

Loop Quantum Gravity

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String 08
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i. Loop Quantum Gravity

- The problem addressed
- Why loops ?

ii. The theory

- Kinematics: spin networks and quanta of space
- Dynamics: spinfoams

iii. Physics

- Cosmology, black holes, singularities, low-energy limit

iv. Conclusion

- What has been achieved? What is missing?
- Future

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i. Loop Quantum Gravity

- Develops since the late 80s.
- About 200 people, 30 research groups.
- Several books (CR: "Quantum Gravity"):



The problem addressed:

How to describe
the fundamental degrees of freedom
when there is no fixed background space

Hypotheses

- A radical conceptual change in our concept of space and time is required.
- The problem can be addressed already in the context of *current* physical theory: general relativity, coupled with the standard model.
- In gravity, (unrenormalizable) UV divergences are consequences of a perturbation expansion around a wrong vacuum. → Confirmed *a posteriori* in LQG.
- Guiding principle: a symmetry: *Diffeomorphism invariance*.

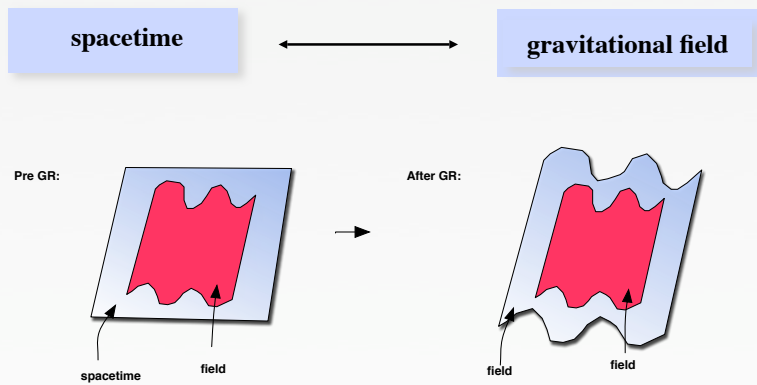
Consistency with quantum mechanics and (in the low-energy limit) General Relativity, and full diff-invariance are extremely strong constraints on the theory.

Main result

- Definition of *Diffeomorphisms invariant* quantum field theory (for gauge fields plus fermions), in canonical and in covariant form.

A comment on general relativity

- GR is: i) A *specific* field theory for the gravitational field $g_{\mu\nu}(x)$: $S[g] = \int d^4x \sqrt{-g} (R + \lambda)$
- ii) A *general* modification of our understanding of spacetime :



This modification is expressed by the invariance of the theory under the active action of the group of the *diffeomorphisms* on all the fields of the theory.

Why Loops ?

“Old” nonperturbative quantum gravity

Canonical (*Wheeler 64, DeWitt 63 ...*)

- *States:* $\Psi(\mathbf{q})$. \mathbf{q} : 3d metric of a $t=0$ surface
- *3d diff:* $\Psi(\mathbf{q}) = \Psi(\mathbf{q}')$ if there is a diff: $\mathbf{q} \rightarrow \mathbf{q}'$
- *Dynamics:* Wheeler-DeWitt eq $\mathbf{H}\Psi(\mathbf{q}) = 0$.

difficulties:

- Which states $\Psi(\mathbf{q})$? Which scalar product $\langle \Psi, \Phi \rangle$?
- Operator \mathbf{H} badly defined and UV divergent.
- No calculation possible.

Covariant (*Misner 57, Hawking 79 ...*)

- $Z = \int Dg e^{iS[g]}$

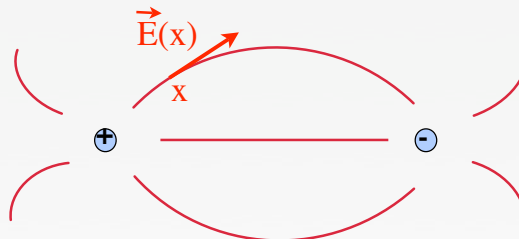
difficulties:

- Integral very badly defined
- Perturbative calculations bring UV divergences back.

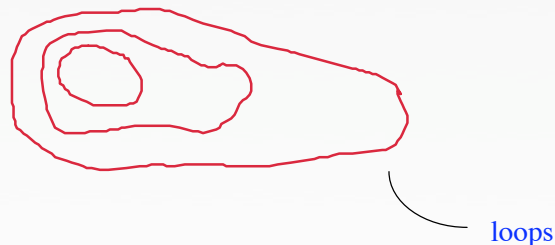
An old idea : Gauge fields are naturally “made up of lines”

(Polyakov, Mandelstam, Wilson, Migdal, ... Faraday)

Faraday lines:



In vacuum:



The gravitational field too, can be described as a gauge field, with a connection as main variable. (Cartan, Weyl, Swinger, Utyama, ..., Ashtekar)

Can we use these loop states as a basis of states in the continuum ?

No !

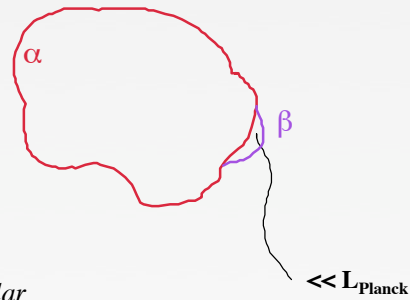
- $\langle \alpha | \alpha \rangle = \infty$

The loop states $|\alpha\rangle$ in the continuum are *too singular*

- $\langle \alpha | \beta \rangle = 0$ for β infinitely close to α

States are “too many”, as a basis of the Hilbert space.

- An *infinitesimal* displacement in space yields a *different* loop state



So far, just difficulties:

- “Old” nonperturbative quantum gravity does not work.
- A loop-state formalism in the continuum for Yang Mills does not work.

But the two ideas provide the solution to each other’s stumbling blocks:

***A loop formulation of gravity
solves both sets of difficulties***

Diffeomorphism invariance :

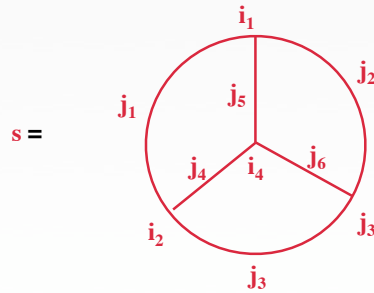
$|S\rangle$ and $|S'\rangle$ are gauge equivalent if S can be transformed into S' by a diffeomorphism !



$$\langle S | \Psi \rangle = \langle S' | \Psi \rangle = \langle S'' | \Psi \rangle \quad \text{for any } \Psi.$$

→ States are determined only by an *abstract* graph γ with j 's and i 's

s-knot states $|s\rangle = |\gamma, j, i\rangle$, where



ii. The theory.

- Start from [GR](#), or [GR+standard model](#), or any other diff-invariant theory, in a formulation where the field is described by a connection A (Ashtekar).
- Do a canonical quantization of the theory, using a basis of [spin network states](#) and operators acting on these.
- Impose [diffeomorphism invariance](#) on the states.
- Study the [Wheeler deWitt equation](#).

Result:

- A (separable) Hilbert space \mathcal{H} of states, and an operator algebra \mathcal{A} .
- Basis of \mathcal{H} : abstract spin network states: graph labelled by spins and intertwiners.
- A well defined UV-finite dynamics.

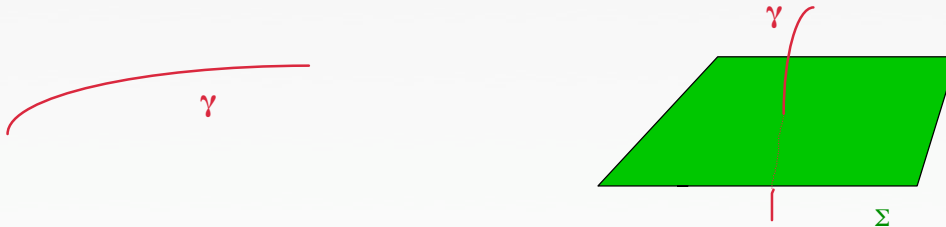
\mathcal{H} :

\mathcal{H}_{ext} : (norm-closure of the) space of the (cylindrical) functionals $\Psi_{\gamma, f}[A] = f(U_{\gamma_1}[A], \dots, U_{\gamma_n}[A])$

where $U_\gamma[A] = e^{\int_\gamma A}$, equipped with the scalar product

$$(\Psi_{\gamma, f}, \Psi_{\gamma', g}) = \int_{SU(2)^n} dU \overline{f(U_1, \dots, U_n)} g(U_1, \dots, U_n)$$

→ This product is SU(2) and diff invariant. Hence \mathcal{H}_{ext} carries a unitary representation of *Diff* and *local SU(2)*. The operation of factoring away the action of these groups is well-defined, and defines $\mathcal{H} = \mathcal{H}_{\text{ext}} / (\text{Diff and local } SU(2))$.



\mathcal{A} :

→ Operators well-defined on \mathcal{H}_{ext} , and self-adjoint: $U_\gamma[A] = e^{\int_\gamma A}$ and $E_\Sigma = \int_\Sigma \mathbf{E}$

where \mathbf{E} is the variable conjugate to A , smeared on a two-surface Σ . This acts as an SU(2) Left-invariant vector field on the SU(2) of each line that intersects Σ .

The LOST uniqueness theorem

(Fleishhack 04; Lewandowski, Okolow, Sahlmann, Thiemann 05)

- $(\mathcal{H}_{\text{ext}}, \mathcal{A})$ provides a representation of the classical poisson algebra of the observables $U_\gamma[A]$ and E_Σ , carrying a unitary representation of *Diff*,

- → this representation is *unique*.

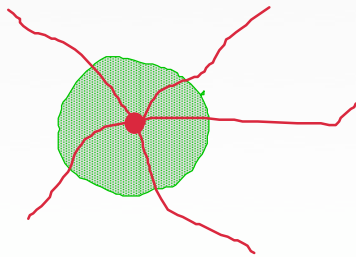
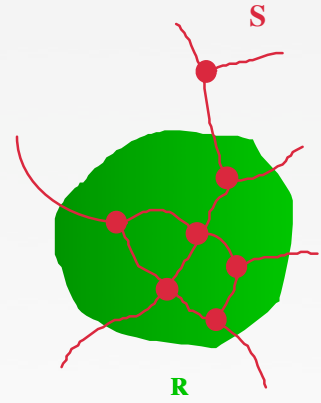
(cfr.: von Neumann theorem in nonrelativistic QM.)

Interpretation of the spin network states $|S\rangle$

Volume: $V(\mathbf{R})$: function of the gravitational field

$$V(\mathbf{R}) = \int_{\mathbf{R}} \sqrt{g} = \int_{\mathbf{R}} \sqrt{EEE} \rightarrow V(\mathbf{R}) \text{ operator}$$

- $V(\mathbf{R})$ is a well-defined self-adjoint operator in \mathcal{H}_{ext} ,
- It has **discrete** spectrum. Eigenstates: spin networks state $|S\rangle$.
- Eigenvalues receive a contribution for each **node** of $|S\rangle$ inside \mathbf{R}



Node = “**Chunk of space**”
with quantized volume

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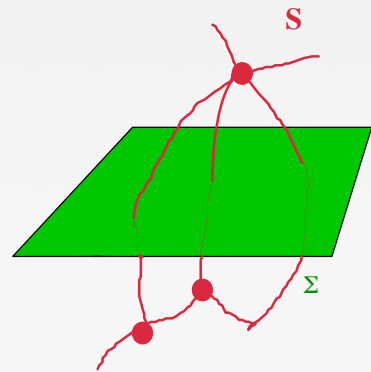
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Area: $A(\Sigma)$

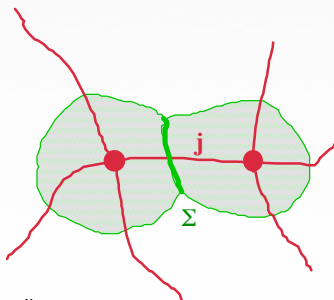
$$A(\Sigma) = \int_{\Sigma} \sqrt{g} = \int_{\Sigma} \sqrt{EE} \rightarrow A(\Sigma) \text{ operator}$$

- $A(\Sigma)$ well-defined selfadjoint operator in \mathcal{H}_{ext} .
- The spectrum is discrete.
- Area gets a contribution for each **link** of $|S\rangle$ that intersects Σ .



Area eigenvalues: $A = 8\pi \hbar G \gamma \sum_i \sqrt{j_i(j_i + 1)}$

(γ = Immirzi parameter)



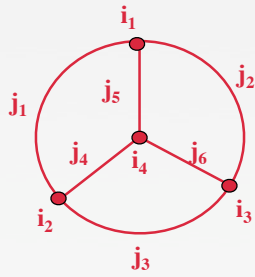
Link = “**Quantum of surface**”
with quantized area

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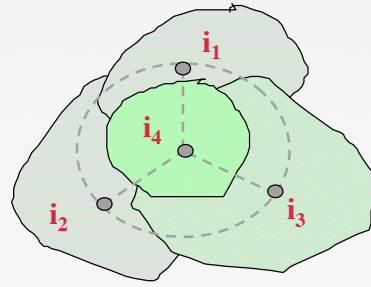
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=



$$|S\rangle = |\gamma, j_l, i_n\rangle$$

s-knot state

connectivity between the elementary quantum chunks of space

quantum numbers of volume

quantum numbers of area

→ Spin network states represent discrete quantum excitations of spacetime

- Spin networks are not excitations **in space**: they are excitations **of space**.
→ Background independent QFT
- → **Discrete structure of space at the Planck scale** - in quantum sense

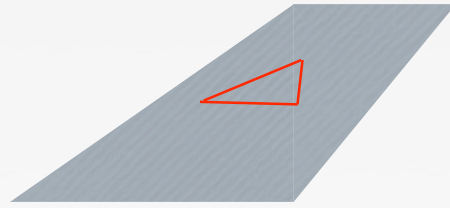
Follows from:

standard quantum theory (cfr granularity of oscillator's energy)
+ standard general relativity (because "space is a field").

Loops & strings: a cartoon comparison

If a **string** is:

a closed line **in space**, that forms matter and forces,



a **loop** is:

a closed line that **forms space itself** as well as matter and forces.



III. The theory - dynamics

$$\mathbf{H}(\mathbf{x}) \begin{array}{c} \circ \\ \mid \\ \bullet \\ \mid \\ \circ \\ \text{x} \end{array} = A \begin{array}{c} \circ \\ \mid \\ \bullet \\ \mid \\ \bullet \\ \mid \\ \bullet \\ \mid \\ \bullet \\ \mid \\ \circ \\ \text{s} \end{array} = A \begin{array}{c} \circ \\ \mid \\ \bullet \\ \mid \\ \bullet \\ \mid \\ \bullet \\ \mid \\ \bullet \\ \mid \\ \circ \\ \text{s}' \end{array}$$

Given by a Wheeler-deWitt operator \mathbf{H} in \mathcal{H} : $\mathbf{H}\Psi = 0$

- \mathbf{H} is defined by a regularization of the classical Hamiltonian constraint. In the limit in which the regularization is removed.

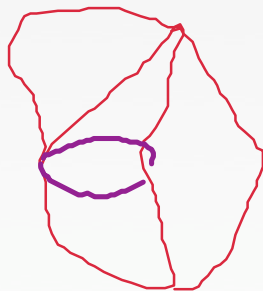
→ \mathbf{H} is a well defined self-adjoint operator, UV finite on diff-invariant states.



$\alpha \rightarrow 0$



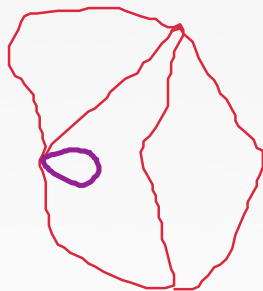
$\alpha \rightarrow 0$



$\alpha \rightarrow 0$



$\alpha \rightarrow 0$



$$\alpha \rightarrow 0$$

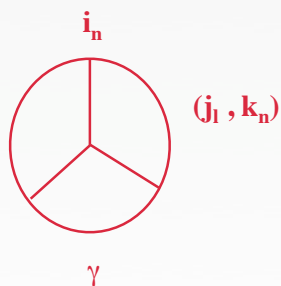


The limit $\alpha \rightarrow 0$ is trivial because there is no short distance structure at all in the theory !

- The theory is naturally ultraviolet finite

Matter

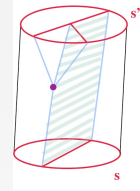
- YM, fermions
- Same techniques: The gravitational field is *not* special
- \rightarrow UV finiteness remains
- **YM and fermions on spin networks = on a Planck scale lattice !**
Notice: no lattice spacing to zero !
- Matter from braiding ?



quantum numbers
of matter

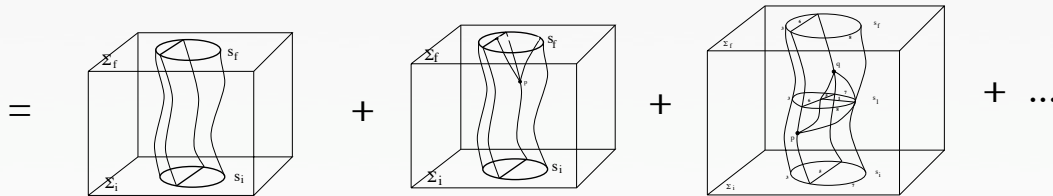
$$|S\rangle = |\gamma, j_l, i_n, k_l\rangle$$

III. The theory - covariant dynamics: spinfoams

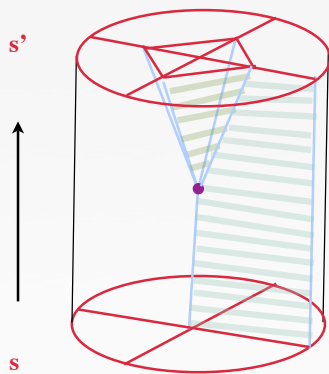
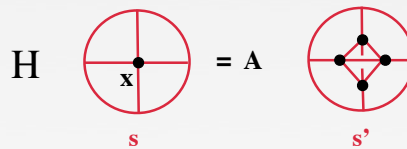


Projector $P = \delta(H)$ on the kernel of H :

$$\begin{aligned} \langle s' | P | s \rangle &= \langle s' | \delta(H) | s \rangle = \langle s' | \int DN e^{iNH} | s \rangle \\ &= c_0 \langle s' | s \rangle + c_1 \langle s' | H | s \rangle + c_2 \langle s' | H^2 | s \rangle + \dots \end{aligned}$$

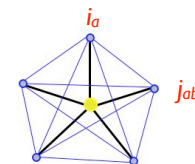


$$\langle s' | P | s \rangle = \sum_{\partial(\sigma, j_f, i_n) = (s \cup s')} \prod_f \dim(j_f) \prod_v A_v(j_f, i_n)$$



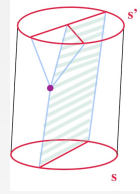
Two-complex, colored with spins and intertwiners = spinfoam

Vertex amplitude: $A(j, i)$



$$\rightarrow A(j_{ab}, i_a) = \sum_{i_a^+, i_a^-} 15j \left(\frac{(1+\gamma)j_{ab}, i_a^+}{2} \right) 15j \left(\frac{|1-\gamma|j_{ab}, i_a^-}{2} \right) \otimes_a f_{i_a^+, i_a^-}^{i_a}$$

Spinfoams



$$Z = \sum_{\sigma_{j_f i_v}} \prod_f \dim(j_f) \prod_v A_v(j_f, i_v)$$

$$A(j_{ab}, i_a) = \sum_{i_a^+, i_a^-} 15j \left(\frac{(1+\gamma)j_{ab}}{2}, i_a^+ \right) 15j \left(\frac{|1-\gamma|j_{ab}}{2}, i_a^- \right) \otimes_a f_{i_a^+, i_a^-}^i$$

- Can be directly derived from a discretization of the action of general relativity, on a *variable lattice*.
- Can be interpreted as a discrete version of Hawking's "integral over geometries"

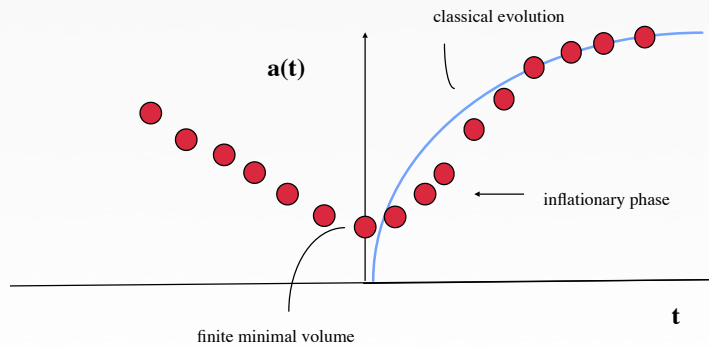
$$Z = \int Dg e^{iS[g]}$$

- → **Z** is generated as the Feynman expansion of an auxiliary field theory defined on a group manifold. This is a 4d generalization of the matrix models (à la Boulatov, Ooguri).
- → In 3d, it gives directly the old Ponzano Regge model ($A=6j$). (With cosmological constant: Turaev-Viro state sum model.)

III. Physics

Loop cosmology

- Discrete cosmological time
- → Big Bang singularity removed (from Planck scale non-locality)
- → Evolution “across” the big bang (robust)
- → (Super-) inflationary behavior at small $a(t)$
- → scale invariant spectrum with the observed spectral index n_s .



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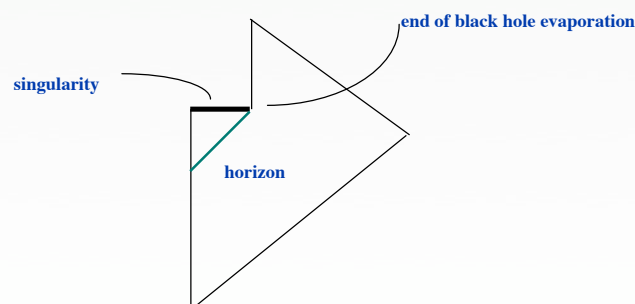
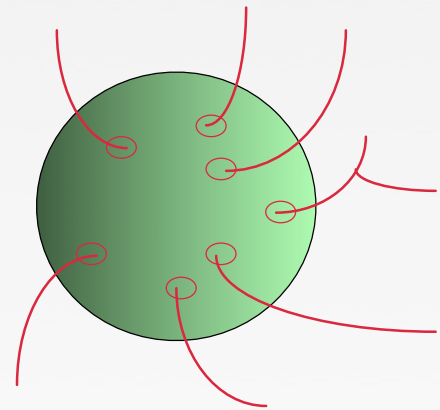
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Black holes

- → Entropy finite, proportional to the area
- → Physical black holes
- $S = A/4$ if a dimensionless free parameter (Immirzi parameter) is fixed



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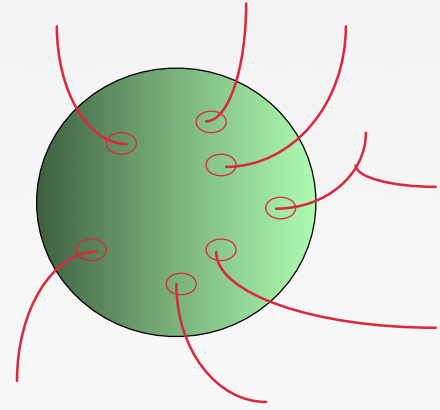
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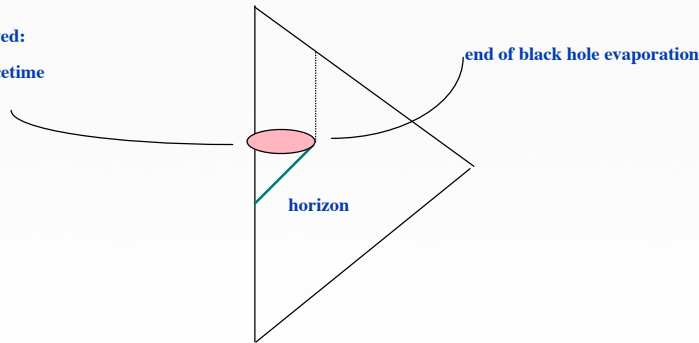
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Black holes

- → Entropy finite, proportional to the area
- → Physical black holes
- $S = A/4$ if a dimensionless free parameter (Immirzi parameter) is fixed
- → $R = 0$ singularity under control:



singularity removed:
non-classical spacetime



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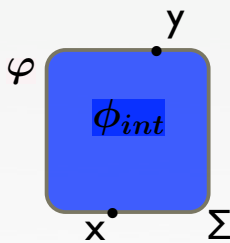
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background-independent n -point functions → low-energy limit

$$W(x, y) = \int D\phi \phi(x)\phi(y)e^{iS[\phi]} \quad \text{is independent from } x \text{ and } y \text{ if } D\phi \text{ and } S[\phi] \text{ are invariant under } Diff.$$



Choose a closed 3-surface where x and y lie, and rewrite W as

$$W(x, y) = \int D\varphi \varphi(x)\varphi(y) W[\varphi] \Psi[\varphi]$$

where

$$W[\varphi] = \int_{\phi_{int}|_{\Sigma}=\varphi} D\phi_{int} e^{iS[\phi_{int}]}$$

Instead of being determined by boundary conditions at infinity, let the boundary state be a state picked on a given boundary geometry q .

$$W[x, y; q] = \int D\varphi \varphi(x)\varphi(y) W[\varphi] \Psi_q[\varphi]$$

This expression is meaningful in a $Diff$ -invariant theory, and reduces to flat space n -point function for appropriate boundary state.

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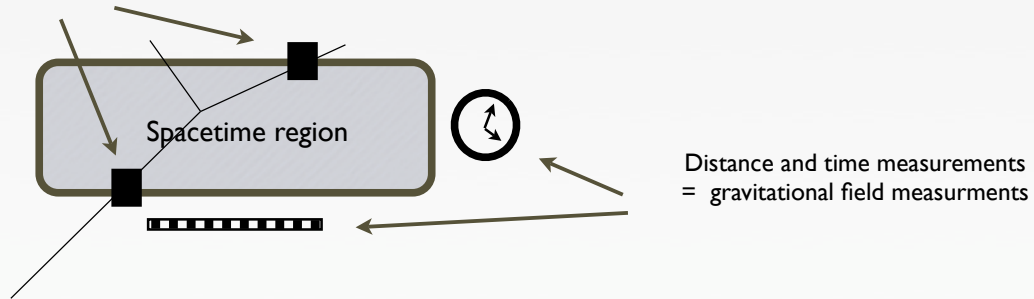
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Boundary values of the gravitational field = geometry of box surface
 = distance and time separation of measurements

Particle detectors = field measurements



In GR, distance and time measurements
 are field measurements like the other ones:
 they are part of the **boundary data** of the problem.

The graviton two-point function in LQC reads

$$\langle 0 | g^{ab}(x) g^{cd}(y) | 0 \rangle = \sum_s W[s] g^{ab}(x) g^{cd}(y) \Psi_q[s]$$

And for large distances this is give at first order by

$$W[s] = \frac{\lambda}{5!} \left(\prod_{n < m} \dim(j_{nm}) \right) A_{vertex}(j_{nm})$$

→ The asymptotic expansion of the vertex gives the low-energy behavior of the theory.

→ Preliminary results yield: - free graviton propagator; - 3 point function; (hence Newton law); - first order corrections to the free graviton propagator. Calculations in progress.

Mathematical developments

- Diffeomorphism invariant measures (Ashtekar-Lewandowski measure)
- C* algebraic techniques
- Category theory
- “Quantum geometry”
- Uniqueness of the representation

IV. Summary

- **Loop quantum gravity** is a technique for defining **Diff-invariant QFT**. It offers a radically new description of space and time by merging in depth QFT with the *diff*-invariance introduced by GR.
- It provides a quantum theory of GR plus the standard model in 4d, which is naturally **UV finite** and has a **discrete structure of space** at Planck scale.
- Has **applications** in cosmology, black hole physics, astrophysics; it resolves black hole and big bang singularities.
 - Unrelated to a natural unification of the forces (we are *not* at the “end of physics”).
 - Different versions of the dynamics exist.
 - Low-energy limit still in progress.
- + Fundamental degrees of freedom explicit.
- + The theory is consistent with today’s physics.
- + No *need* of higher dimensions (high-d formulation possible).
- + No *need* of supersymmetry (supersymmetric theories possible).
- + Consistent with, and based on, basic QM and GR insights.