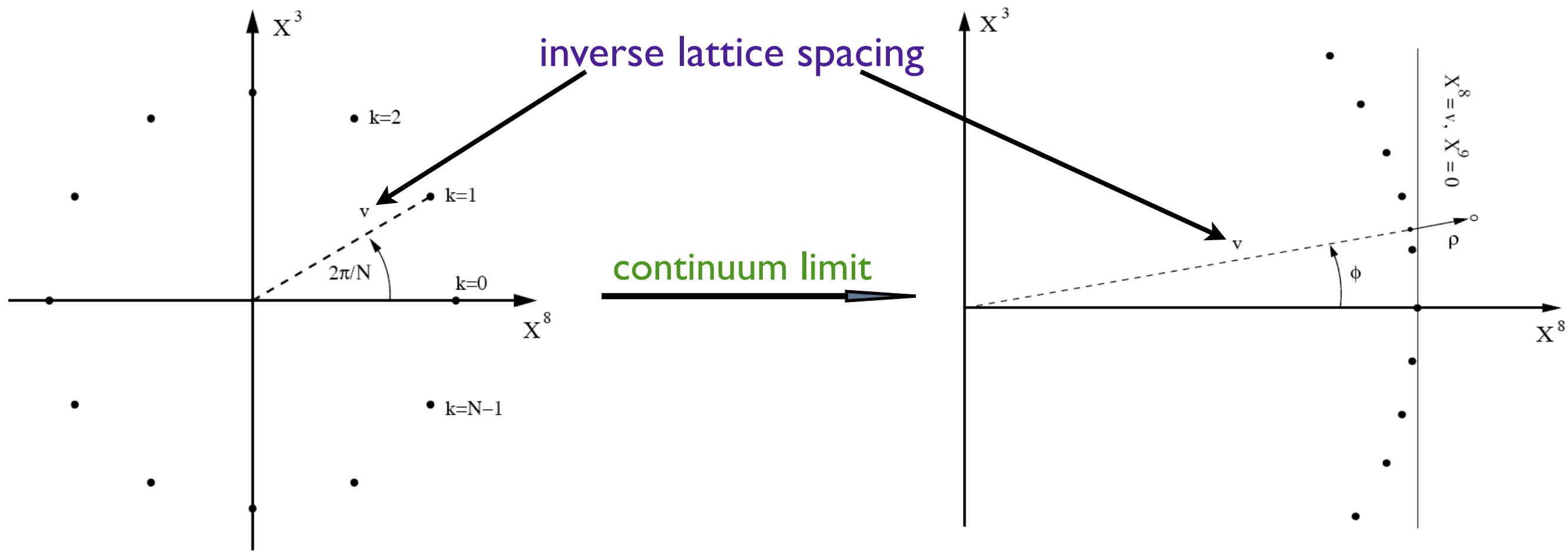
tip of cone ≡ orbifold fixed point
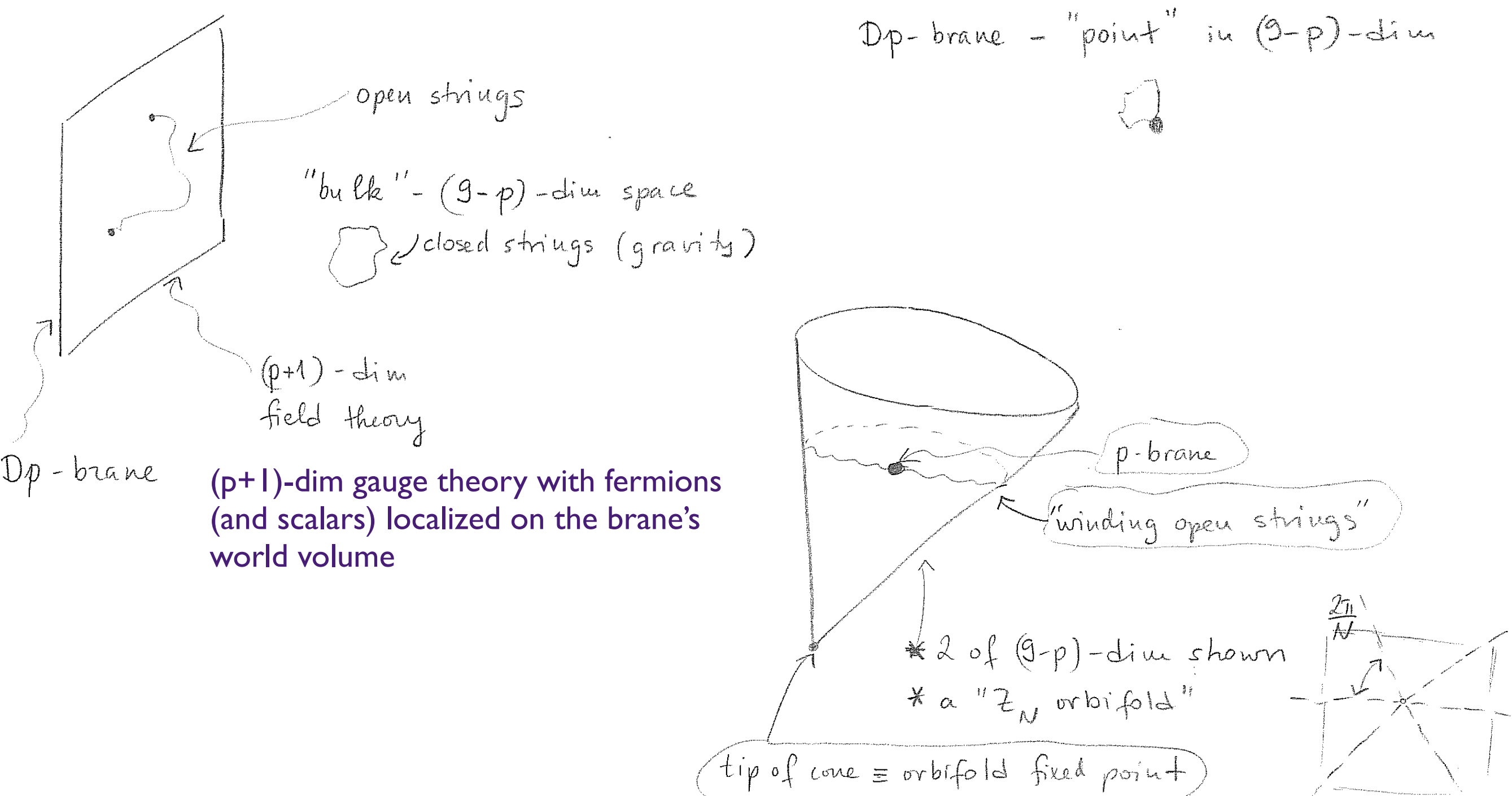
i.) lattice SUSY and D-brane orbifolds

Douglas, Moore `96
Lykken, Trivedi, EP `97, `98
Cohen, Kaplan, Katz, Unsal `02
Giedt, Rozali, EP `03
Kaplan, Unsal `06
(Giedt, EP `04,
 Giedt, Koniuk, Yavin, EP `04
 and many other non-brane papers)

Ex.: 1-d lattice SUSY: D-branes on a Z_N orbifold

number of lattice sites

inverse lattice spacing

continuum limit

"winding" open strings between images = lattice momentum modes (1dim Wilson r=1)

lattice theory preserves supersymmetry, which does not involve translations in the lattice directions - but enhances to full SUSY in continuum limit

i.) lattice SUSY and D-brane orbifolds

Lykken, Trivedi, EP `97   realized this 1-dim lattice (of N=2 SYM)
"Chiral gauge theories from D-branes"   in retrospect! - had a non-lattice motivation

Cohen, Kaplan, Katz, Unsal `02   constructed 2,3,4-dim lattices with SUSY

studies of quantum continuum limit, anomalies, fine-tuning issues, etc....

Giedt, Rozali, EP `03;  Giedt, EP `04; Giedt, Koniuk, Yavin, EP `04
and many other brane and non-brane papers

moral/summary of what transpired since:

- branes on orbifolds motivated the construction of (non)gauge actions that preserve some SUSY; quite nontrivial to come up with in field theory (a version of Kaehler-Dirac fermions appears; all unified under "topological twisting" umbrella - Giedt, EP `04; Unsal `06)

•• actually useful for simulations in low-dim SUSY systems (e.g., Catterall et al, KEK group) to, e.g., study aspects of low-dim AdS/CFT

••• all of this works only for vectorlike (extended) SUSY; no clue about chiral!
   (but recall interest: superpartner masses; supersymmetry breaking in
    chiral SUSY theories?)

ii.) non-SUSY chiral lattice gauge theories?

*A fascinating development also influenced by strings (extra dimensions) ultimately led to the understanding of how to **explicitly** realize exact chiral symmetries on the lattice.*

Ginsparg, Wilson `82- forgotten! - 'til `97
Callan, Harvey, `85
Kaplan `92
Narayanan, Neuberger `94
Neuberger `97
P. Hasenfratz, Laliena, Niedermayer `98
Neuberger `98
Luscher `98-`99

With *few exceptions*, however, only L-R symmetric (vectorlike) theories with chiral symmetry are nonperturbatively defined on the lattice
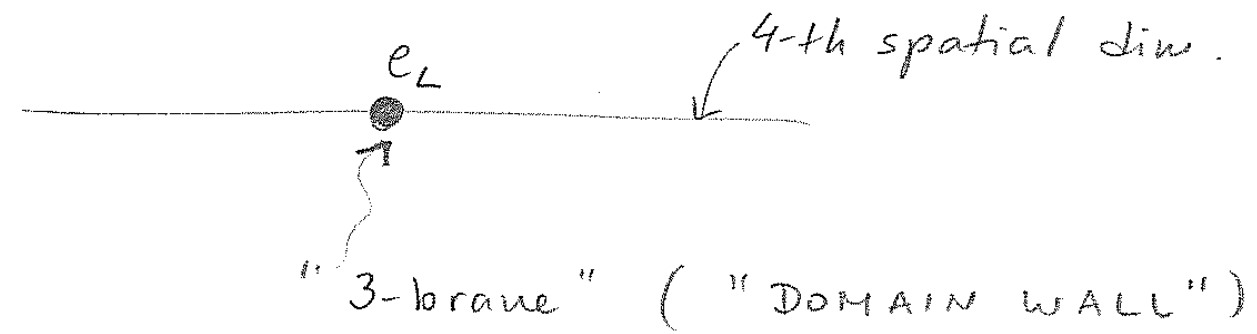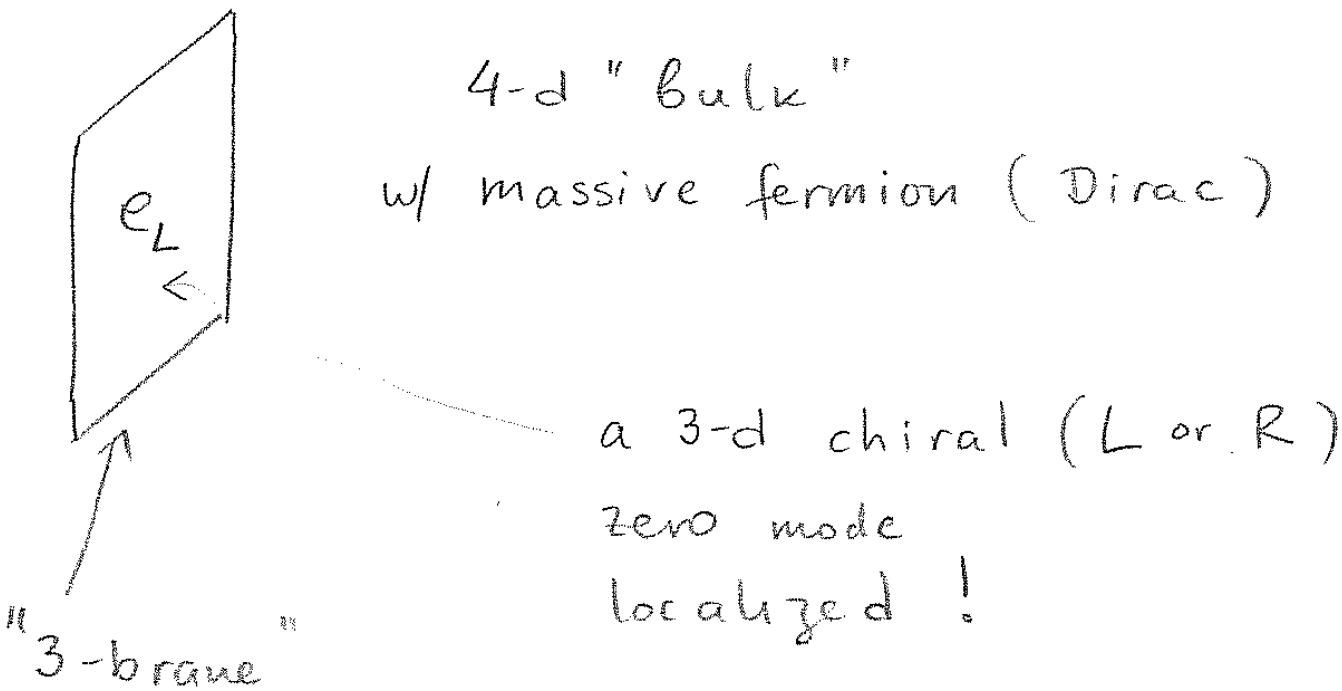
- great progress for QCD and QCD-like theories, even if it takes $$$ to use.

Where is the challenge?
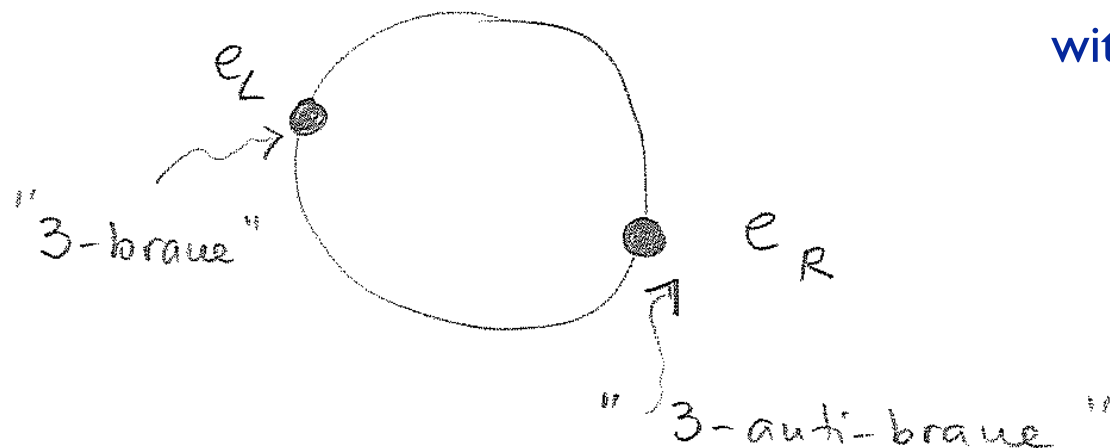
ii.) non-SUSY chiral lattice gauge theories?

Kaplan's `92 construction:

i.) an infinite 4-th(space)dimension
where Dirac fermion propagates

ii.)chiral zero mode localized on
3-dim defect

4-d "bulk"
w/ massive fermion (Dirac)

e_L

a 3-d chiral (L or R)
zero mode
localized !

"3-brane"

e_L

4-th spatial dim.

"3-brane" ("DOMAIN WALL")

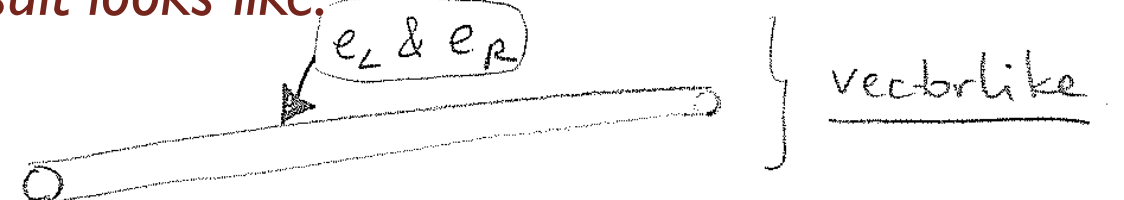to understand the "infinity" - make dimension compact -
but this brings in "anti-wall" with mirror (R-handed) fermion -
-- end up with vectorlike theory.

with Neuberger `97 operator etc...

e_L

"3-brane"

e_R

"3-anti-brane"

to an observer who doesn't resolve extra dimension,
the end result looks like:

e_L & e_R

vectorlike

(to the experts: apologies for somewhat misleading picture)

ii.) non-SUSY chiral lattice gauge theories?

Where is the challenge? - get rid of the mirror partners

Luscher `98-`99 (also, in part, Neuberger `98) – implicitly did so by showing how to define chiral partition function for anomaly free U(1). Mathematically, generalization to, say, SU(N), poses a difficult problem.

Can this hard (for me) math be avoided?

In string theory, constructing chiral gauge theories often proceeds from higher-dim theories by "orbifolding" (or a "blown-up" version) of L-R symmetric (vectorlike) theories - by "erasing" the mirrors, or, more colloquially speaking, decoupling them... (OFTEN CALLED "GEOMETRIC ENGINEERING")

e.g., Lykken, Trivedi, EP `97
+ many others

In each case, more structure than a simple domain wall is required and anomaly cancellation is always built-in the construction (which otherwise "falls apart").

Even stranger things involving chirality can happen: chiral-nonchiral "transitions" as parameters (moduli) are varied (Douglas, Zhou `04) **& chiral-nonchiral Seiberg dualities in N=1 SUSY** (Pouliot; Strassler and Pouliot `96).

These hints are perhaps useful to think about... for now, follow a simple-minded, "experimentalist" - or "engineering"- approach, which is, however, testable.

ii.) non-SUSY chiral lattice gauge theories?

Where is the challenge? - get rid of the mirror partners

"light" fermions
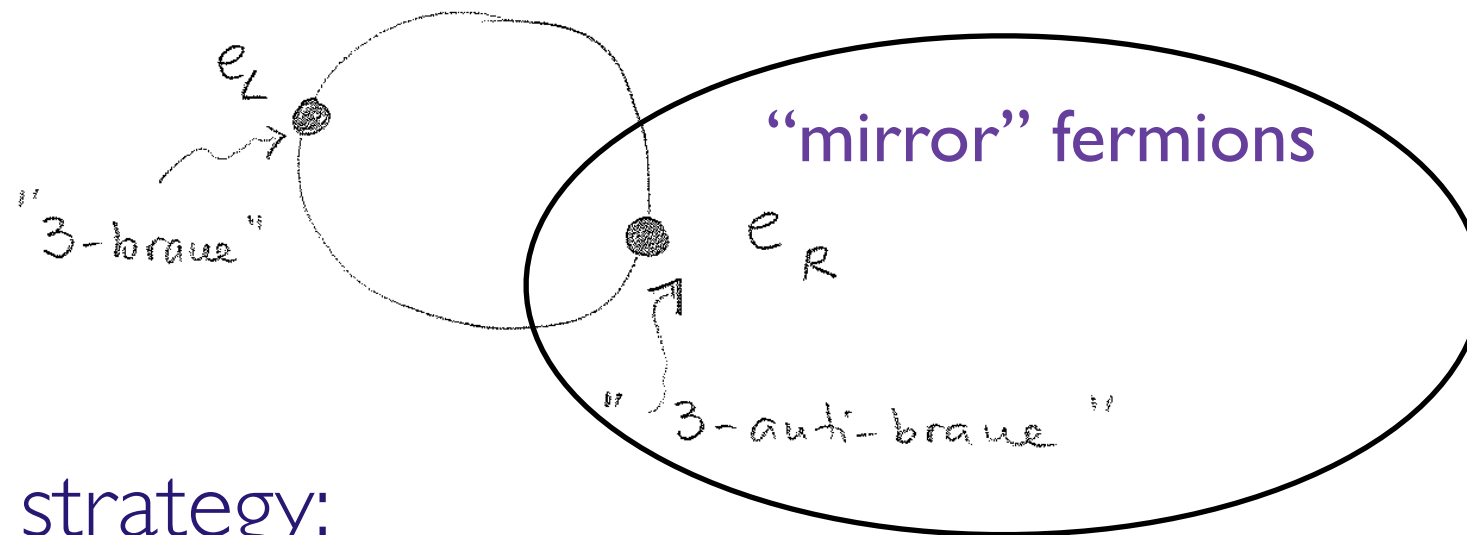
Bhattacharya, Martin, EP `06
Giedt, EP `07
Shang, EP `07
Shang, EP `09

earlier idea of similar flavor
(not pursued & independent)
Creutz, Rebbi, Tytgat, Xue `96

our strategy:

add gauge invariant multi-fermion (or Yukawa) interactions between mirror fermions

break all anomalous and anomaly free global symmetries of the mirror
since 't Hooft anomaly matching on the lattice applies to these interactions
(Shang, EP `07,`09)

drive the mirror interactions into a strong-coupling symmetric phase (exists: Giedt, EP `07)

at strong mirror couplings expect all mirrors to decouple, as there is no symmetry
reason for light mirrors at strong coupling - *a field theory "orbifold" realization!?*

note that this story has all the flavor of Eichten, Preskill `86 with one crucial difference:
we now have exact chiral symmetries and anomaly matching arguments on the lattice!

ii.) non-SUSY chiral lattice gauge theories?

I told you our strategy – but "Does it work?"

apart from symmetry arguments, must

hang on... "experiment" is slow (mainly by Joel Giedt, on BlueGene/L, Rennselaer PI)

but we have not seen reasons to give up -

so far, our results on anomaly matching go in the right direction:
min number of fermions needed to match anomalies remains massless
when anomalous mirrors (cheap!) are studied –

- but, at the same time, we don't know if we have
succeeded or "not failed", yet!

I have not failed. I've just found 10,000 ways that won't work.

Thomas A. Edison

CONCLUSION - studying "gauge theory space" is:

a.) of theoretical interest

   and

b.) it may help us understand what the LHC will be trying
   to tell us about the short-distance properties of nature

SUSY

pure YM

**QCD-like
(vectorlike)**

**non-SUSY chiral
gauge theories**

?

⟷

Higgs boson.
...
Non-SM Higgs boson.
...
New Beyond SM Particles.
...
Strong Interactions.
...
Dark matter.
...
Little Higgs and friends.
...
Supersymmetry.
...

(an aside:

microscopic world: only weak interactions "chiral," but "chirality" is common in the **macroscopic** world

it's natural to ask: could there be any connection?

(I think) it is hard to imagine how the small chirality violation in the weak interactions could bias chemical/bio processes to lead to the observed "homochirality" of macro world

- recall electromagnetism, most relevant for chemistry/bio is L-R symmetric

- energy splitting between L- vs. R- molecules due to weak interactions is tiny $E/kT \sim 10^{-17}$ (kT = 0.025 eV, splitting between l- and d-amino acids)

but, needless to say, speculations trying to amplify this tiny effect exist
- and are subject of debate... ("aminoacidgenesis" vs "baryogenesis"?)

- end of aside **)**