

Nuclear “pasta” in supernova cores

Hidetaka Sonoda
University of Tokyo

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Collaborators

G. Watanabe, K.Sato, T. Takiwaki, K. Yasuoka, T. Ebisuzaki

Outline

- Introduction
- Reproduction of nuclear pasta by QMD
- Neutrino opacity of nuclear pasta
- Summary

Introduction

Supernova and nuclear pasta

Supernova Explosion

Not yet perfectly understood

Understanding interaction between neutrino and matter

One of promising ingredient for success of SN

One of yet understood related physics is ...

Nuclear shape transitions in dense matter ($\sim 10^{14}$ g/cc) (**Nuclear pasta**)

Impact on scattering cross section of neutrino and matter

(by factor of several tens or hundreds relative to uniform matter)

Very important for more accurate simulations of SN

What is nuclear “pasta” ?

Non-spherical nuclei in dense matter $\sim 10^{14}$ g/cc
Caused by surface and Coulomb energy

Sphere

Rod

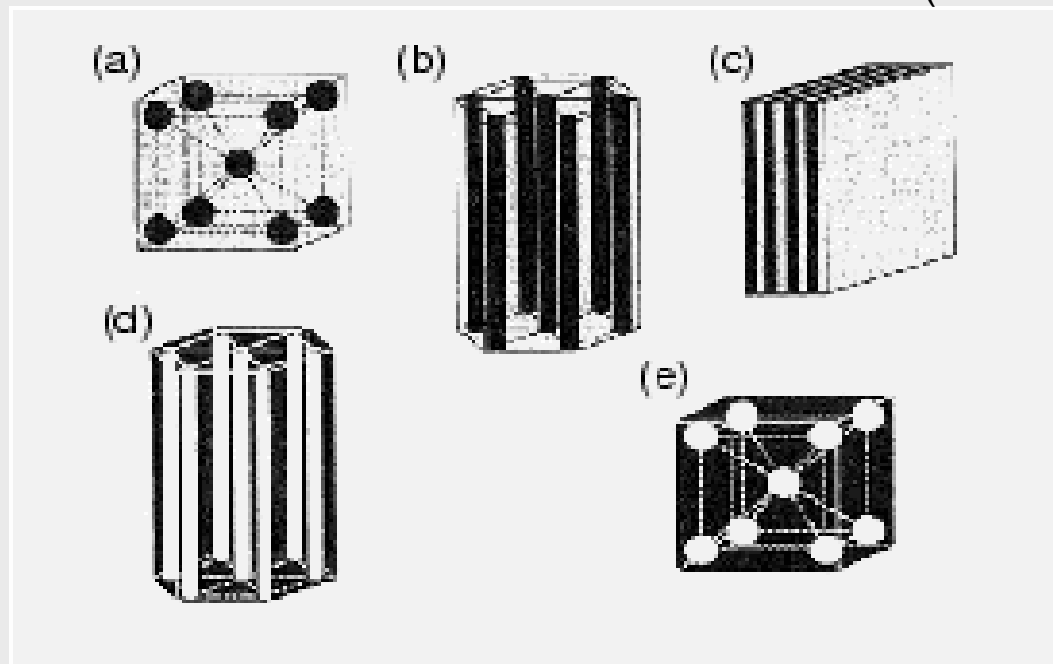
Slab

Rod-like Bubbles

Spherical Bubbles

Uniform Nuclear Matter

(Ravenhall et al. 1983, Hashimoto et al. 1984)



Meatball
Spaghetti
Lasagna
Anti-spaghetti
Cheese
“Pasta” Phases

(K.Oyamatsu, Nucl.Phys.**A561**,431(1993))

Estimated abundance of “pasta” in SN cores

Conditions for non-spherical nuclei

(Pethick & Ravenhall (1995))

Volume fraction of nuclei $> 1/8$

spherical nuclei become unstable

(Bohr-Wheeler condition; $E_{\text{coul}} > 2E_{\text{surf}}$)

Possible region of “pasta” phases

In simulation of core-collapse SN with typical parameters

These region can occupy 10%~20% of core mass

Our study

1. Reproduction of pasta phases

Pasta phases really appear in supernova cores ?

How does uncertainties of nuclear force affect ?

2. Effects of “pasta” phases in supernova cores

How is neutrino opacity changed by pasta?

**To approach to these problems,
we use Quantum Molecular Dynamics (QMD)**

1. Reproduction of Pasta

Method

Why QMD ?

Quantum Molecular Dynamics (QMD)

We can observe how nuclei melts into uniform

QMD is suitable for observing structural transitions

- 1. No assumptions on nuclear shapes**
- 2. Natural treatments of thermal fluctuations**

Framework of QMD

Model Hamiltonian (Maruyama *et al.* Phys. Rev. C 57, 655 (1998))
(Chikazumi *et al.* Phys. Rev. C 63, 024602(2001))

$$\mathcal{H} = \sum_i \frac{P_i^2}{2m_i} + V_{\text{Pauli}} + V_{\text{Skyrme}} + V_{\text{sym}} + V_{\text{surface}} + V_{\text{MD}} + V_{\text{Coulomb}}$$

The diagram shows the Hamiltonian equation with color-coded boxes around each term. Lines connect these boxes to labels below the equation:

- $\sum_i \frac{P_i^2}{2m_i}$ (red box) connects to **Kinetic energy** (red box)
- V_{Pauli} (green box) connects to **Pauli Potential** (green box)
- $V_{\text{Skyrme}} + V_{\text{sym}} + V_{\text{surface}} + V_{\text{MD}}$ (blue box) connects to **Nuclear force** (blue box)
- V_{Coulomb} (purple box) connects to **Coulomb energy** (purple box)

Determination of parameters by reproducing ...

- Kinetic energy of ideal Fermi gas
- Properties of symmetric nuclear matter at saturation density
- Optical potential of pN elastic scattering
- Binding energy and rms. radius of heavy stable nuclei

Time evolution calculated by solving EoM of QMD

Simulations at zero and nonzero temperatures

Simulation settings

2048 or 10976 or 16384 nucleons (n,p,e system)

Periodic boundary condition

Proton fraction $x=0.3$

Uniform background of electrons

Zero temperature: frictional relaxation

$$\begin{aligned}\dot{\mathbf{R}}_i &= \frac{\partial \mathcal{H}}{\partial \mathbf{P}_i} - \xi_R \frac{\partial \mathcal{H}}{\partial \mathbf{R}_i} \\ \dot{\mathbf{P}}_i &= -\frac{\partial \mathcal{H}}{\partial \mathbf{R}_i} - \xi_P \frac{\partial \mathcal{H}}{\partial \mathbf{P}_i}\end{aligned} \quad (\text{in timescale of } \sim O(10^4)\text{fm/c})$$

Nonzero temperature: Nose-Hoover thermostat

Result

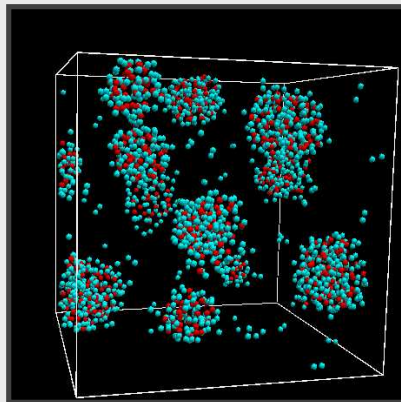
Pasta phases at zero temperature

Cooling of Hot nuclear matter (~10 MeV) down to 0.1 MeV

Snapshots calculated for model 2

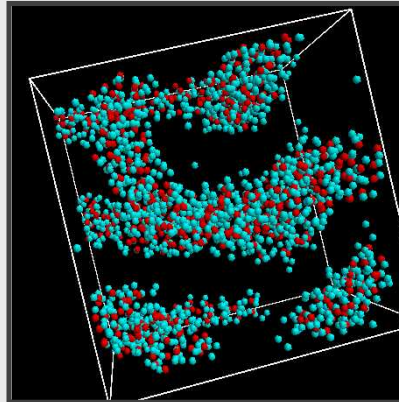
Sphere

0.100 ρ_0



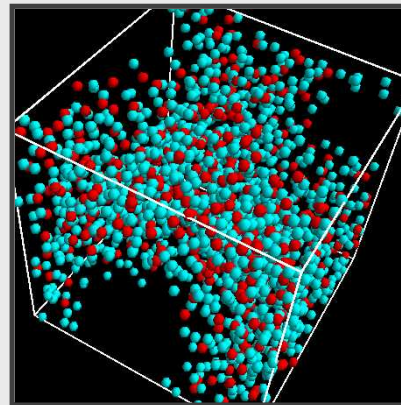
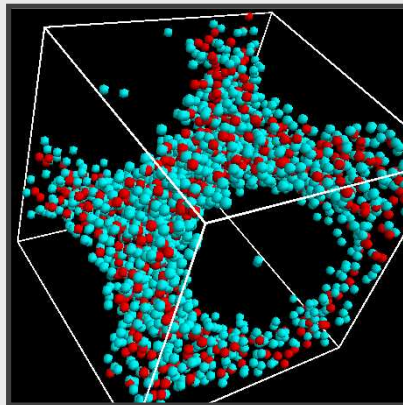
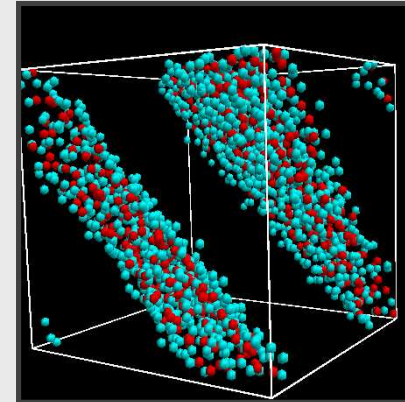
Rod

0.200 ρ_0



Slab

0.393 ρ_0



Red: Proton

Blue: Neutron

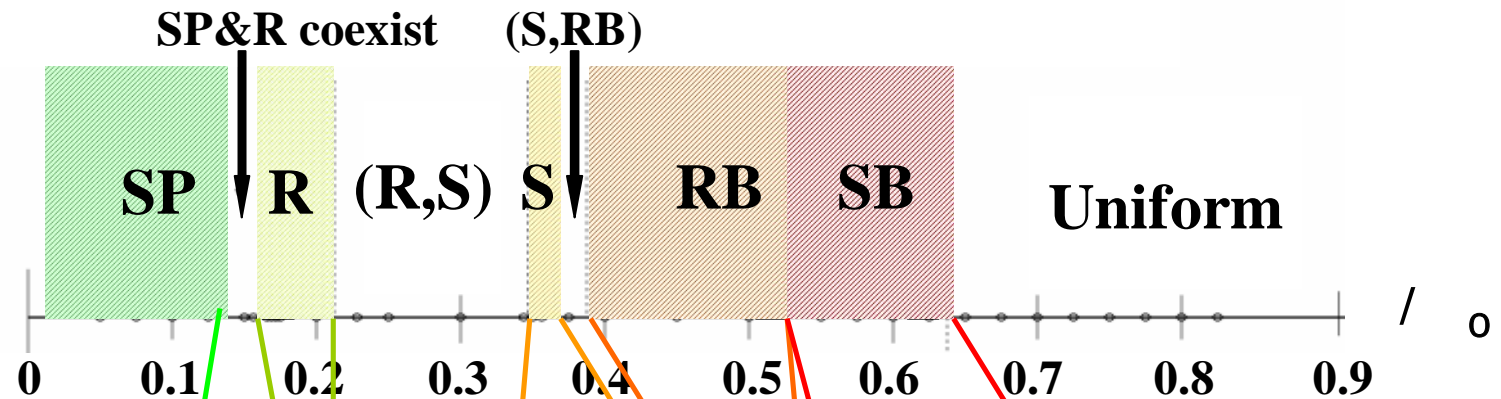
$\rho_0 = 0.168 \text{ fm}^{-3}$
(Nuclear density)

Rod-like Bubbles 0.490 ρ_0

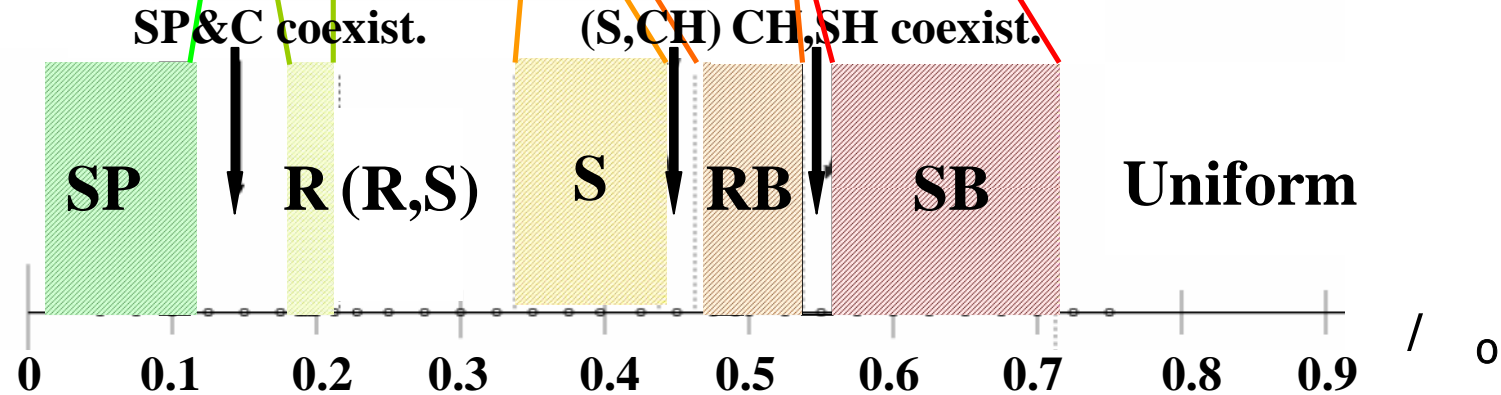
Spherical Bubbles 0.575 ρ_0

Phase diagram at zero-temperature

Model 1



Model 2



SP: sphere S: slab SB: spherical bubble
 R: rod RB: rodlike bubble (,): intermediate

Sphere

Rod

Slab

Rod-like bubbles

Spherical bubbles

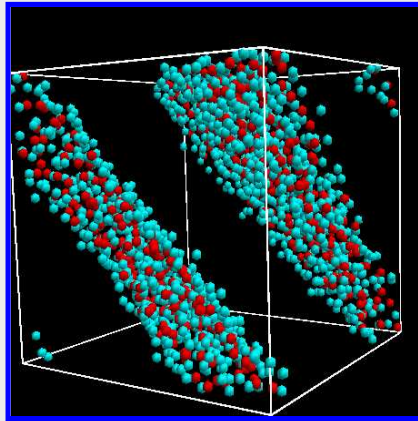
Uniform matter

Pasta phase at nonzero temperatures

Transitions by increasing temperatures

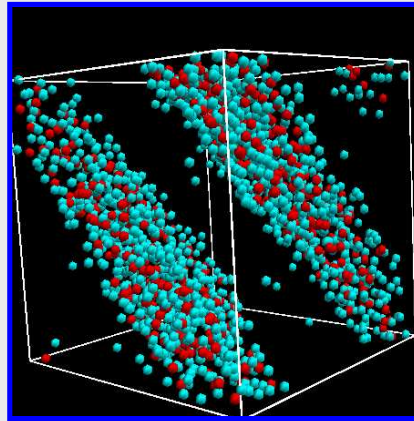
e.g. : $0.393 \rho_0$ (slablike nuclei at $T=0$ MeV) for model 2

T= 0 MeV



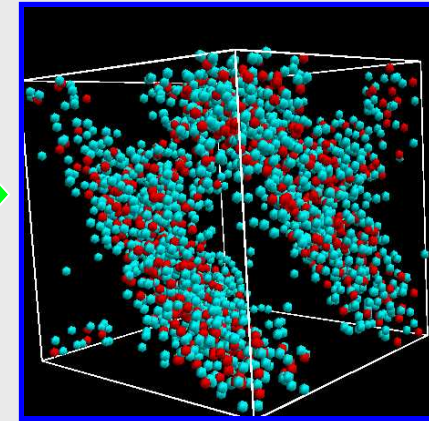
Slab-like nuclei

T= 1 MeV



Increasing dripped
neutrons

T= 2 MeV



Connected slabs

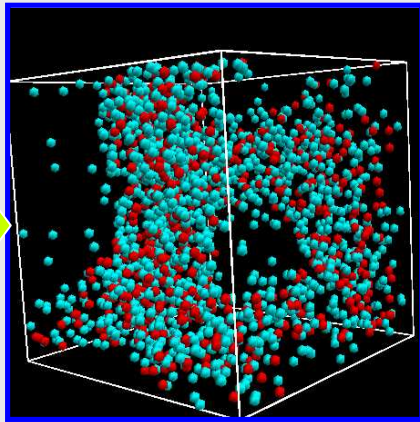


Increasing dripped neutrons
Nuclear surface diffusive

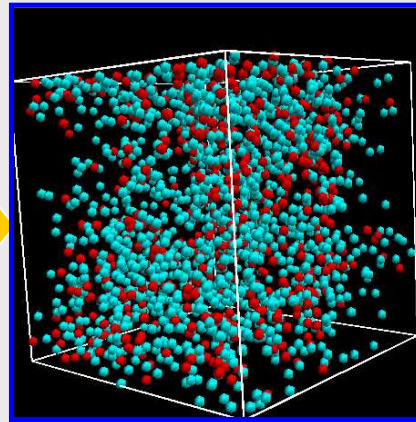
Increasing temperature

e.g: $0.393 \rho_0$ (slablike nuclei at zero temperature)

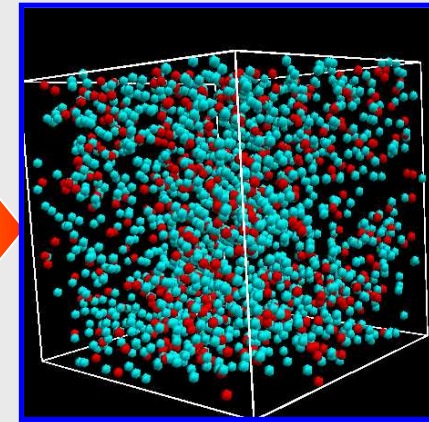
T=3MeV



T=5MeV



T=6MeV



Approaching to
rod-like bubble phase

Nuclear surface
cannot identified

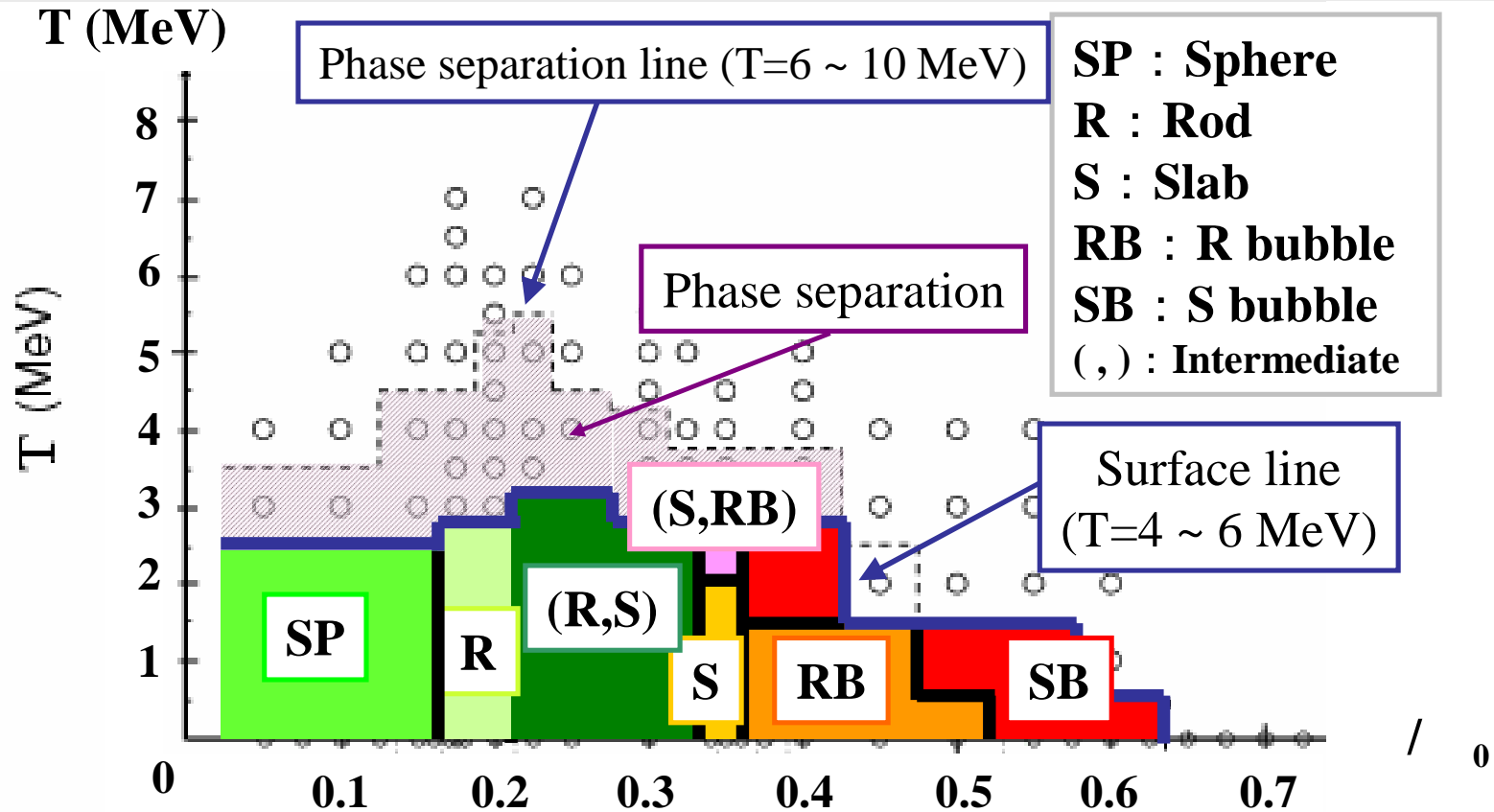
Phase separation
disappeared



**Structural transition, diffusive surface, dripped protons,
Disappearance of liquid-gas phase separation**

Phase diagram at nonzero temperature

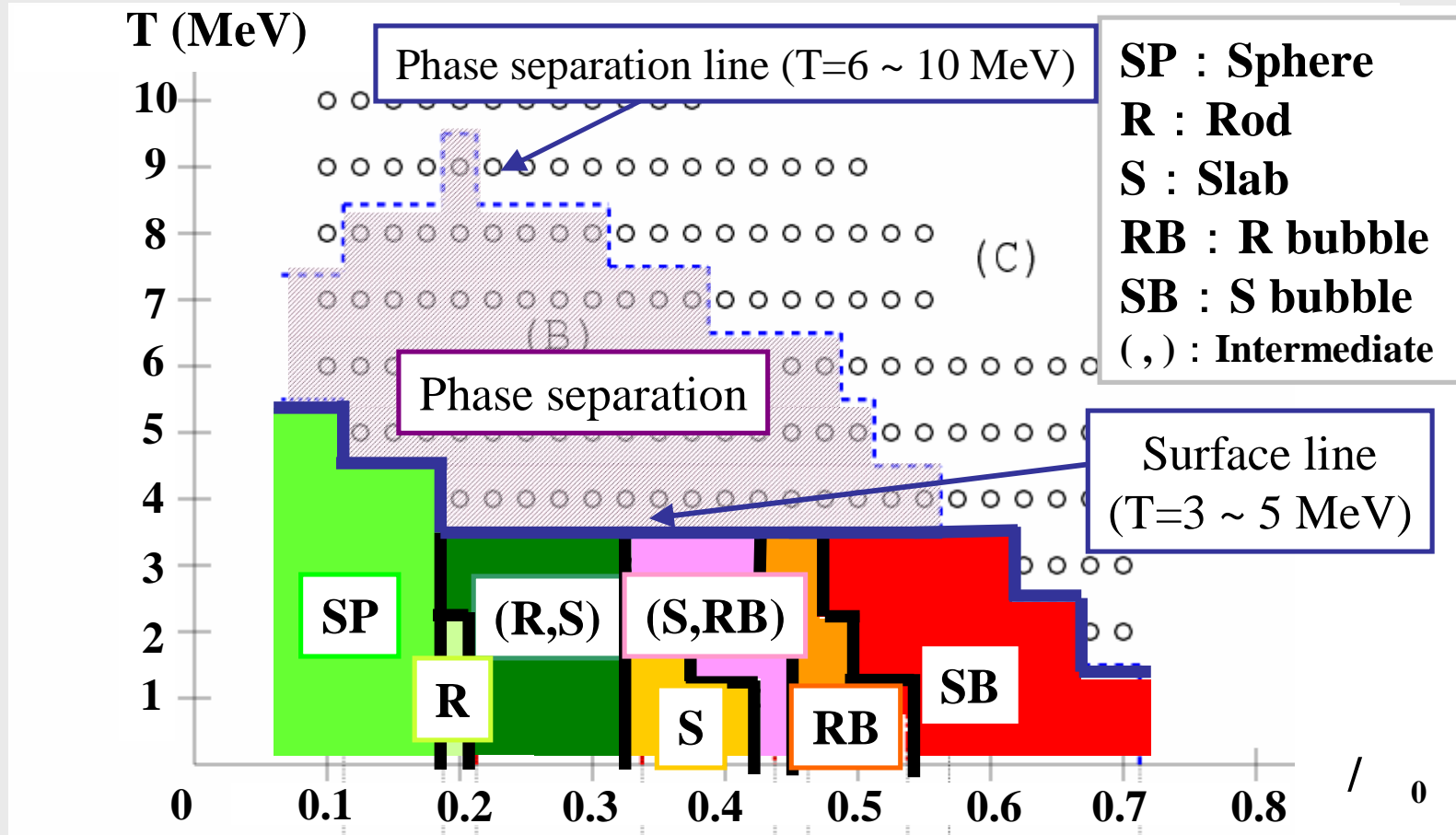
Model 1



$T=2 \sim 3$ MeV; nuclear surface cannot be identified
 $T=3 \sim 5$ MeV; uniform matter appears

Phase diagram at nonzero temperature

Model 2



$T = 3 \sim 5$ MeV; nuclear surface cannot be identified
 $T = 6 \sim 10$ MeV; uniform matter appears

Model dependence of Phase diagram

Pasta region in ρ - T plane for model 2 larger than model 1

Symmetry energy parameter L is helpful for understanding uncertainty of pasta region in neutron star
(Oyamatsu and Iida (2007))

L : density derivative of energy of pure neutron matter at normal nuclear density

L would be also helpful for pasta region in supernova

Model dependence of Phase diagram

High L = low symmetric energy at subnuclear densities

Low L = high symmetric energy at subnuclear densities

Suppose decreasing density or temperature of uniform matter, proton clustering instability occurs fast in case of low L

This extends pasta region for low L

Saturation density of asymmetric nuclear matter is high for low L

This also extends pasta region for low L

Estimated value of L is 93 MeV for model 1

80 MeV for model 2

Consistent with QMD results

(L of various nuclear forces in the range of 10~120 MeV (Oyamatsu & Iida (2007)))²¹

Summary of phase diagram

- QMD simulation of nucleon many body system at subnuclear density gives “pasta” phases
- Density dependence of symmetry energy L would be helpful for understanding phase diagram

2. Neutrino Opacity

Motivation

Neutrino transport

--- important for success of SN

Trapped when collapsing phases and effects on EOS

How does pasta phases affect neutrino opacity?

Neutrino scattering cross section

Differential scattering cross section of neutrino-nucleon systems

$$\frac{1}{N} \frac{d\sigma}{d\Omega}(\mathbf{q}) = \frac{G_F^2 E_\nu^2}{4\pi^2} (1 + \cos\theta) \cdot c_\nu^{(n)^2} \times \bar{S}(\mathbf{q})$$

Neutrino-one neutron cross section

Magnification factor
(static structure factor)

Angle averaged total transport cross section

$$\sigma_t = \langle \bar{S}(E_\nu) \rangle \sigma_t^0$$

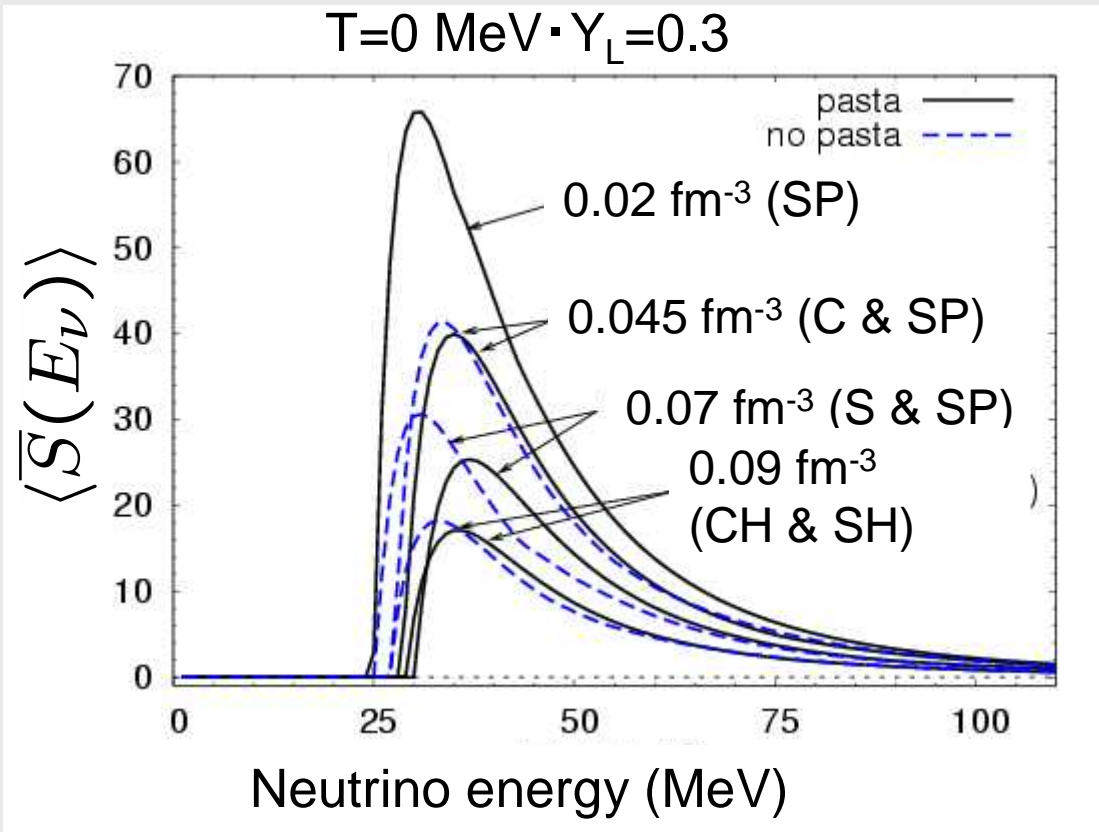
$\langle \bar{S}(E_\nu) \rangle$ magnification of total transport cross section

Results

- Understanding effects of pasta phases on neutrino opacity by liquid drop model
- Calculation of cross section by results of QMD

Scattering by pasta phases

Comparison of cases with and without pasta by BBP liquid drop model



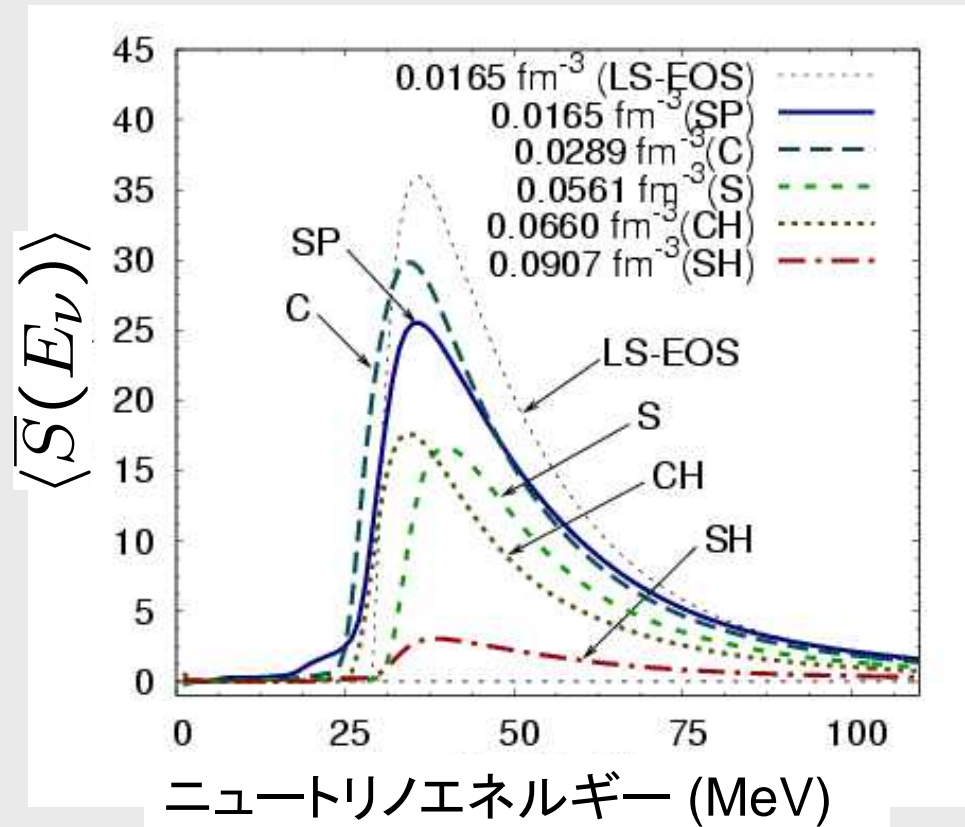
Black: Pasta phases
Blue: spherical nuclei

Peak around 30~40 MeV
Peak declines with density
Incoherent $E < 25 \text{ MeV}$

Existence of pasta phases
increases peak energy up to $E = 30\text{-}40 \text{ MeV}$
slightly increases for $E > 40 \text{ MeV}$

QMD results (Density dependence)

$Y_e=0.3, T=1 \text{ MeV}$



Compared with LD

- lower peak height
- broad peak width
- tail at low energy

Due to **thermal fluctuations**

Height reverse

SP C

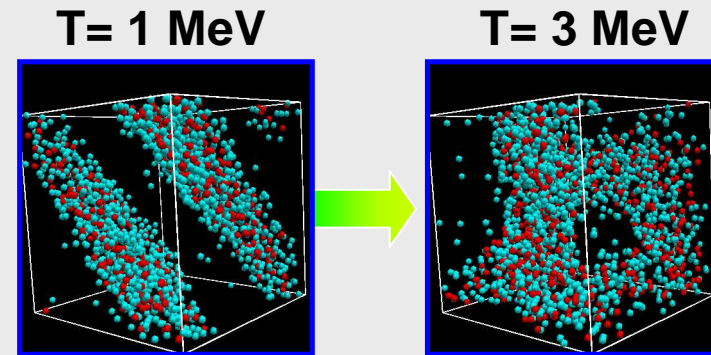
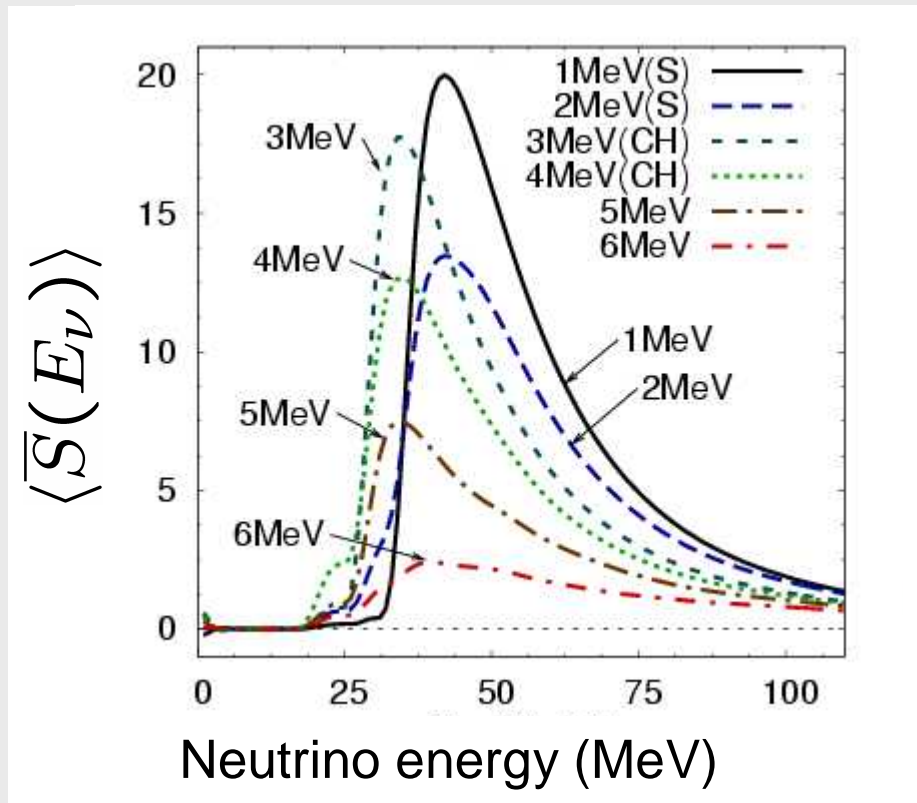
Thermal fluctuations ?

S CH

**number of reciprocal lattice
vectors ?**

QMD results (temperature dependence)

$Y_e=0.3$, $\rho=0.0660\text{fm}^{-3}$ (Slab at $T=0$)



- Peak decline with temperature
- Phase transition triggers peak energy transition reverse of peak height

Phase transitions highly complicate cross section for intermediate energy neutrinos ($E \sim 25\text{-}50$ MeV)

Summary of neutrino opacity

- Pasta phases increases peak energy (30-40 MeV)
- Increasing temperature basically lowers neutrino opacity
- Phase transitions complicate neutrino opacity for intermediate energy

Summary

- Pasta phases are reproduced in the framework of QMD
- Pasta region in ρ - T plane is understood in terms of symmetry energy parameter L ; low L enlarges pasta region
- Neutrino opacity is greatly affected by pasta phases, especially for intermediate energy neutrinos (25-50 MeV)