## Advanced Differential Geometry for Theoreticians

Fiber Bundles, Jet Manifolds and Lagrangian Theory


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## Contents

Introduction ..... 5
1 Geometry of fibre bundles ..... 7
1.1 Fibre bundles ..... 7
1.2 Vector and affine bundles ..... 15
1.3 Vector fields ..... 22
1.4 Exterior and tangent-valued forms ..... 25
2 Jet manifolds ..... 33
2.1 First order jet manifolds ..... 33
2.2 Higher order jet manifolds ..... 35
2.3 Differential operators and equations ..... 42
2.4 Infinite order jet formalism ..... 45
3 Connections on fibre bundles ..... 51
3.1 Connections as tangent-valued forms ..... 51
3.2 Connections as jet bundle sections ..... 54
3.3 Curvature and torsion ..... 57
3.4 Linear and affine connections ..... 58
3.5 Flat connections ..... 63
3.6 Connections on composite bundles ..... 64
4 Geometry of principal bundles ..... 69
4.1 Geometry of Lie groups ..... 69
4.2 Bundles with structure groups ..... 73
4.3 Principal bundles ..... 76
4.4 Principal connections ..... 81
4.5 Canonical principal connection ..... 86
4.6 Gauge transformations ..... 88
4.7 Geometry of associated bundles ..... 91
4.8 Reduced structure ..... 95
5 Geometry of natural bundles ..... 99
5.1 Natural bundles ..... 99
5.2 Linear world connections ..... 103
5.3 Affine world connections ..... 108
6 Geometry of graded manifolds ..... 113
6.1 Grassmann-graded algebraic calculus ..... 113
6.2 Grassmann-graded differential calculus ..... 117
6.3 Graded manifolds ..... 122
6.4 Graded differential forms ..... 128
7 Lagrangian theory ..... 131
7.1 Variational bicomplex ..... 131
7.2 Lagrangian theory on fibre bundles ..... 133
7.3 Grassmann-graded Lagrangian theory ..... 142
7.4 Noether identities ..... 152
7.5 Gauge symmetries ..... 159
8 Topics on commutative geometry ..... 163
8.1 Commutative algebra ..... 163
8.2 Differential operators on modules ..... 169
8.3 Homology and cohomology of complexes ..... 172
8.4 Differential calculus over a commutative ring ..... 175
8.5 Sheaf cohomology ..... 179
8.6 Local-ringed spaces ..... 188
Bibliography ..... 195
Index ..... 198

## Introduction

In contrast with quantum field theory, classical field theory can be formulated in a strict mathematical way by treating classical fields as sections of smooth fibre bundles $[9,17,21,24]$. This also is the case of timedependent non-relativistic mechanics on fibre bundles over $\mathbb{R}[10,16,19]$.

This book aim to compile the relevant material on fibre bundles, jet manifolds, connections, graded manifolds and Lagrangian theory [9, 17, $22]$.

The book is based on the graduate and post graduate courses of lectures given at the Department of Theoretical Physics of Moscow State University (Russia). It addresses to a wide audience of mathematicians, mathematical physicists and theoreticians. It is tacitly assumed that the reader has some familiarity with the basics of differential geometry [11, 13, 26].

Throughout the book, all morphisms are smooth (i.e. of class $C^{\infty}$ ) and manifolds are smooth real and finite-dimensional. A smooth real manifold is customarily assumed to be Hausdorff and second-countable (i.e., it has a countable base for topology). Consequently, it is a locally compact space which is a union of a countable number of compact subsets, a separable space (i.e., it has a countable dense subset), a paracompact and completely regular space. Being paracompact, a smooth manifold admits a partition of unity by smooth real functions. Unless otherwise stated, manifolds are assumed to be connected (and, consequently, arcwise connected). We follow the notion of a manifold without boundary.

## Chapter 1

## Geometry of fibre bundles

Throughout the book, fibre bundles are smooth finite-dimensional and locally-trivial.

### 1.1 Fibre bundles

Let $Z$ be a manifold. By

$$
\pi_{Z}: T Z \rightarrow Z, \quad \pi_{Z}^{*}: T^{*} Z \rightarrow Z
$$

are denoted its tangent and cotangent bundles, respectively. Given coordinates ( $z^{\alpha}$ ) on $Z$, they are equipped with the holonomic coordinates

$$
\begin{array}{ll}
\left(z^{\lambda}, \dot{z}^{\lambda}\right), & \dot{z}^{\prime \lambda}=\frac{\partial z^{\prime \lambda}}{\partial z^{\mu}} \dot{z}^{\mu}, \\
\left(z^{\lambda}, \dot{z}_{\lambda}\right), & \dot{z}_{\lambda}^{\prime}=\frac{\partial z^{\prime \mu}}{\partial z^{\lambda}} \dot{\mu}_{\mu},
\end{array}
$$

with respect to the holonomic frames $\left\{\partial_{\lambda}\right\}$ and coframes $\left\{d z^{\lambda}\right\}$ in the tangent and cotangent spaces to $Z$, respectively. Any manifold morphism $f: Z \rightarrow Z^{\prime}$ yields the tangent morphism

$$
T f: T Z \rightarrow T Z^{\prime}, \quad \dot{z}^{\prime \lambda} \circ T f=\frac{\partial f^{\lambda}}{\partial x^{\mu}} \dot{z}^{\mu} .
$$

Let us consider manifold morphisms of maximal rank. They are immersions (in particular, imbeddings) and submersions. An injective immersion is a submanifold, and a surjective submersion is a fibred manifold (in particular, a fibre bundle).

Given manifolds $M$ and $N$, by the rank of a morphism $f: M \rightarrow N$ at a point $p \in M$ is meant the rank of the linear morphism

$$
T_{p} f: T_{p} M \rightarrow T_{f(p)} N
$$

For instance, if $f$ is of maximal rank at $p \in M$, then $T_{p} f$ is injective when $\operatorname{dim} M \leq \operatorname{dim} N$ and surjective when $\operatorname{dim} N \leq \operatorname{dim} M$. In this case, $f$ is called an immersion and a submersion at a point $p \in M$, respectively.

Since $p \rightarrow \operatorname{rank}_{p} f$ is a lower semicontinuous function, then the morphism $T_{p} f$ is of maximal rank on an open neighbourhood of $p$, too. It follows from the inverse function theorem that:

- if $f$ is an immersion at $p$, then it is locally injective around $p$.
- if $f$ is a submersion at $p$, it is locally surjective around $p$.

If $f$ is both an immersion and a submersion, it is called a local diffeomorphism at $p$. In this case, there exists an open neighbourhood $U$ of $p$ such that $f: U \rightarrow f(U)$ is a diffeomorphism onto an open set $f(U) \subset N$.

A manifold morphism $f$ is called the immersion (resp. submersion) if it is an immersion (resp. submersion) at all points of $M$. A submersion is necessarily an open map, i.e., it sends open subsets of $M$ onto open subsets of $N$. If an immersion $f$ is open (i.e., $f$ is a homeomorphism onto $f(M)$ equipped with the relative topology from $N)$, it is called the imbedding.

A pair $(M, f)$ is called a submanifold of $N$ if $f$ is an injective immersion. A submanifold $(M, f)$ is an imbedded submanifold if $f$ is an imbedding. For the sake of simplicity, we usually identify $(M, f)$ with $f(M)$. If $M \subset N$, its natural injection is denoted by $i_{M}: M \rightarrow N$.

There are the following criteria for a submanifold to be imbedded.
Theorem 1.1.1: Let $(M, f)$ be a submanifold of $N$.
(i) The map $f$ is an imbedding iff, for each point $p \in M$, there exists a (cubic) coordinate chart $(V, \psi)$ of $N$ centered at $f(p)$ so that $f(M) \cap V$ consists of all points of $V$ with coordinates $\left(x^{1}, \ldots, x^{m}, 0, \ldots, 0\right)$.
(ii) Suppose that $f: M \rightarrow N$ is a proper map, i.e., the pre-images of compact sets are compact. Then $(M, f)$ is a closed imbedded submanifold of $N$. In particular, this occurs if $M$ is a compact manifold.
(iii) If $\operatorname{dim} M=\operatorname{dim} N$, then $(M, f)$ is an open imbedded submanifold of $N$.

A triple

$$
\begin{equation*}
\pi: Y \rightarrow X, \quad \operatorname{dim} X=n>0 \tag{1.1.1}
\end{equation*}
$$

is called a fibred manifold if a manifold morphism $\pi$ is a surjective submersion, i.e., the tangent morphism $T \pi: T Y \rightarrow T X$ is a surjection. One says that $Y$ is a total space of a fibred manifold (1.1.1), $X$ is its base, $\pi$ is a fibration, and $Y_{x}=\pi^{-1}(x)$ is a fibre over $x \in X$.

Any fibre is an imbedded submanifold of $Y$ of dimension $\operatorname{dim} Y-$ $\operatorname{dim} X$.

Theorem 1.1.2: A surjection (1.1.1) is a fired manifold iff a manifold $Y$ admits an atlas of coordinate charts $\left(U_{Y} ; x^{\lambda}, y^{i}\right)$ such that $\left(x^{\lambda}\right)$ are coordinates on $\pi\left(U_{Y}\right) \subset X$ and coordinate transition functions read

$$
x^{\prime \lambda}=f^{\lambda}\left(x^{\mu}\right), \quad y^{\prime i}=f^{i}\left(x^{\mu}, y^{j}\right)
$$

These coordinates are called fibred coordinates compatible with a fibration $\pi$.

By a local section of a surjection (1.1.1) is meant an injection $s: U \rightarrow$ $Y$ of an open subset $U \subset X$ such that $\pi \circ s=\operatorname{Id} U$, i.e., a section sends any point $x \in X$ into the fibre $Y_{x}$ over this point. A local section also is defined over any subset $N \in X$ as the restriction to $N$ of a local section over an open set containing $N$. If $U=X$, one calls $s$ the global section. Hereafter, by a section is meant both a global section and a local section (over an open subset).

Theorem 1.1.3: A surjection $\pi$ (1.1.1) is a fibred manifold iff, for each point $y \in Y$, there exists a local section $s$ of $\pi: Y \rightarrow X$ passing through $y$.

The range $s(U)$ of a local section $s: U \rightarrow Y$ of a fibred manifold $Y \rightarrow X$ is an imbedded submanifold of $Y$. It also is a closed map, which sends closed subsets of $U$ onto closed subsets of $Y$. If $s$ is a global section, then $s(X)$ is a closed imbedded submanifold of $Y$. Global sections of a fibred manifold need not exist.

Theorem 1.1.4: Let $Y \rightarrow X$ be a fibred manifold whose fibres are diffeomorphic to $\mathbb{R}^{m}$. Any its section over a closed imbedded submanifold (e.g., a point) of $X$ is extended to a global section. In particular, such a fibred manifold always has a global section.

Given fibred coordinates $\left(U_{Y} ; x^{\lambda}, y^{i}\right)$, a section $s$ of a fibred manifold $Y \rightarrow X$ is represented by collections of local functions $\left\{s^{i}=y^{i} \circ s\right\}$ on $\pi\left(U_{Y}\right)$.

A fibred manifold $Y \rightarrow X$ is called a fibre bundle if admits a fibred coordinate atlas $\left\{\left(\pi^{-1}\left(U_{\xi}\right) ; x^{\lambda}, y^{i}\right)\right\}$ over a cover $\left\{\pi^{-1}\left(U_{\iota}\right)\right\}$ of $Y$ which is the inverse image of a cover $\mathfrak{U}=\left\{U_{\xi}\right\}$ is a cover of $X$. In this case, there exists a manifold $V$, called a typical fibre, such that $Y$ is locally diffeomorphic to the splittings

$$
\begin{equation*}
\psi_{\xi}: \pi^{-1}\left(U_{\xi}\right) \rightarrow U_{\xi} \times V, \tag{1.1.2}
\end{equation*}
$$

glued together by means of transition functions

$$
\begin{equation*}
\varrho_{\xi \zeta}=\psi_{\xi} \circ \psi_{\zeta}^{-1}: U_{\xi} \cap U_{\zeta} \times V \rightarrow U_{\xi} \cap U_{\zeta} \times V \tag{1.1.3}
\end{equation*}
$$

on overlaps $U_{\xi} \cap U_{\zeta}$. Transition functions $\varrho_{\xi \zeta}$ fulfil the cocycle condition

$$
\begin{equation*}
\varrho_{\xi \zeta} \circ \varrho_{\zeta \iota}=\varrho_{\xi \iota} \tag{1.1.4}
\end{equation*}
$$

on all overlaps $U_{\xi} \cap U_{\zeta} \cap U_{\iota}$. Restricted to a point $x \in X$, trivialization morphisms $\psi_{\xi}$ (1.1.2) and transition functions $\varrho_{\xi \zeta}$ (1.1.3) define diffeomorphisms of fibres

$$
\begin{array}{ll}
\psi_{\xi}(x): Y_{x} \rightarrow V, & x \in U_{\xi} \\
\varrho_{\xi \zeta}(x): V \rightarrow V, & x \in U_{\xi} \cap U_{\zeta} \tag{1.1.6}
\end{array}
$$

Trivialization charts $\left(U_{\xi}, \psi_{\xi}\right)$ together with transition functions $\varrho_{\xi \zeta}$ (1.1.3) constitute a bundle atlas

$$
\begin{equation*}
\Psi=\left\{\left(U_{\xi}, \psi_{\xi}\right), \varrho_{\xi \zeta}\right\} \tag{1.1.7}
\end{equation*}
$$

of a fibre bundle $Y \rightarrow X$. Two bundle atlases are said to be equivalent if their union also is a bundle atlas, i.e., there exist transition functions between trivialization charts of different atlases.

A fibre bundle $Y \rightarrow X$ is uniquely defined by a bundle atlas. Given an atlas $\Psi$ (1.1.7), there is a unique manifold structure on $Y$ for which $\pi: Y \rightarrow X$ is a fibre bundle with the typical fibre $V$ and the bundle atlas $\Psi$. All atlases of a fibre bundle are equivalent.

Remark 1.1.1: The notion of a fibre bundle introduced above is the notion of a smooth locally trivial fibre bundle. In a general setting, a continuous fibre bundle is defined as a continuous surjective submersion of topological spaces $Y \rightarrow X$. A continuous map $\pi: Y \rightarrow X$ is called a submersion if, for any point $y \in Y$, there exists an open neighborhood $U$ of the point $\pi(y)$ and a right inverse $\sigma: U \rightarrow Y$ of $\pi$ such that $\sigma \circ \pi(y)=y$, i.e., there exists a local section of $\pi$. The notion of a locally trivial continuous fibre bundle is a repetition of that of a smooth fibre bundle, where trivialization morphisms $\psi_{\xi}$ and transition functions $\varrho_{\xi \zeta}$ are continuous.

We have the following useful criteria for a fibred manifold to be a fibre bundle.

Theorem 1.1.5: If a fibration $\pi: Y \rightarrow X$ is a proper map, then $Y \rightarrow X$ is a fibre bundle. In particular, a fibred manifold with a compact total space is a fibre bundle.

Theorem 1.1.6: A fibred manifold whose fibres are diffeomorphic either to a compact manifold or $\mathbb{R}^{r}$ is a fibre bundle.

A comprehensive relation between fibred manifolds and fibre bundles
is given in Remark 3.1.2. It involves the notion of an Ehresmann connection.

Unless otherwise stated, we restrict our consideration to fibre bundles. Without a loss of generality, we further assume that a cover $\mathfrak{U}$ for a bundle atlas of $Y \rightarrow X$ also is a cover for a manifold atlas of the base $X$. Then, given a bundle atlas $\Psi$ (1.1.7), a fibre bundle $Y$ is provided with the associated bundle coordinates

$$
x^{\lambda}(y)=\left(x^{\lambda} \circ \pi\right)(y), \quad y^{i}(y)=\left(y^{i} \circ \psi_{\xi}\right)(y), \quad y \in \pi^{-1}\left(U_{\xi}\right)
$$

where $x^{\lambda}$ are coordinates on $U_{\xi} \subset X$ and $y^{i}$, called fibre coordinates, are coordinates on a typical fibre $V$.

The forthcoming Theorems 1.1.7-1.1.9 describe the particular covers which one can choose for a bundle atlas. Throughout the book, only proper covers of manifolds are considered, i.e., $U_{\xi} \neq U_{\zeta}$ if $\zeta \neq \xi$. Recall that a cover $\mathfrak{U}^{\prime}$ is a refinement of a cover $\mathfrak{U}$ if, for each $U^{\prime} \in \mathfrak{U}^{\prime}$, there exists $U \in \mathfrak{U}$ such that $U^{\prime} \subset U$. If a fibre bundle $Y \rightarrow X$ has a bundle atlas over a cover $\mathfrak{U}$ of $X$, it admits a bundle atlas over any refinement of $\mathfrak{U}$.

A fibred manifold $Y \rightarrow X$ is called trivial if $Y$ is diffeomorphic to the product $X \times V$. Different trivializations of $Y \rightarrow X$ differ from each other in surjections $Y \rightarrow V$.

Theorem 1.1.7: Any fibre bundle over a contractible base is trivial.
However, a fibred manifold over a contractible base need not be trivial, even its fibres are mutually diffeomorphic.

It follows from Theorem 1.1.7 that any cover of a base $X$ consisting of domains (i.e., contractible open subsets) is a bundle cover.

Theorem 1.1.8: Every fibre bundle $Y \rightarrow X$ admits a bundle atlas over a countable cover $\mathfrak{U}$ of $X$ where each member $U_{\xi}$ of $\mathfrak{U}$ is a domain whose closure $\bar{U}_{\xi}$ is compact.

If a base $X$ is compact, there is a bundle atlas of $Y$ over a finite cover of $X$ which obeys the condition of Theorem 1.1.8.

Theorem 1.1.9: Every fibre bundle $Y \rightarrow X$ admits a bundle atlas over a finite cover $\mathfrak{U}$ of $X$, but its members need not be contractible and connected.

Morphisms of fibre bundles, by definition, are fibrewise morphisms, sending a fibre to a fibre. Namely, a bundle morphism of a fibre bundle $\pi: Y \rightarrow X$ to a fibre bundle $\pi^{\prime}: Y^{\prime} \rightarrow X^{\prime}$ is defined as a pair $(\Phi, f)$ of manifold morphisms which form a commutative diagram

$$
\begin{aligned}
& Y \xrightarrow{\Phi} Y^{\prime} \\
& \pi \\
& X \xrightarrow{f}{ }_{\square} X^{\prime}
\end{aligned}
$$

Bundle injections and surjections are called bundle monomorphisms and epimorphisms, respectively. A bundle diffeomorphism is called a bundle isomorphism, or a bundle automorphism if it is an isomorphism to itself. For the sake of brevity, a bundle morphism over $f=\operatorname{Id} X$ is often said to be a bundle morphism over $X$, and is denoted by $Y \underset{X}{\longrightarrow} Y^{\prime}$. In particular, a bundle automorphism over $X$ is called a vertical automorphism.

A bundle monomorphism $\Phi: Y \rightarrow Y^{\prime}$ over $X$ is called a subbundle of a fibre bundle $Y^{\prime} \rightarrow X$ if $\Phi(Y)$ is a submanifold of $Y^{\prime}$. There is the following useful criterion for an image and an inverse image of a bundle morphism to be subbundles.

Theorem 1.1.10: Let $\Phi: Y \rightarrow Y^{\prime}$ be a bundle morphism over $X$. Given a global section $s$ of the fibre bundle $Y^{\prime} \rightarrow X$ such that $s(X) \subset \Phi(Y)$, by the kernel of a bundle morphism $\Phi$ with respect to a section $s$ is meant the inverse image

$$
\operatorname{Ker}_{s} \Phi=\Phi^{-1}(s(X))
$$

of $s(X)$ by $\Phi$. If $\Phi: Y \rightarrow Y^{\prime}$ is a bundle morphism of constant rank over $X$, then $\Phi(Y)$ and $\operatorname{Ker}_{s} \Phi$ are subbundles of $Y^{\prime}$ and $Y$, respectively.

Let us describe the following standard constructions of new fibre bundles from the old ones.

- Given a fibre bundle $\pi: Y \rightarrow X$ and a manifold morphism $f: X^{\prime} \rightarrow$ $X$, the pull-back of $Y$ by $f$ is called the manifold

$$
\begin{equation*}
f^{*} Y=\left\{\left(x^{\prime}, y\right) \in X^{\prime} \times Y: \pi(y)=f\left(x^{\prime}\right)\right\} \tag{1.1.8}
\end{equation*}
$$

together with the natural projection $\left(x^{\prime}, y\right) \rightarrow x^{\prime}$. It is a fibre bundle over $X^{\prime}$ such that the fibre of $f^{*} Y$ over a point $x^{\prime} \in X^{\prime}$ is that of $Y$ over the point $f\left(x^{\prime}\right) \in X$. There is the canonical bundle morphism

$$
\begin{equation*}
f_{Y}:\left.f^{*} Y \ni\left(x^{\prime}, y\right)\right|_{\pi(y)=f\left(x^{\prime}\right)} \vec{f} y \in Y \tag{1.1.9}
\end{equation*}
$$

Any section $s$ of a fibre bundle $Y \rightarrow X$ yields the pull-back section

$$
f^{*} s\left(x^{\prime}\right)=\left(x^{\prime}, s\left(f\left(x^{\prime}\right)\right)\right.
$$

of $f^{*} Y \rightarrow X^{\prime}$.

- If $X^{\prime} \subset X$ is a submanifold of $X$ and $i_{X^{\prime}}$ is the corresponding natural injection, then the pull-back bundle

$$
i_{X^{\prime}}^{*} Y=\left.Y\right|_{X^{\prime}}
$$

is called the restriction of a fibre bundle $Y$ to the submanifold $X^{\prime} \subset X$. If $X^{\prime}$ is an imbedded submanifold, any section of the pull-back bundle

$$
\left.Y\right|_{X^{\prime}} \rightarrow X^{\prime}
$$

is the restriction to $X^{\prime}$ of some section of $Y \rightarrow X$.

- Let $\pi: Y \rightarrow X$ and $\pi^{\prime}: Y^{\prime} \rightarrow X$ be fibre bundles over the same base $X$. Their bundle product $Y \times_{X} Y^{\prime}$ over $X$ is defined as the pull-back

$$
\underset{X}{\times} Y^{\prime}=\pi^{*} Y^{\prime} \quad \text { or } \quad Y \underset{X}{\times} Y^{\prime}=\pi^{\prime *} Y
$$

together with its natural surjection onto $X$. Fibres of the bundle product $Y \times Y^{\prime}$ are the Cartesian products $Y_{x} \times Y_{x}^{\prime}$ of fibres of fibre bundles $Y$ and $Y^{\prime}$.

Let us consider the composition

$$
\begin{equation*}
\pi: Y \rightarrow \Sigma \rightarrow X \tag{1.1.10}
\end{equation*}
$$

of fibre bundles

$$
\begin{align*}
& \pi_{Y \Sigma}: Y \rightarrow \Sigma  \tag{1.1.11}\\
& \pi_{\Sigma X}: \Sigma \rightarrow X \tag{1.1.12}
\end{align*}
$$

One can show that it is a fibre bundle, called the composite bundle. It is provided with bundle coordinates $\left(x^{\lambda}, \sigma^{m}, y^{i}\right)$, where $\left(x^{\lambda}, \sigma^{m}\right)$ are bundle coordinates on the fibre bundle (1.1.12), i.e., transition functions of coordinates $\sigma^{m}$ are independent of coordinates $y^{i}$.

Theorem 1.1.11: Given a composite bundle (1.1.10), let $h$ be a global section of the fibre bundle $\Sigma \rightarrow X$. Then the restriction

$$
\begin{equation*}
Y^{h}=h^{*} Y \tag{1.1.13}
\end{equation*}
$$

of the fibre bundle $Y \rightarrow \Sigma$ to $h(X) \subset \Sigma$ is a subbundle of the fibre bundle $Y \rightarrow X$.

Theorem 1.1.12: Given a section $h$ of the fibre bundle $\Sigma \rightarrow X$ and a section $s_{\Sigma}$ of the fibre bundle $Y \rightarrow \Sigma$, their composition $s=s_{\Sigma} \circ h$ is a section of the composite bundle $Y \rightarrow X$ (1.1.10). Conversely, every section $s$ of the fibre bundle $Y \rightarrow X$ is a composition of the section $h=\pi_{Y \Sigma} \circ s$ of the fibre bundle $\Sigma \rightarrow X$ and some section $s_{\Sigma}$ of the fibre bundle $Y \rightarrow \Sigma$ over the closed imbedded submanifold $h(X) \subset \Sigma$.

### 1.2 Vector and affine bundles

A vector bundle is a fibre bundle $Y \rightarrow X$ such that:

- its typical fibre $V$ and all the fibres $Y_{x}=\pi^{-1}(x), x \in X$, are real finite-dimensional vector spaces;
- there is a bundle atlas $\Psi(1.1 .7)$ of $Y \rightarrow X$ whose trivialization morphisms $\psi_{\xi}$ (1.1.5) and transition functions $\varrho_{\xi \zeta}$ (1.1.6) are linear isomorphisms.
Accordingly, a vector bundle is provided with linear bundle coordinates $\left(y^{i}\right)$ possessing linear transition functions $y^{i}=A_{j}^{i}(x) y^{j}$. We have

$$
\begin{equation*}
y=y^{i} e_{i}(\pi(y))=y^{i} \psi_{\xi}(\pi(y))^{-1}\left(e_{i}\right), \quad \pi(y) \in U_{\xi} \tag{1.2.1}
\end{equation*}
$$

where $\left\{e_{i}\right\}$ is a fixed basis for the typical fibre $V$ of $Y$, and $\left\{e_{i}(x)\right\}$ are the fibre bases (or the frames) for the fibres $Y_{x}$ of $Y$ associated to the bundle atlas $\Psi$.

By virtue of Theorem 1.1.4, any vector bundle has a global section, e.g., the canonical global zero-valued section $\widehat{0}(x)=0$. Global sections of a vector bundle $Y \rightarrow X$ constitute a projective $C^{\infty}(X)$-module $Y(X)$ of finite rank. It is called the structure module of a vector bundle.

Theorem 1.2.1: Let a vector bundle $Y \rightarrow X$ admit $m=\operatorname{dim} V$ nowhere vanishing global sections $s_{i}$ which are linearly independent, i.e., $\stackrel{m}{\wedge} s_{i} \neq 0$. Then $Y$ is trivial.

Remark 1.2.1: Theorem 8.6.3 state the categorial equivalence between the vector bundles over a smooth manifold $X$ and projective $C^{\infty}(X)$ modules of finite rank. Therefore, the differential calculus (including linear differential operators, linear connections) on vector bundles can be algebraically formulated as the differential calculus on these modules.

By a morphism of vector bundles is meant a linear bundle morphism, which is a fibrewise map whose restriction to each fibre is a linear map.

Given a linear bundle morphism $\Phi: Y^{\prime} \rightarrow Y$ of vector bundles over $X$, its kernel $\operatorname{Ker} \Phi$ is defined as the inverse image $\Phi^{-1}(\widehat{0}(X))$ of the canonical zero-valued section $\widehat{0}(X)$ of $Y$. By virtue of Theorem 1.1.10, if $\Phi$ is of constant rank, its kernel and its range are vector subbundles of
the vector bundles $Y^{\prime}$ and $Y$, respectively. For instance, monomorphisms and epimorphisms of vector bundles fulfil this condition.

There are the following particular constructions of new vector bundles from the old ones.

- Let $Y \rightarrow X$ be a vector bundle with a typical fibre $V$. By $Y^{*} \rightarrow X$ is denoted the dual vector bundle with the typical fibre $V^{*}$ dual of $V$. The interior product of $Y$ and $Y^{*}$ is defined as a fibred morphism

$$
\rfloor: Y \otimes Y^{*} \underset{X}{\longrightarrow} X \times \mathbb{R}
$$

- Let $Y \rightarrow X$ and $Y^{\prime} \rightarrow X$ be vector bundles with typical fibres $V$ and $V^{\prime}$, respectively. Their Whitney $\operatorname{sum} Y \underset{X}{\oplus} Y^{\prime}$ is a vector bundle over $X$ with the typical fibre $V \oplus V^{\prime}$.
- Let $Y \rightarrow X$ and $Y^{\prime} \rightarrow X$ be vector bundles with typical fibres $V$ and $V^{\prime}$, respectively. Their tensor product $Y \underset{X}{\otimes} Y^{\prime}$ is a vector bundle over $X$ with the typical fibre $V \otimes V^{\prime}$. Similarly, the exterior product of vector bundles $Y \underset{X}{\wedge} Y^{\prime}$ is defined. The exterior product

$$
\begin{equation*}
\wedge Y=X \times \mathbb{R} \oplus_{X} Y \oplus_{X}{ }_{\wedge}^{2} Y \underset{X}{\oplus} \cdots \stackrel{k}{\wedge} Y, \quad k=\operatorname{dim} Y-\operatorname{dim} X \tag{1.2.2}
\end{equation*}
$$

is called the exterior bundle.
Remark 1.2.2: Given vector bundles $Y$ and $Y^{\prime}$ over the same base $X$, every linear bundle morphism

$$
\Phi: Y_{x} \ni\left\{e_{i}(x)\right\} \rightarrow\left\{\Phi_{i}^{k}(x) e_{k}^{\prime}(x)\right\} \in Y_{x}^{\prime}
$$

over $X$ defines a global section

$$
\Phi: x \rightarrow \Phi_{i}^{k}(x) e^{i}(x) \otimes e_{k}^{\prime}(x)
$$

of the tensor product $Y \otimes Y^{*}$, and vice versa.
A sequence

$$
Y^{\prime} \xrightarrow{i} Y \xrightarrow{j} Y^{\prime \prime}
$$

of vector bundles over the same base $X$ is called exact at $Y$ if $\operatorname{Ker} j=$ $\operatorname{Im} i$. A sequence of vector bundles

$$
\begin{equation*}
0 \rightarrow Y^{\prime} \xrightarrow{i} Y \xrightarrow{j} Y^{\prime \prime} \rightarrow 0 \tag{1.2.3}
\end{equation*}
$$

over $X$ is said to be a short exact sequence if it is exact at all terms $Y^{\prime}$, $Y$, and $Y^{\prime \prime}$. This means that $i$ is a bundle monomorphism, $j$ is a bundle epimorphism, and $\operatorname{Ker} j=\operatorname{Im} i$. Then $Y^{\prime \prime}$ is the factor bundle $Y / Y^{\prime}$ whose structure module is the quotient $Y(X) / Y^{\prime}(X)$ of the structure modules of $Y$ and $Y^{\prime}$. Given an exact sequence of vector bundles (1.2.3), there is the exact sequence of their duals

$$
0 \rightarrow Y^{\prime \prime *} \xrightarrow{j^{*}} Y^{*} \xrightarrow{i^{*}} Y^{\prime *} \rightarrow 0 .
$$

One says that an exact sequence (1.2.3) is split if there exists a bundle monomorphism $s: Y^{\prime \prime} \rightarrow Y$ such that $j \circ s=\operatorname{Id} Y^{\prime \prime}$ or, equivalently,

$$
Y=i\left(Y^{\prime}\right) \oplus s\left(Y^{\prime \prime}\right)=Y^{\prime} \oplus Y^{\prime \prime}
$$

Theorem 1.2.2: Every exact sequence of vector bundles (1.2.3) is split.

The tangent bundle $T Z$ and the cotangent bundle $T^{*} Z$ of a manifold $Z$ exemplify vector bundles.

Remark 1.2.3: Given an atlas $\Psi_{Z}=\left\{\left(U_{\iota}, \phi_{\iota}\right)\right\}$ of a manifold $Z$, the tangent bundle is provided with the holonomic bundle atlas

$$
\begin{equation*}
\Psi_{T}=\left\{\left(U_{\iota}, \psi_{\iota}=T \phi_{\iota}\right)\right\} \tag{1.2.4}
\end{equation*}
$$

where $T \phi_{\iota}$ is the tangent morphism to $\phi_{\iota}$. The associated linear bundle coordinates are holonomic (or induced) coordinates ( $\dot{z}^{\lambda}$ ) with respect to the holonomic frames $\left\{\partial_{\lambda}\right\}$ in tangent spaces $T_{z} Z$.

The tensor product of tangent and cotangent bundles

$$
\begin{equation*}
T=(\stackrel{m}{\otimes} T Z) \otimes\left(\stackrel{k}{\otimes} T^{*} Z\right), \quad m, k \in \mathbb{N} \tag{1.2.5}
\end{equation*}
$$

is called a tensor bundle, provided with holonomic bundle coordinates $\dot{x}_{\beta_{1} \cdots \beta_{k}}^{\alpha_{1} \cdots \alpha_{m}}$ possessing transition functions

$$
\dot{x}_{\beta_{1} \cdots \beta_{k}}^{\alpha_{1} \cdots \alpha_{m}}=\frac{\partial x^{\prime \alpha_{1}}}{\partial x^{\mu_{1}}} \cdots \frac{\partial x^{\prime \alpha_{m}}}{\partial x^{\mu_{m}}} \frac{\partial x^{\nu_{1}}}{\partial x^{\prime \beta_{1}}} \cdots \frac{\partial x^{\nu_{k}}}{\partial x^{\prime \beta_{k}}} \dot{x}_{\nu_{1} \cdots \nu_{k}}^{\mu_{1} \cdots \mu_{m}} .
$$

Let $\pi_{Y}: T Y \rightarrow Y$ be the tangent bundle of a fibre bundle $\pi: Y \rightarrow$ $X$. Given bundle coordinates $\left(x^{\lambda}, y^{i}\right)$ on $Y$, it is equipped with the holonomic coordinates $\left(x^{\lambda}, y^{i}, \dot{x}^{\lambda}, \dot{y}^{i}\right)$. The tangent bundle $T Y \rightarrow Y$ has the subbundle $V Y=\operatorname{Ker}(T \pi)$, which consists of the vectors tangent to fibres of $Y$. It is called the vertical tangent bundle of $Y$ and is provided with the holonomic coordinates $\left(x^{\lambda}, y^{i}, \dot{y}^{i}\right)$ with respect to the vertical frames $\left\{\partial_{i}\right\}$. Every bundle morphism $\Phi: Y \rightarrow Y^{\prime}$ yields the linear bundle morphism over $\Phi$ of the vertical tangent bundles

$$
\begin{equation*}
V \Phi: V Y \rightarrow V Y^{\prime}, \quad \dot{y}^{\prime i} \circ V \Phi=\frac{\partial \Phi^{i}}{\partial y^{j}} \dot{y}^{j} . \tag{1.2.6}
\end{equation*}
$$

It is called the vertical tangent morphism.
In many important cases, the vertical tangent bundle $V Y \rightarrow Y$ of a fibre bundle $Y \rightarrow X$ is trivial, and is isomorphic to the bundle product

$$
\begin{equation*}
V Y=\underset{X}{\times} \underset{Y}{\bar{Y}} \tag{1.2.7}
\end{equation*}
$$

where $\bar{Y} \rightarrow X$ is some vector bundle. It follows that $V Y$ can be provided with bundle coordinates $\left(x^{\lambda}, y^{i}, \bar{y}^{i}\right)$ such that transition functions of coordinates $\bar{y}^{i}$ are independent of coordinates $y^{i}$. One calls (1.2.7) the vertical splitting. For instance, every vector bundle $Y \rightarrow X$ admits the canonical vertical splitting

$$
\begin{equation*}
V Y=Y \underset{X}{\oplus} Y \tag{1.2.8}
\end{equation*}
$$

The vertical cotangent bundle $V^{*} Y \rightarrow Y$ of a fibre bundle $Y \rightarrow X$ is defined as the dual of the vertical tangent bundle $V Y \rightarrow Y$. It is not a subbundle of the cotangent bundle $T^{*} Y$, but there is the canonical surjection

$$
\begin{equation*}
\zeta: T^{*} Y \ni \dot{x}_{\lambda} d x^{\lambda}+\dot{y}_{i} d y^{i} \rightarrow \dot{y}_{i} \bar{d} y^{i} \in V^{*} Y \tag{1.2.9}
\end{equation*}
$$

where $\left\{\bar{d} y^{i}\right\}$, possessing transition functions

$$
\bar{d} y^{\prime i}=\frac{\partial y^{\prime i}}{\partial y^{j}} \bar{d} y^{j}
$$

are the duals of the holonomic frames $\left\{\partial_{i}\right\}$ of $V Y$.
For any fibre bundle $Y$, there exist the exact sequences of vector bundles

$$
\begin{align*}
& 0 \rightarrow V Y \rightarrow T Y \xrightarrow{\pi_{T}} Y \underset{X}{\times} T X \rightarrow 0  \tag{1.2.10}\\
& 0 \rightarrow Y \underset{X}{\times} T^{*} X \rightarrow T^{*} Y \rightarrow V^{*} Y \rightarrow 0 \tag{1.2.11}
\end{align*}
$$

Their splitting, by definition, is a connection on $Y \rightarrow X$.
For the sake of simplicity, we agree to denote the pull-backs

$$
\underset{X}{\times} T X, \quad Y \underset{X}{\times} T^{*} X
$$

by $T X$ and $T^{*} X$, respectively.
Let $\bar{\pi}: \bar{Y} \rightarrow X$ be a vector bundle with a typical fibre $\bar{V}$. An affine bundle modelled over the vector bundle $\bar{Y} \rightarrow X$ is a fibre bundle $\pi: Y \rightarrow X$ whose typical fibre $V$ is an affine space modelled over $\bar{V}$ such that the following conditions hold.

- All the fibres $Y_{x}$ of $Y$ are affine spaces modelled over the corresponding fibres $\bar{Y}_{x}$ of the vector bundle $\bar{Y}$.
- There is an affine bundle atlas

$$
\Psi=\left\{\left(U_{\alpha}, \psi_{\chi}\right), \varrho_{\chi \zeta}\right\}
$$

of $Y \rightarrow X$ whose local trivializations morphisms $\psi_{\chi}(1.1 .5)$ and transition functions $\varrho_{\chi \zeta}$ (1.1.6) are affine isomorphisms.

Dealing with affine bundles, we use only affine bundle coordinates ( $y^{i}$ ) associated to an affine bundle atlas $\Psi$. There are the bundle morphisms

$$
\begin{array}{ll}
Y \times \bar{Y} \bar{X} \\
\underset{X}{\times} Y, & \left(y^{i}, \bar{y}^{i}\right) \rightarrow y^{i}+\bar{y}^{i} \\
Y \underset{X}{\times} \\
Y & \left(y^{i}, y^{\prime i}\right) \rightarrow y^{i}-y^{\prime i}
\end{array}
$$

where $\left(\bar{y}^{i}\right)$ are linear coordinates on the vector bundle $\bar{Y}$.

By virtue of Theorem 1.1.4, affine bundles have global sections, but in contrast with vector bundles, there is no canonical global section of an affine bundle. Let $\pi: Y \rightarrow X$ be an affine bundle. Every global section $s$ of an affine bundle $Y \rightarrow X$ modelled over a vector bundle $\bar{Y} \rightarrow X$ yields the bundle morphisms

$$
\begin{align*}
& Y \ni y \rightarrow y-s(\pi(y)) \in \bar{Y}  \tag{1.2.12}\\
& \bar{Y} \ni \bar{y} \rightarrow s(\pi(y))+\bar{y} \in Y \tag{1.2.13}
\end{align*}
$$

In particular, every vector bundle $Y$ has a natural structure of an affine bundle due to the morphisms (1.2.13) where $s=\hat{0}$ is the canonical zerovalued section of $Y$. For instance, the tangent bundle $T X$ of a manifold $X$ is naturally an affine bundle $A T X$ called the affine tangent bundle.

Theorem 1.2.3: Any affine bundle $Y \rightarrow X$ admits bundle coordinates $\left(x^{\lambda}, \tilde{y}^{i}\right)$ with linear transition functions $\widetilde{y}^{i}=A_{j}^{i}(x) \tilde{y}^{j}$ (see Example 4.8.2).

By a morphism of affine bundles is meant a bundle morphism $\Phi$ : $Y \rightarrow Y^{\prime}$ whose restriction to each fibre of $Y$ is an affine map. It is called an affine bundle morphism. Every affine bundle morphism $\Phi$ : $Y \rightarrow Y^{\prime}$ of an affine bundle $Y$ modelled over a vector bundle $\bar{Y}$ to an affine bundle $Y^{\prime}$ modelled over a vector bundle $\bar{Y}^{\prime}$ yields an unique linear bundle morphism

$$
\begin{equation*}
\bar{\Phi}: \bar{Y} \rightarrow \bar{Y}^{\prime}, \quad \bar{y}^{i} \circ \bar{\Phi}=\frac{\partial \Phi^{i}}{\partial y^{j}} \bar{y}^{j} \tag{1.2.14}
\end{equation*}
$$

called the linear derivative of $\Phi$.
Similarly to vector bundles, if $\Phi: Y \rightarrow Y^{\prime}$ is an affine morphism of affine bundles of constant rank, then $\Phi(Y)$ and $\operatorname{Ker} \Phi$ are affine subbundles of $Y^{\prime}$ and $Y$, respectively.

Every affine bundle $Y \rightarrow X$ modelled over a vector bundle $\bar{Y} \rightarrow X$ admits the canonical vertical splitting

$$
\begin{equation*}
V Y=\underset{X}{\times} \bar{Y} \tag{1.2.15}
\end{equation*}
$$

Note that Theorems 1.1.8 and 1.1.9 on a particular cover for bundle atlases remain true in the case of linear and affine atlases of vector and affine bundles.

### 1.3 Vector fields

Vector fields on a manifold $Z$ are global sections of the tangent bundle $T Z \rightarrow Z$.

The set $\mathcal{T}(Z)$ of vector fields on $Z$ is both a $C^{\infty}(Z)$-module and a real Lie algebra with respect to the Lie bracket

$$
\begin{aligned}
& u=u^{\lambda} \partial_{\lambda}, \quad v=v^{\lambda} \partial_{\lambda} \\
& {[v, u]=\left(v^{\lambda} \partial_{\lambda} u^{\mu}-u^{\lambda} \partial_{\lambda} v^{\mu}\right) \partial_{\mu}}
\end{aligned}
$$

Given a vector field $u$ on $X$, a curve

$$
c: \mathbb{R} \supset(,) \rightarrow Z
$$

in $Z$ is said to be an integral curve of $u$ if $T c=u(c)$. Every vector field $u$ on a manifold $Z$ can be seen as an infinitesimal generator of a local one-parameter group of diffeomorphisms (a flow), and vice versa. Onedimensional orbits of this group are integral curves of $u$. A vector field is called complete if its flow is a one-parameter group of diffeomorphisms of $Z$. For instance, every vector field on a compact manifold is complete.

A vector field $u$ on a fibre bundle $Y \rightarrow X$ is called projectable if it projects onto a vector field on $X$, i.e., there exists a vector field $\tau$ on $X$ such that

$$
\tau \circ \pi=T \pi \circ u
$$

A projectable vector field takes the coordinate form

$$
\begin{equation*}
u=u^{\lambda}\left(x^{\mu}\right) \partial_{\lambda}+u^{i}\left(x^{\mu}, y^{j}\right) \partial_{i}, \quad \tau=u^{\lambda} \partial_{\lambda} \tag{1.3.1}
\end{equation*}
$$

Its flow is a local one-parameter group of automorphisms of $Y \rightarrow X$ over a local one-parameter group of diffeomorphisms of $X$ whose generator
is $\tau$. A projectable vector field is called vertical if its projection onto $X$ vanishes, i.e., if it lives in the vertical tangent bundle $V Y$.

A vector field $\tau=\tau^{\lambda} \partial_{\lambda}$ on a base $X$ of a fibre bundle $Y \rightarrow X$ gives rise to a vector field on $Y$ by means of a connection on this fibre bundle (see the formula (3.1.6)). Nevertheless, every tensor bundle (1.2.5) admits the canonical lift of vector fields

$$
\begin{equation*}
\tilde{\tau}=\tau^{\mu} \partial_{\mu}+\left[\partial_{\nu} \tau^{\alpha_{1}} \dot{x}_{\beta_{1} \cdots \beta_{k}}^{\nu \alpha_{2} \cdots \alpha_{m}}+\ldots-\partial_{\beta_{1}} \tau^{\nu} \dot{x}_{\nu \beta_{2} \cdots \beta_{k}}^{\alpha_{1} \cdots \alpha_{m}}-\ldots\right] \dot{\partial}_{\alpha_{1} \cdots \alpha_{m}}^{\beta_{1} \cdots \beta_{k}} \tag{1.3.2}
\end{equation*}
$$

where we employ the compact notation

$$
\begin{equation*}
\dot{\partial}_{\lambda}=\frac{\partial}{\partial \dot{x}^{\lambda}} . \tag{1.3.3}
\end{equation*}
$$

This lift is functorial, i.e., it is an $\mathbb{R}$-linear monomorphism of the Lie algebra $\mathcal{T}(X)$ of vector fields on $X$ to the Lie algebra $\mathcal{T}(Y)$ of vector fields on $Y$ (see Section 5.1). In particular, we have the functorial lift

$$
\begin{equation*}
\widetilde{\tau}=\tau^{\mu} \partial_{\mu}+\partial_{\nu} \tau^{\alpha} \dot{x}^{\nu} \frac{\partial}{\partial \dot{x}^{\alpha}} \tag{1.3.4}
\end{equation*}
$$

of vector fields on $X$ onto the tangent bundle $T X$ and their functorial lift

$$
\begin{equation*}
\widetilde{\tau}=\tau^{\mu} \partial_{\mu}-\partial_{\beta} \tau^{\nu} \dot{x}_{\nu} \frac{\partial}{\partial \dot{x}_{\beta}} \tag{1.3.5}
\end{equation*}
$$

onto the cotangent bundle $T^{*} X$.
A fibre bundle admitting functorial lift of vector fields on its base is called the natural bundle (see Chapter 5).

A subbundle $\mathbf{T}$ of the tangent bundle $T Z$ of a manifold $Z$ is called a regular distribution (or, simply, a distribution). A vector field $u$ on $Z$ is said to be subordinate to a distribution $\mathbf{T}$ if it lives in $\mathbf{T}$. A distribution $\mathbf{T}$ is called involutive if the Lie bracket of $\mathbf{T}$-subordinate vector fields also is subordinate to $\mathbf{T}$.

A subbundle of the cotangent bundle $T^{*} Z$ of $Z$ is called a codistribution $\mathbf{T}^{*}$ on a manifold $Z$. For instance, the annihilator Ann $\mathbf{T}$ of a distribution $\mathbf{T}$ is a codistribution whose fibre over $z \in Z$ consists of covectors $w \in T_{z}^{*}$ such that $\left.v\right\rfloor w=0$ for all $v \in \mathbf{T}_{z}$.

Theorem 1.3.1: Let $\mathbf{T}$ be a distribution and Ann $\mathbf{T}$ its annihilator. Let $\wedge$ Ann $\mathbf{T}(Z)$ be the ideal of the exterior algebra $\mathcal{O}^{*}(Z)$ which is generated by sections of Ann $\mathbf{T} \rightarrow Z$. A distribution $\mathbf{T}$ is involutive iff the ideal $\wedge \operatorname{Ann} \mathbf{T}(Z)$ is a differential ideal, i.e.,

$$
d(\wedge \operatorname{Ann} \mathbf{T}(Z)) \subset \wedge \operatorname{Ann} \mathbf{T}(Z)
$$

The following local coordinates can be associated to an involutive distribution.

Theorem 1.3.2: Let $\mathbf{T}$ be an involutive $r$-dimensional distribution on a manifold $Z, \operatorname{dim} Z=k$. Every point $z \in Z$ has an open neighborhood $U$ which is a domain of an adapted coordinate chart $\left(z^{1}, \ldots, z^{k}\right)$ such that, restricted to $U$, the distribution $\mathbf{T}$ and its annihilator Ann $\mathbf{T}$ are spanned by the local vector fields $\partial / \partial z^{1}, \cdots, \partial / \partial z^{r}$ and the one-forms $d z^{r+1}, \ldots, d z^{k}$, respectively.

A connected submanifold $N$ of a manifold $Z$ is called an integral manifold of a distribution $\mathbf{T}$ on $Z$ if $T N \subset \mathbf{T}$. Unless otherwise stated, by an integral manifold is meant an integral manifold of dimension of T. An integral manifold is called maximal if no other integral manifold contains it. The following is the classical theorem of Frobenius.

Theorem 1.3.3: Let $\mathbf{T}$ be an involutive distribution on a manifold $Z$. For any $z \in Z$, there exists a unique maximal integral manifold of $\mathbf{T}$ through $z$, and any integral manifold through $z$ is its open subset.

Maximal integral manifolds of an involutive distribution on a manifold $Z$ are assembled into a regular foliation $\mathcal{F}$ of $Z$. A regular $r$-dimensional foliation (or, simply, a foliation) $\mathcal{F}$ of a $k$-dimensional manifold $Z$ is defined as a partition of $Z$ into connected $r$-dimensional submanifolds (the leaves of a foliation) $F_{\iota}, \iota \in I$, which possesses the following properties. A foliated manifold $(Z, \mathcal{F})$ admits an adapted coordinate atlas

$$
\begin{equation*}
\left\{\left(U_{\xi} ; z^{\lambda} ; z^{i}\right)\right\}, \quad \lambda=1, \ldots, n-r, \quad i=1, \ldots, r \tag{1.3.6}
\end{equation*}
$$

such that transition functions of coordinates $z^{\lambda}$ are independent of the remaining coordinates $z^{i}$ and, for each leaf $F$ of a foliation $\mathcal{F}$, the connected components of $F \cap U_{\xi}$ are given by the equations $z^{\lambda}=$ const. These connected components and coordinates $\left(z^{i}\right)$ on them make up a coordinate atlas of a leaf $F$.

It should be emphasized that leaves of a foliation need not be closed or imbedded submanifolds. Every leaf has an open tubular neighborhood $U$, i.e., if $z \in U$, then a leaf through $z$ also belongs to $U$.

A pair $(Z, \mathcal{F})$ where $\mathcal{F}$ is a foliation of $Z$ is called a foliated manifold. For instance, any submersion $f: Z \rightarrow M$ yields a foliation

$$
\mathcal{F}=\left\{F_{p}=f^{-1}(p)\right\}_{p \in f(Z)}
$$

of $Z$ indexed by elements of $f(Z)$, which is an open submanifold of $M$, i.e., $Z \rightarrow f(Z)$ is a fibred manifold. Leaves of this foliation are closed imbedded submanifolds. Such a foliation is called simple. It is a fibred manifold over $f(Z)$. Any (regular) foliation is locally simple.

### 1.4 Exterior and tangent-valued forms

An exterior $r$-form on a manifold $Z$ is a section

$$
\phi=\frac{1}{r!} \phi_{\lambda_{1} \ldots \lambda_{r}} d z^{\lambda_{1}} \wedge \cdots \wedge d z^{\lambda_{r}}
$$

of the exterior product $\stackrel{r}{\wedge} T^{*} Z \rightarrow Z$, where

$$
\begin{aligned}
& d z^{\lambda_{1}} \wedge \cdots \wedge d z^{\lambda_{r}}=\frac{1}{r!} \epsilon^{\lambda_{1} \ldots \lambda_{r}}{ }_{\mu_{1} \ldots \mu_{r}} d x^{\mu_{1}} \otimes \cdots \otimes d x^{\mu_{r}}, \\
& \epsilon^{\ldots \lambda_{i} \ldots \lambda_{j} \ldots} \ldots \mu_{p \ldots \mu_{k} \ldots}=-\epsilon^{\ldots \lambda_{j} \ldots \lambda_{i} \ldots \ldots \mu_{p} \ldots \mu_{k} \ldots}=-\epsilon^{\ldots \lambda_{i} \ldots \lambda_{j} \ldots}{ }_{\lambda_{1} \ldots \mu_{k} \ldots \mu_{p} \ldots}, \\
& \epsilon^{\lambda_{1} \ldots \lambda_{r}}{ }_{\lambda_{1} \ldots \lambda_{r}}=1 .
\end{aligned}
$$

Let $\mathcal{O}^{r}(Z)$ denote the vector space of exterior $r$-forms on a manifold $Z$. By definition, $\mathcal{O}^{0}(Z)=C^{\infty}(Z)$ is the ring of smooth real functions on $Z$. All exterior forms on $Z$ constitute the $\mathbb{N}$-graded commutative algebra
$\mathcal{O}^{*}(Z)$ of global sections of the exterior bundle $\wedge T^{*} Z$ (1.2.2) endowed with the exterior product

$$
\begin{aligned}
& \phi= \frac{1}{r!} \phi_{\lambda_{1} \ldots \lambda_{r}} d z^{\lambda_{1}} \wedge \cdots \wedge d z^{\lambda_{r}}, \quad \sigma=\frac{1}{s!} \sigma_{\mu_{1} \ldots \mu_{s}} d z^{\mu_{1}} \wedge \cdots \wedge d z^{\mu_{s}}, \\
& \phi \wedge \sigma=\frac{1}{r!s!} \phi_{\nu_{1} \ldots \nu_{r}} \sigma_{\nu_{r+1} \ldots \nu_{r+s}} d z^{\nu_{1}} \wedge \cdots \wedge d z^{\nu_{r+s}}= \\
& \frac{1}{r!s!(r+s)!} \epsilon^{\nu_{1} \ldots \nu_{r+s}}{ }_{\alpha_{1} \ldots \alpha_{r+s}} \phi_{\nu_{1} \ldots \nu_{r}} \sigma_{\nu_{r+1} \ldots \nu_{r+s}} d z^{\alpha_{1}} \wedge \cdots \wedge d z^{\alpha_{r+s}},
\end{aligned}
$$

such that

$$
\phi \wedge \sigma=(-1)^{|\phi||\sigma|} \sigma \wedge \phi,
$$

where the symbol $|\phi|$ stands for the form degree. An algebra $\mathcal{O}^{*}(Z)$ also is provided with the exterior differential

$$
d \phi=d z^{\mu} \wedge \partial_{\mu} \phi=\frac{1}{r!} \partial_{\mu} \phi_{\lambda_{1} \ldots \lambda_{r}} d z^{\mu} \wedge d z^{\lambda_{1}} \wedge \cdots \wedge d z^{\lambda_{r}}
$$

which obeys the relations

$$
d \circ d=0, \quad d(\phi \wedge \sigma)=d(\phi) \wedge \sigma+(-1)^{|\phi|} \phi \wedge d(\sigma)
$$

The exterior differential $d$ makes $\mathcal{O}^{*}(Z)$ into a differential graded algebra (henceforth DGA) which is the minimal Chevalley-Eilenberg differential calculus $\mathcal{O}^{*} \mathcal{A}$ over the real ring $\mathcal{A}=C^{\infty}(Z)$. Its de Rham complex is (8.6.5).

Given a manifold morphism $f: Z \rightarrow Z^{\prime}$, any exterior $k$-form $\phi$ on $Z^{\prime}$ yields the pull-back exterior form $f^{*} \phi$ on $Z$ given by the condition

$$
f^{*} \phi\left(v^{1}, \ldots, v^{k}\right)(z)=\phi\left(T f\left(v^{1}\right), \ldots, T f\left(v^{k}\right)\right)(f(z))
$$

for an arbitrary collection of tangent vectors $v^{1}, \cdots, v^{k} \in T_{z} Z$. We have the relations

$$
f^{*}(\phi \wedge \sigma)=f^{*} \phi \wedge f^{*} \sigma, \quad d f^{*} \phi=f^{*}(d \phi)
$$

In particular, given a fibre bundle $\pi: Y \rightarrow X$, the pull-back onto $Y$ of exterior forms on $X$ by $\pi$ provides the monomorphism of graded
commutative algebras $\mathcal{O}^{*}(X) \rightarrow \mathcal{O}^{*}(Y)$. Elements of its range $\pi^{*} \mathcal{O}^{*}(X)$ are called basic forms. Exterior forms

$$
\phi: Y \rightarrow \stackrel{r}{\wedge} T^{*} X, \quad \phi=\frac{1}{r!} \phi_{\lambda_{1} \ldots \lambda_{r}} d x^{\lambda_{1}} \wedge \cdots \wedge d x^{\lambda_{r}}
$$

on $Y$ such that $u\rfloor \phi=0$ for an arbitrary vertical vector field $u$ on $Y$ are said to be horizontal forms. Horizontal forms of degree $n=\operatorname{dim} X$ are called densities. We use for them the compact notation

$$
\begin{align*}
& L=\frac{1}{n!} L_{\mu_{1} \ldots \mu_{n}} d x^{\mu_{1}} \wedge \cdots \wedge d x^{\mu_{n}}=\mathcal{L} \omega, \quad \mathcal{L}=L_{1 \ldots n}, \\
& \omega=d x^{1} \wedge \cdots \wedge d x^{n}=\frac{1}{n!} \epsilon_{\mu_{1} \ldots \mu_{n}} d x^{\mu_{1}} \wedge \cdots \wedge d x^{\mu_{n}},  \tag{1.4.1}\\
& \left.\left.\left.\omega_{\lambda}=\partial_{\lambda}\right\rfloor \omega, \quad \omega_{\mu \lambda}=\partial_{\mu}\right\rfloor \partial_{\lambda}\right\rfloor \omega
\end{align*}
$$

where $\epsilon$ is the skew-symmetric Levi-Civita symbol with the component $\epsilon_{\mu_{1} \ldots \mu_{n}}=1$.

The interior product (or contraction) of a vector field $u$ and an exterior $r$-form $\phi$ on a manifold $Z$ is given by the coordinate expression

$$
\begin{aligned}
u\rfloor \phi= & \sum_{k=1}^{r} \frac{(-1)^{k-1}}{r!} u^{\lambda_{k}} \phi_{\lambda_{1} \ldots \lambda_{k} \ldots \lambda_{r}} d z^{\lambda_{1}} \wedge \cdots \wedge \widehat{d z}^{\lambda_{k}} \wedge \cdots \wedge d z^{\lambda_{r}}= \\
& \frac{1}{(r-1)!} u^{\mu} \phi_{\mu \alpha_{2} \ldots \alpha_{r}} d z^{\alpha_{2}} \wedge \cdots \wedge d z^{\alpha_{r}}
\end{aligned}
$$

where the caret ^ denotes omission. It obeys the relations

$$
\begin{aligned}
& \left.\left.\phi\left(u_{1}, \ldots, u_{r}\right)=u_{r}\right\rfloor \cdots u_{1}\right\rfloor \phi \\
& \left.u\rfloor(\phi \wedge \sigma)=u\rfloor \phi \wedge \sigma+(-1)^{|\phi|} \phi \wedge u\right\rfloor \sigma
\end{aligned}
$$

The Lie derivative of an exterior form $\phi$ along a vector field $u$ is

$$
\begin{aligned}
& \left.\left.\mathbf{L}_{u} \phi=u\right\rfloor d \phi+d(u\rfloor \phi\right) \\
& \mathbf{L}_{u}(\phi \wedge \sigma)=\mathbf{L}_{u} \phi \wedge \sigma+\phi \wedge \mathbf{L}_{u} \sigma
\end{aligned}
$$

It is a derivation of the graded algebra $\mathcal{O}^{*}(Z)$ such that

$$
\mathbf{L}_{u} \circ \mathbf{L}_{u^{\prime}}-\mathbf{L}_{u^{\prime}} \circ \mathbf{L}_{u}=\mathbf{L}_{\left[u, u^{\prime}\right]} .
$$

In particular, if $f$ is a function, then

$$
\left.\mathbf{L}_{u} f=u(f)=u\right\rfloor d f
$$

An exterior form $\phi$ is invariant under a local one-parameter group of diffeomorphisms $G(t)$ of $Z$ (i.e., $G(t)^{*} \phi=\phi$ ) iff its Lie derivative along the infinitesimal generator $u$ of this group vanishes, i.e., $\mathbf{L}_{u} \phi=0$.

A tangent-valued $r$-form on a manifold $Z$ is a section

$$
\begin{equation*}
\phi=\frac{1}{r!} \phi_{\lambda_{1} \ldots \lambda_{r}}^{\mu} d z^{\lambda_{1}} \wedge \cdots \wedge d z^{\lambda_{r}} \otimes \partial_{\mu} \tag{1.4.2}
\end{equation*}
$$

of the tensor bundle

$$
\stackrel{r}{\wedge} T^{*} Z \otimes T Z \rightarrow Z
$$

Remark 1.4.1: There is one-to-one correspondence between the tangentvalued one-forms $\phi$ on a manifold $Z$ and the linear bundle endomorphisms

$$
\begin{align*}
& \left.\hat{\phi}: T Z \rightarrow T Z, \quad \hat{\phi}: T_{z} Z \ni v \rightarrow v\right\rfloor \phi(z) \in T_{z} Z,  \tag{1.4.3}\\
& \left.\hat{\phi}^{*}: T^{*} Z \rightarrow T^{*} Z, \quad \hat{\phi}^{*}: T_{z}^{*} Z \ni v^{*} \rightarrow \phi(z)\right\rfloor v^{*} \in T_{z}^{*} Z, \tag{1.4.4}
\end{align*}
$$

over $Z$ (see Remark 1.2.2). For instance, the canonical tangent-valued one-form

$$
\begin{equation*}
\theta_{Z}=d z^{\lambda} \otimes \partial_{\lambda} \tag{1.4.5}
\end{equation*}
$$

on $Z$ corresponds to the identity morphisms (1.4.3) and (1.4.4).
Remark 1.4.2: Let $Z=T X$, and let $T T X$ be the tangent bundle of $T X$. There is the bundle endomorphism

$$
\begin{equation*}
J\left(\partial_{\lambda}\right)=\dot{\partial}_{\lambda}, \quad J\left(\dot{\partial}_{\lambda}\right)=0 \tag{1.4.6}
\end{equation*}
$$

of $T T X$ over $X$. It corresponds to the canonical tangent-valued form

$$
\begin{equation*}
\theta_{J}=d x^{\lambda} \otimes \dot{\partial}_{\lambda} \tag{1.4.7}
\end{equation*}
$$

on the tangent bundle $T X$. It is readily observed that $J \circ J=0$.

The space $\mathcal{O}^{*}(Z) \otimes \mathcal{T}(Z)$ of tangent-valued forms is provided with the Frölicher-Nijenhuis bracket

$$
\begin{aligned}
& {[,]_{\mathrm{FN}}: \mathcal{O}^{r}(Z) \otimes \mathcal{T}(Z) \times \mathcal{O}^{s}(Z) \otimes \mathcal{T}(Z) \rightarrow \mathcal{O}^{r+s}(Z) \otimes \mathcal{T}(Z),} \\
& {[\alpha \otimes u, \beta \otimes v]_{\mathrm{FN}}=(\alpha \wedge \beta) \otimes[u, v]+\left(\alpha \wedge \mathbf{L}_{u} \beta\right) \otimes v-} \\
& \left.\left.\quad\left(\mathbf{L}_{v} \alpha \wedge \beta\right) \otimes u+(-1)^{r}(d \alpha \wedge u\rfloor \beta\right) \otimes v+(-1)^{r}(v\rfloor \alpha \wedge d \beta\right) \otimes u, \\
& \alpha \in \mathcal{O}^{r}(Z), \quad \beta \in \mathcal{O}^{s}(Z), \quad u, v \in \mathcal{T}(Z) .
\end{aligned}
$$

Its coordinate expression is

$$
\begin{aligned}
& {[\phi, \sigma]_{\mathrm{FN}}=\frac{1}{r!s!}\left(\phi_{\lambda_{1} \ldots \lambda_{r}}^{\nu} \partial_{\nu} \sigma_{\lambda_{r+1} \ldots \lambda_{r+s}}^{\mu}-\sigma_{\lambda_{r+1} \ldots \lambda_{r+s}}^{\nu} \partial_{\nu} \phi_{\lambda_{1} \ldots \lambda_{r}}^{\mu}-\right.} \\
& \left.\quad r \phi_{\lambda_{1} \ldots \lambda_{r-1} \nu}^{\mu} \partial_{\lambda_{r}} \sigma_{\lambda_{r+1} \ldots \lambda_{r+s}}^{\nu}+s \sigma_{\nu \lambda_{r+2} \ldots \lambda_{r+s}}^{\mu} \partial_{\lambda_{r+1}} \phi_{\lambda_{1} \ldots \lambda_{r}}^{\nu}\right) \\
& \quad d z^{\lambda_{1}} \wedge \cdots \wedge d z^{\lambda_{r+s}} \otimes \partial_{\mu}, \\
& \phi \in \mathcal{O}^{r}(Z) \otimes \mathcal{T}(Z), \quad \sigma \in \mathcal{O}^{s}(Z) \otimes \mathcal{T}(Z) .
\end{aligned}
$$

There are the relations

$$
\begin{align*}
& {[\phi, \sigma]_{\mathrm{FN}}=(-1)^{|\phi||\psi|+1}[\sigma, \phi]_{\mathrm{FN}},}  \tag{1.4.9}\\
& {\left[\phi,[\sigma, \theta]_{\mathrm{FN}},\right.}  \tag{1.4.10}\\
& \phi, \sigma, \theta \in]_{\mathrm{FN}}=\left[[\phi, \sigma]_{\mathrm{FN}}, \theta\right]_{\mathrm{FN}}+(-1)^{|\phi| \sigma \mid}\left[\sigma,[\phi, \theta]_{\mathrm{FN}}\right]_{\mathrm{FN}}, \\
& \mathcal{T}(Z)
\end{align*}
$$

Given a tangent-valued form $\theta$, the Nijenhuis differential on $\mathcal{O}^{*}(Z) \otimes$ $\mathcal{T}(Z)$ is defined as the morphism

$$
d_{\theta}: \psi \rightarrow d_{\theta} \psi=[\theta, \psi]_{\mathrm{FN}}, \quad \psi \in \mathcal{O}^{*}(Z) \otimes \mathcal{T}(Z) .
$$

By virtue of (1.4.10), it has the property

$$
d_{\phi}[\psi, \theta]_{\mathrm{FN}}=\left[d_{\phi} \psi, \theta\right]_{\mathrm{FN}}+(-1)^{|\phi||\psi|}\left[\psi, d_{\phi} \theta\right]_{\mathrm{FN}} .
$$

In particular, if $\theta=u$ is a vector field, the Nijenhuis differential is the Lie derivative of tangent-valued forms

$$
\begin{aligned}
& \mathbf{L}_{u} \sigma=d_{u} \sigma=[u, \sigma]_{\mathrm{FN}}=\frac{1}{s!}\left(u^{\nu} \partial_{\nu} \sigma_{\lambda_{1} \ldots \lambda_{s}}^{\mu}-\sigma_{\lambda_{1} \ldots \lambda_{s}}^{\nu} \partial_{\nu} u^{\mu}+\right. \\
& \left.\quad s \sigma_{\nu \lambda_{2} \ldots \lambda_{s}}^{\mu} \partial_{\lambda_{1}} u^{\nu}\right) d x^{\lambda_{1}} \wedge \cdots \wedge d x^{\lambda_{s}} \otimes \partial_{\mu}, \quad \sigma \in \mathcal{O}^{s}(Z) \otimes \mathcal{T}(Z) .
\end{aligned}
$$

Let $Y \rightarrow X$ be a fibre bundle. We consider the following subspaces of the space $\mathcal{O}^{*}(Y) \otimes \mathcal{T}(Y)$ of tangent-valued forms on $Y$ :

- horizontal tangent-valued forms

$$
\begin{aligned}
& \phi: Y \rightarrow \stackrel{r}{\wedge} T^{*} X \underset{Y}{\otimes} T Y \\
& \phi=d x^{\lambda_{1}} \wedge \cdots \wedge d x^{\lambda_{r}} \otimes \frac{1}{r!}\left[\phi_{\lambda_{1} \ldots \lambda_{r}}^{\mu}(y) \partial_{\mu}+\phi_{\lambda_{1} \ldots \lambda_{r}}^{i}(y) \partial_{i}\right]
\end{aligned}
$$

- projectable horizontal tangent-valued forms

$$
\phi=d x^{\lambda_{1}} \wedge \cdots \wedge d x^{\lambda_{r}} \otimes \frac{1}{r!}\left[\phi_{\lambda_{1} \ldots \lambda_{r}}^{\mu}(x) \partial_{\mu}+\phi_{\lambda_{1} \ldots \lambda_{r}}^{i}(y) \partial_{i}\right]
$$

- vertical-valued form

$$
\begin{aligned}
& \phi: Y \rightarrow \stackrel{r}{\wedge} T^{*} X \underset{Y}{\otimes} V Y \\
& \phi=\frac{1}{r!} \phi_{\lambda_{1} \ldots \lambda_{r}}^{i}(y) d x^{\lambda_{1}} \wedge \cdots \wedge d x^{\lambda_{r}} \otimes \partial_{i}
\end{aligned}
$$

- vertical-valued one-forms, called soldering forms,

$$
\sigma=\sigma_{\lambda}^{i}(y) d x^{\lambda} \otimes \partial_{i}
$$

- basic soldering forms

$$
\sigma=\sigma_{\lambda}^{i}(x) d x^{\lambda} \otimes \partial_{i}
$$

Remark 1.4.3: The tangent bundle $T X$ is provided with the canonical soldering form $\theta_{J}$ (1.4.7). Due to the canonical vertical splitting

$$
\begin{equation*}
V T X=T X \underset{X}{\times} T X \tag{1.4.11}
\end{equation*}
$$

the canonical soldering form (1.4.7) on $T X$ defines the canonical tangentvalued form $\theta_{X}(1.4 .5)$ on $X$. By this reason, tangent-valued one-forms on a manifold $X$ also are called soldering forms.

Remark 1.4.4: Let $Y \rightarrow X$ be a fibre bundle, $f: X^{\prime} \rightarrow X$ a manifold morphism, $f^{*} Y \rightarrow X^{\prime}$ the pull-back of $Y$ by $f$, and

$$
f_{Y}: f^{*} Y \rightarrow Y
$$

the corresponding bundle morphism (1.1.9). Since

$$
V f^{*} Y=f^{*} V Y=f_{Y}^{*} V Y, \quad V_{y^{\prime}} Y^{\prime}=V_{f_{Y}\left(y^{\prime}\right)} Y
$$

one can define the pull-back $f^{*} \phi$ onto $f^{*} Y$ of any vertical-valued form $\phi$ on $Y$ in accordance with the relation

$$
f^{*} \phi\left(v^{1}, \ldots, v^{r}\right)\left(y^{\prime}\right)=\phi\left(T f_{Y}\left(v^{1}\right), \ldots, T f_{Y}\left(v^{r}\right)\right)\left(f_{Y}\left(y^{\prime}\right)\right) .
$$

We also mention the $T X$-valued forms

$$
\begin{align*}
& \phi: Y \rightarrow \stackrel{r}{\wedge} T^{*} X \underset{Y}{\otimes} T X  \tag{1.4.12}\\
& \phi=\frac{1}{r!} \phi_{\lambda_{1} \ldots \lambda_{r}}^{\mu} d x^{\lambda_{1}} \wedge \cdots \wedge d x^{\lambda_{r}} \otimes \partial_{\mu}
\end{align*}
$$

and $V^{*} Y$-valued forms

$$
\begin{align*}
& \phi: Y \rightarrow \stackrel{r}{\wedge} T^{*} X \underset{Y}{\otimes} V^{*} Y,  \tag{1.4.13}\\
& \phi=\frac{1}{r!} \phi_{\lambda_{1} \ldots \lambda_{r} i} d x^{\lambda_{1}} \wedge \cdots \wedge d x^{\lambda_{r}} \otimes \bar{d} y^{i} .
\end{align*}
$$

It should be emphasized that (1.4.12) are not tangent-valued forms, while (1.4.13) are not exterior forms. They exemplify vector-valued forms. Given a vector bundle $E \rightarrow X$, by a $E$-valued $k$-form on $X$, is meant a section of the fibre bundle

$$
\left(\stackrel{k}{\wedge} T^{*} X\right) \underset{X}{\otimes} E^{*} \rightarrow X
$$

## Chapter 2

## Jet manifolds

There are different notions of jets. Here we are concerned with jets of sections of fibre bundles. They are the particular jets of maps and the jets of submanifolds. Let us also mention the jets of modules over a commutative ring. In particular, given a smooth manifold $X$, the jets of a projective $C^{\infty}(X)$-module $P$ of finite rank are exactly the jets of sections of the vector bundle over $X$ whose module of sections is $P$ in accordance with the Serre-Swan Theorem 8.6.3. The notion of jets is extended to modules over graded commutative rings. However, the jets of modules over a noncommutative ring can not be defined.

### 2.1 First order jet manifolds

Given a fibre bundle $Y \rightarrow X$ with bundle coordinates $\left(x^{\lambda}, y^{i}\right)$, let us consider the equivalence classes $j_{x}^{1} s$ of its sections $s$, which are identified by their values $s^{i}(x)$ and the values of their partial derivatives $\partial_{\mu} s^{i}(x)$ at a point $x \in X$. They are called the first order jets of sections at $x$. One can justify that the definition of jets is coordinate-independent. The key point is that the set $J^{1} Y$ of first order jets $j_{x}^{1} s, x \in X$, is a smooth manifold with respect to the adapted coordinates $\left(x^{\lambda}, y^{i}, y_{\lambda}^{i}\right)$ such that

$$
\begin{equation*}
y_{\lambda}^{i}\left(j_{x}^{1} s\right)=\partial_{\lambda} s^{i}(x), \quad y_{\lambda}^{\prime i}=\frac{\partial x^{\mu}}{\partial x^{\prime \lambda}}\left(\partial_{\mu}+y_{\mu}^{j} \partial_{j}\right) y^{\prime i} . \tag{2.1.1}
\end{equation*}
$$

It is called the first order jet manifold of a fibre bundle $Y \rightarrow X$. We call $\left(y_{\lambda}^{i}\right)$ the jet coordinates.

The jet manifold $J^{1} Y$ admits the natural fibrations

$$
\begin{align*}
& \pi^{1}: J^{1} Y \ni j_{x}^{1} s \rightarrow x \in X  \tag{2.1.2}\\
& \pi_{0}^{1}: J^{1} Y \ni j_{x}^{1} s \rightarrow s(x) \in Y \tag{2.1.3}
\end{align*}
$$

A glance at the transformation law (2.1.1) shows that $\pi_{0}^{1}$ is an affine bundle modelled over the vector bundle

$$
\begin{equation*}
T^{*} X \underset{Y}{\otimes} V Y \rightarrow Y \tag{2.1.4}
\end{equation*}
$$

It is convenient to call $\pi^{1}(2.1 .2)$ the jet bundle, while $\pi_{0}^{1}(2.1 .3)$ is said to be the affine jet bundle.

Let us note that, if $Y \rightarrow X$ is a vector or an affine bundle, the jet bundle $\pi_{1}$ (2.1.2) is so.

Jets can be expressed in terms of familiar tangent-valued forms as follows. There are the canonical imbeddings

$$
\begin{align*}
& \lambda_{(1)}: J^{1} Y \underset{Y}{\rightarrow} T^{*} X \underset{Y}{\otimes} T Y, \\
& \lambda_{(1)}=d x^{\lambda} \otimes\left(\partial_{\lambda}+y_{\lambda}^{i} \partial_{i}\right)=d x^{\lambda} \otimes d_{\lambda},  \tag{2.1.5}\\
& \theta_{(1)}: J^{1} Y \underset{Y}{\rightarrow} T^{*} Y \underset{Y}{\otimes} V Y, \\
& \theta_{(1)}=\left(d y^{i}-y_{\lambda}^{i} d x^{\lambda}\right) \otimes \partial_{i}=\theta^{i} \otimes \partial_{i}, \tag{2.1.6}
\end{align*}
$$

where $d_{\lambda}$ are called total derivatives, and $\theta^{i}$ are local contact forms.
Remark 2.1.1: We further identify the jet manifold $J^{1} Y$ with its images under the canonical morphisms (2.1.5) and (2.1.6), and represent the jets $j_{x}^{1} s=\left(x^{\lambda}, y^{i}, y_{\mu}^{i}\right)$ by the tangent-valued forms $\lambda_{(1)}$ (2.1.5) and $\theta_{(1)}(2.1 .6)$.

Any section $s$ of $Y \rightarrow X$ has the jet prolongation to the section

$$
\left(J^{1} s\right)(x)=j_{x}^{1} s, \quad y_{\lambda}^{i} \circ J^{1} s=\partial_{\lambda} s^{i}(x)
$$

of the jet bundle $J^{1} Y \rightarrow X$. A section of the jet bundle $J^{1} Y \rightarrow X$ is called integrable if it is the jet prolongation of some section of $Y \rightarrow X$.

Remark 2.1.2: By virtue of Theorem 1.1.4, the affine jet bundle $J^{1} Y \rightarrow Y$ admits global sections. If $Y$ is trivial, there is the canonical zero section $\widehat{0}(Y)$ of $J^{1} Y \rightarrow Y$ taking its values into centers of its affine fibres.

Any bundle morphism $\Phi: Y \rightarrow Y^{\prime}$ over a diffeomorphism $f$ admits a jet prolongation to a bundle morphism of affine jet bundles

$$
\begin{equation*}
J^{1} \Phi: J^{1} Y \underset{\Phi}{\longrightarrow} J^{1} Y^{\prime}, \quad y_{\lambda}^{\prime i} \circ J^{1} \Phi=\frac{\partial\left(f^{-1}\right)^{\mu}}{\partial x^{\prime \lambda}} d_{\mu} \Phi^{i} . \tag{2.1.7}
\end{equation*}
$$

Any projectable vector field $u$ (1.3.1) on a fibre bundle $Y \rightarrow X$ has a jet prolongation to the projectable vector field

$$
\begin{align*}
& J^{1} u=r_{1} \circ J^{1} u: J^{1} Y \rightarrow J^{1} T Y \rightarrow T J^{1} Y, \\
& J^{1} u=u^{\lambda} \partial_{\lambda}+u^{i} \partial_{i}+\left(d_{\lambda} u^{i}-y_{\mu}^{i} \partial_{\lambda} u^{\mu}\right) \partial_{i}^{\lambda}, \tag{2.1.8}
\end{align*}
$$

on the jet manifold $J^{1} Y$. To obtain (2.1.8), the canonical bundle morphism

$$
r_{1}: J^{1} T Y \rightarrow T J^{1} Y, \quad \dot{y}_{\lambda}^{i} \circ r_{1}=\left(\dot{y}^{i}\right)_{\lambda}-y_{\mu}^{i} \dot{x}_{\lambda}^{\mu}
$$

is used. In particular, there is the canonical isomorphism

$$
\begin{equation*}
V J^{1} Y=J^{1} V Y, \quad \dot{y}_{\lambda}^{i}=\left(\dot{y}^{i}\right)_{\lambda} . \tag{2.1.9}
\end{equation*}
$$

### 2.2 Higher order jet manifolds

The notion of first jets of sections of a fibre bundle is naturally extended to higher order jets.

Let $Y \rightarrow X$ be a fibre bundle. Given its bundle coordinates $\left(x^{\lambda}, y^{i}\right)$, a multi-index $\Lambda$ of the length $|\Lambda|=k$ throughout denotes a collection of indices $\left(\lambda_{1} \ldots \lambda_{k}\right)$ modulo permutations. By $\Lambda+\Sigma$ is meant a multi-index $\left(\lambda_{1} \ldots \lambda_{k} \sigma_{1} \ldots \sigma_{r}\right)$. For instance $\lambda+\Lambda=\left(\lambda \lambda_{1} \ldots \lambda_{r}\right)$. By $\Lambda \Sigma$ is denoted the union of collections ( $\lambda_{1} \ldots \lambda_{k} ; \sigma_{1} \ldots \sigma_{r}$ ) where the indices $\lambda_{i}$ and $\sigma_{j}$
are not permitted. Summation over a multi-index $\Lambda$ means separate summation over each its index $\lambda_{i}$. We use the compact notation

$$
\partial_{\Lambda}=\partial_{\lambda_{k}} \circ \cdots \circ \partial_{\lambda_{1}}, \quad \Lambda=\left(\lambda_{1} \ldots \lambda_{k}\right)
$$

The $r$-order jet manifold $J^{r} Y$ of sections of a bundle $Y \rightarrow X$ is defined as the disjoint union of equivalence classes $j_{x}^{r} s$ of sections $s$ of $Y \rightarrow X$ such that sections $s$ and $s^{\prime}$ belong to the same equivalence class $j_{x}^{r} s$ iff

$$
s^{i}(x)=s^{\prime i}(x), \quad \partial_{\Lambda} s^{i}(x)=\partial_{\Lambda} s^{\prime i}(x), \quad 0<|\Lambda| \leq r
$$

In brief, one can say that sections of $Y \rightarrow X$ are identified by the $r+1$ terms of their Taylor series at points of $X$. The particular choice of coordinates does not matter for this definition. The equivalence classes $j_{x}^{r} s$ are called the $r$-order jets of sections. Their set $J^{r} Y$ is endowed with an atlas of the adapted coordinates

$$
\begin{align*}
& \left(x^{\lambda}, y_{\Lambda}^{i}\right), \quad y_{\Lambda}^{i} \circ s=\partial_{\Lambda} s^{i}(x), \quad 0 \leq|\Lambda| \leq r  \tag{2.2.1}\\
& y_{\lambda+\Lambda}^{\prime i}=\frac{\partial x^{\mu}}{\partial^{\prime} x^{\lambda}} d_{\mu} y_{\Lambda}^{\prime i} \tag{2.2.2}
\end{align*}
$$

where the symbol $d_{\lambda}$ stands for the higher order total derivative

$$
\begin{equation*}
d_{\lambda}=\partial_{\lambda}+\sum_{0 \leq|\Lambda| \leq r-1} y_{\Lambda+\lambda}^{i} \partial_{i}^{\Lambda}, \quad d_{\lambda}^{\prime}=\frac{\partial x^{\mu}}{\partial x^{\prime \lambda}} d_{\mu} \tag{2.2.3}
\end{equation*}
$$

These derivatives act on exterior forms on $J^{r} Y$ and obey the relations

$$
\begin{aligned}
& {\left[d_{\lambda}, d_{\mu}\right]=0, \quad d_{\lambda} \circ d=d \circ d_{\lambda}} \\
& d_{\lambda}(\phi \wedge \sigma)=d_{\lambda}(\phi) \wedge \sigma+\phi \wedge d_{\lambda}(\sigma), \quad d_{\lambda}(d \phi)=d\left(d_{\lambda}(\phi)\right) \\
& d_{\lambda}\left(d x^{\mu}\right)=0, \quad d_{\lambda}\left(d y_{\Lambda}^{i}\right)=d y_{\lambda+\Lambda}^{i}
\end{aligned}
$$

We use the compact notation

$$
d_{\Lambda}=d_{\lambda_{r}} \circ \cdots \circ d_{\lambda_{1}}, \quad \Lambda=\left(\lambda_{r} \ldots \lambda_{1}\right)
$$

The coordinates (2.2.1) bring the set $J^{r} Y$ into a finite-dimensional manifold. The coordinates (2.2.1) are compatible with the natural surjections

$$
\pi_{k}^{r}: J^{r} Y \rightarrow J^{k} Y, \quad r>k
$$

which form the composite bundle

$$
\pi^{r}: J^{r} Y \xrightarrow{\pi_{r-1}^{r}} J^{r-1} Y \xrightarrow{\pi_{r-2}^{r-1}} \cdots \xrightarrow{\pi_{0}^{1}} Y \quad{ }^{\pi}
$$

with the properties

$$
\pi_{s}^{k} \circ \pi_{k}^{r}=\pi_{s}^{r}, \quad \pi^{s} \circ \pi_{s}^{r}=\pi^{r}
$$

A glance at the transition functions (2.2.2) shows that the fibration

$$
\pi_{r-1}^{r}: J^{r} Y \rightarrow J^{r-1} Y
$$

is an affine bundle modelled over the vector bundle

$$
\begin{equation*}
\stackrel{r}{\vee} T^{*} X \underset{J^{r-1} Y}{\otimes} V Y \rightarrow J^{r-1} Y . \tag{2.2.4}
\end{equation*}
$$

Remark 2.2.1: Let us recall that a base of any affine bundle is a strong deformation retract of its total space. Consequently, $Y$ is a strong deformation retract of $J^{1} Y$, which in turn is a strong deformation retract of $J^{2} Y$, and so on. It follows that a fibre bundle $Y$ is a strong deformation retract of any finite order jet manifold $J^{r} Y$. Therefore, by virtue of the Vietoris-Begle theorem, there is an isomorphism

$$
\begin{equation*}
H^{*}\left(J^{r} Y ; \mathbb{R}\right)=H^{*}(Y ; \mathbb{R}) \tag{2.2.5}
\end{equation*}
$$

of cohomology of $J^{r} Y$ and $Y$ with coefficients in the constant sheaf $\mathbb{R}$.

Remark 2.2.2: To introduce higher order jet manifolds, one can use the construction of repeated jet manifolds. Let us consider the $r$-order jet manifold $J^{r} J^{k} Y$ of a jet bundle $J^{k} Y \rightarrow X$. It is coordinated by $\left(x^{\mu}, y_{\Sigma \Lambda}^{i}\right),|\Lambda| \leq k,|\Sigma| \leq r$. There is a canonical monomorphism

$$
\sigma_{r k}: J^{r+k} Y \rightarrow J^{r} J^{k} Y, \quad y_{\Sigma \Lambda}^{i} \circ \sigma_{r k}=y_{\Sigma+\Lambda}^{i}
$$

In the calculus in higher order jets, we have the $r$-order jet prolongation functor such that, given fibre bundles $Y$ and $Y^{\prime}$ over $X$, every
bundle morphism $\Phi: Y \rightarrow Y^{\prime}$ over a diffeomorphism $f$ of $X$ admits the $r$-order jet prolongation to a morphism of $r$-order jet manifolds

$$
\begin{equation*}
J^{r} \Phi: J^{r} Y \ni j_{x}^{r} s \rightarrow j_{f(x)}^{r}\left(\Phi \circ s \circ f^{-1}\right) \in J^{r} Y^{\prime} \tag{2.2.6}
\end{equation*}
$$

The jet prolongation functor is exact. If $\Phi$ is an injection or a surjection, so is $J^{r} \Phi$. It also preserves an algebraic structure. In particular, if $Y \rightarrow X$ is a vector bundle, $J^{r} Y \rightarrow X$ is well. If $Y \rightarrow X$ is an affine bundle modelled over the vector bundle $\bar{Y} \rightarrow X$, then $J^{r} Y \rightarrow X$ is an affine bundle modelled over the vector bundle $J^{r} \bar{Y} \rightarrow X$.

Every section $s$ of a fibre bundle $Y \rightarrow X$ admits the $r$-order jet prolongation to the integrable section $\left(J^{r} s\right)(x)=j_{x}^{r} s$ of the jet bundle $J^{r} Y \rightarrow X$.

Let $\mathcal{O}_{k}^{*}=\mathcal{O}^{*}\left(J^{k} Y\right)$ be the DGA of exterior forms on a jet manifold $J^{k} Y$. Every exterior form $\phi$ on a jet manifold $J^{k} Y$ gives rise to the pullback form $\pi_{k}^{k+i *} \phi$ on a jet manifold $J^{k+i} Y$. We have the direct sequence of $C^{\infty}(X)$-algebras

$$
\mathcal{O}^{*}(X) \xrightarrow{\pi^{*}} \mathcal{O}^{*}(Y) \xrightarrow{\pi_{0}^{1 *}} \mathcal{O}_{1}^{*} \xrightarrow{\pi_{1}^{2 *}} \cdots \xrightarrow{\pi_{r-1}^{r}{ }^{*}} \mathcal{O}_{r}^{*}
$$

Remark 2.2.3: By virtue of de Rham Theorem 8.6.4, the cohomology of the de Rham complex of $\mathcal{O}_{k}^{*}$ equals the cohomology $H^{*}\left(J^{k} Y ; \mathbb{R}\right)$ of $J^{k} Y$ with coefficients in the constant sheaf $\mathbb{R}$. The latter in turn coincides with the sheaf cohomology $H^{*}(Y ; \mathbb{R})$ of $Y$ (see Remark 2.2.1) and, thus, it equals the de Rham cohomology $H_{\mathrm{DR}}^{*}(Y)$ of $Y$.

Given a $k$-order jet manifold $J^{k} Y$ of $Y \rightarrow X$, there exists the canonical bundle morphism

$$
r_{(k)}: J^{k} T Y \rightarrow T J^{k} Y
$$

over a surjection

$$
J^{k} Y \underset{X}{\times} J^{k} T X \rightarrow J^{k} Y \underset{X}{\times} T X
$$

whose coordinate expression is

$$
\dot{y}_{\Lambda}^{i} \circ r_{(k)}=\left(\dot{y}^{i}\right)_{\Lambda}-\sum\left(\dot{y}^{i}\right)_{\mu+\Sigma}\left(\dot{x}^{\mu}\right)_{\Xi}, \quad 0 \leq|\Lambda| \leq k
$$

where the sum is taken over all partitions $\Sigma+\Xi=\Lambda$ and $0<|\Xi|$. In particular, we have the canonical isomorphism over $J^{k} Y$

$$
\begin{equation*}
r_{(k)}: J^{k} V Y \rightarrow V J^{k} Y, \quad\left(\dot{y}^{i}\right)_{\Lambda}=\dot{y}_{\Lambda}^{i} \circ r_{(k)} \tag{2.2.7}
\end{equation*}
$$

As a consequence, every projectable vector field $u$ (1.3.1) on a fibre bundle $Y \rightarrow X$ has the following $k$-order jet prolongation to a vector field on $J^{k} Y$ :

$$
\begin{align*}
& J^{k} u=r_{(k)} \circ J^{k} u: J^{k} Y \rightarrow T J^{k} Y \\
& J^{k} u=u^{\lambda} \partial_{\lambda}+u^{i} \partial_{i}+\sum_{0<|\Lambda| \leq k}\left(d_{\Lambda}\left(u^{i}-y_{\mu}^{i} u^{\mu}\right)+y_{\mu+\Lambda}^{i} u^{\mu}\right) \partial_{i}^{\Lambda} \tag{2.2.8}
\end{align*}
$$

(cf. (2.1.8) for $k=1$ ). In particular, the $k$-order jet prolongation (2.2.8) of a vertical vector field $u=u^{i} \partial_{i}$ on $Y \rightarrow X$ is a vertical vector field

$$
\begin{equation*}
J^{k} u=u^{i} \partial_{i}+\sum_{0<\backslash \Lambda \mid \leq k} d_{\Lambda} u^{i} \partial_{i}^{\Lambda} \tag{2.2.9}
\end{equation*}
$$

on $J^{k} Y \rightarrow X$ due to the isomorphism (2.2.7).
A vector field $u_{r}$ on an $r$-order jet manifold $J^{r} Y$ is called projectable if, for any $k<r$, there exists a projectable vector field $u_{k}$ on $J^{k} Y$ such that

$$
u_{k} \circ \pi_{k}^{r}=T \pi_{k}^{r} \circ u_{r}
$$

A projectable vector field $u_{k}$ on $J^{k} Y$ has the coordinate expression

$$
u_{k}=u^{\lambda} \partial_{\lambda}+\sum_{0 \leq|\Lambda| \leq k} u_{\Lambda}^{i} \partial_{i}^{\Lambda}
$$

such that $u_{\lambda}$ depends only on coordinates $x^{\mu}$ and every component $u_{\Lambda}^{i}$ is independent of coordinates $y_{\Xi}^{i},|\Xi|>|\Lambda|$. In particular, the $k$-order jet prolongation $J^{k} u(2.2 .8)$ of a projectable vector field on $Y$ is a projectable vector field on $J^{k} Y$. It is called an integrable vector field.

Let $\mathcal{P}^{k}$ denote a vector space of projectable vector fields on a jet manifold $J^{k} Y$. It is easily seen that $\mathcal{P}^{r}$ is a real Lie algebra and that the morphisms $T \pi_{k}^{r}, k<r$, constitute the inverse system

$$
\begin{equation*}
\mathcal{P}^{0} \stackrel{T \pi_{0}^{1}}{\leftrightarrows} \mathcal{P}^{1} \stackrel{T \pi_{1}^{2}}{\leftrightarrows} \cdots \stackrel{T \pi_{r-2}^{r-1}}{\leftrightarrows} \mathcal{P}^{r-1} \stackrel{T \pi_{r-1}^{r}}{\leftrightarrows} \mathcal{P}^{r} \tag{2.2.10}
\end{equation*}
$$

of these Lie algebras. One can show the following.
Theorem 2.2.1: The $k$-order jet prolongation (2.2.8) is a Lie algebra monomorphism of the Lie algebra $\mathcal{P}^{0}$ of projectable vector fields on $Y \rightarrow$ $X$ to the Lie algebra $\mathcal{P}^{k}$ of projectable vector fields on $J^{k} Y$ such that

$$
\begin{equation*}
T \pi_{k}^{r}\left(J^{r} u\right)=J^{k} u \circ \pi_{k}^{r} \tag{2.2.11}
\end{equation*}
$$

Every projectable vector field $u_{k}$ on $J^{k} Y$ is decomposed into the sum

$$
\begin{equation*}
u_{k}=J^{k}\left(T \pi_{0}^{k}\left(u_{k}\right)\right)+v_{k} \tag{2.2.12}
\end{equation*}
$$

of the integrable vector field $J^{k}\left(T \pi_{0}^{k}\left(u_{k}\right)\right)$ and a projectable vector field $v_{k}$ which is vertical with respect to a fibration $J^{k} Y \rightarrow Y$.

Similarly to the canonical monomorphisms (2.1.5) - (2.1.6), there are the canonical bundle monomorphisms over $J^{k} Y$ :

$$
\begin{align*}
& \lambda_{(k)}: J^{k+1} Y \longrightarrow T^{*} X \underset{J^{k} Y}{\otimes} T J^{k} Y, \\
& \lambda_{(k)}=d x^{\lambda} \otimes d_{\lambda},  \tag{2.2.13}\\
& \theta_{(k)}: J^{k+1} Y \longrightarrow T^{*} J^{k} Y \underset{J^{k} Y}{\otimes} V J^{k} Y, \\
& \theta_{(k)}=\sum_{|\Lambda| \leq k}\left(d y_{\Lambda}^{i}-y_{\lambda+\Lambda}^{i} d x^{\lambda}\right) \otimes \partial_{i}^{\Lambda} . \tag{2.2.14}
\end{align*}
$$

The one-forms

$$
\begin{equation*}
\theta_{\Lambda}^{i}=d y_{\Lambda}^{i}-y_{\lambda+\Lambda}^{i} d x^{\lambda} \tag{2.2.15}
\end{equation*}
$$

are called the local contact forms. Monomorphisms (2.2.13) - (2.2.14) yield the bundle monomorphisms over $J^{k+1} Y$ :

$$
\begin{aligned}
& \hat{\lambda}_{(k)}: T X \underset{X}{\times} J^{k+1} Y \longrightarrow T J^{k} Y \underset{J^{k} Y}{\times} J^{k+1} Y, \\
& \hat{\theta}_{(k)}: V^{*} J^{k} Y \underset{J^{k} Y}{\times} \longrightarrow T^{*} J^{k} Y \underset{J^{k} Y}{\times} J^{k+1} Y
\end{aligned}
$$

(cf. (3.2.1) - (3.2.2) for $k=1$ ). These monomorphisms in turn define the canonical horizontal splittings of the pull-back bundles

$$
\begin{align*}
& \pi_{k}^{k+1 *} T J^{k} Y=\hat{\lambda}_{(k)}\left(T X \underset{X}{\times} J^{k+1} Y\right) \underset{J^{k+1} Y}{\oplus} V J^{k} Y,  \tag{2.2.16}\\
& \dot{x}^{\lambda} \partial_{\lambda}+\sum_{|\Lambda| \leq k} \dot{y}_{\Lambda}^{i} \partial_{i}^{\Lambda}=\dot{x}^{\lambda} d_{\lambda}+\sum_{|\Lambda| \leq k}\left(\dot{y}_{\Lambda}^{i}-\dot{x}^{\lambda} y_{\lambda+\Lambda}^{i}\right) \partial_{i}^{\Lambda}, \\
& \pi_{k}^{k+1 *} T^{*} J^{k} Y=T^{*} X \underset{J^{k+1} Y}{\oplus} \hat{\theta}_{(k)}\left(V^{*} J^{k} Y \underset{J^{k} Y}{\times} J^{k+1} Y\right),  \tag{2.2.17}\\
& \dot{x}_{\lambda} d x^{\lambda}+\sum_{|\Lambda| \leq k} \dot{y}_{i}^{\Lambda} d y_{\Lambda}^{i}=\left(\dot{x}_{\lambda}+\sum_{|\Lambda| \leq k} \dot{y}_{i}^{\Lambda} y_{\lambda+\Lambda}^{i}\right) d x^{\lambda}+\sum \dot{y}_{i}^{\Lambda} \theta_{\Lambda}^{i}
\end{align*}
$$

For instance, it follows from the canonical horizontal splitting (2.2.16) that any vector field $u_{k}$ on $J^{k} Y$ admits the canonical decomposition

$$
\begin{align*}
u_{k}= & u_{H}+u_{V}=\left(u^{\lambda} \partial_{\lambda}+\sum_{|\Lambda| \leq k} y_{\lambda+\Lambda}^{i} \partial_{i}^{\Lambda}\right)+  \tag{2.2.18}\\
& \sum_{|\Lambda| \leq k}\left(u_{\Lambda}^{i}-u^{\lambda} y_{\lambda+\Lambda}^{i}\right) \partial_{i}^{\Lambda}
\end{align*}
$$

over $J^{k+1} Y$ into the horizontal and vertical parts.
By virtue of the canonical horizontal splitting (2.2.17), every exterior one-form $\phi$ on $J^{k} Y$ admits the canonical splitting of its pull-back onto $J^{k+1} Y$ into the horizontal and vertical parts:

$$
\begin{equation*}
\pi_{k}^{k+1 *} \phi=\phi_{H}+\phi_{V}=h_{0} \phi+\left(\phi-h_{0}(\phi)\right) \tag{2.2.19}
\end{equation*}
$$

where $h_{0}$ is the horizontal projection

$$
h_{0}\left(d x^{\lambda}\right)=d x^{\lambda}, \quad h_{0}\left(d y_{\lambda_{1} \cdots \lambda_{k}}^{i}\right)=y_{\mu \lambda_{1} \ldots \lambda_{k}}^{i} d x^{\mu}
$$

The vertical part of the splitting is called a contact one-form on $J^{k+1} Y$.
Let us consider an ideal of the algebra $\mathcal{O}_{k}^{*}$ of exterior forms on $J^{k} Y$ which is generated by the contact one-forms on $J^{k} Y$. This ideal, called the ideal of contact forms, is locally generated by the contact forms $\theta_{\Lambda}^{i}$ (2.2.15). One can show that an exterior form $\phi$ on the a manifold $J^{k} Y$ is a contact form iff its pull-back $\bar{s}^{*} \phi$ onto a base $X$ by means of any integrable section $\bar{s}$ of $J^{k} Y \rightarrow X$ vanishes.

### 2.3 Differential operators and equations

Jet manifolds provides the conventional language of theory of differential equations and differential operators if they need not be linear.

Definition 2.3.1: A system of $k$-order partial differential equations on a fibre bundle $Y \rightarrow X$ is defined as a closed subbundle $\mathfrak{E}$ of a jet bundle $J^{k} Y \rightarrow X$. For the sake of brevity, we agree to call $\mathfrak{E}$ a differential equation.

By a classical solution of a differential equation $\mathfrak{E}$ on $Y \rightarrow X$ is meant a section $s$ of $Y \rightarrow X$ such that its $k$-order jet prolongation $J^{k} s$ lives in $\mathfrak{E}$.

Let $J^{k} Y$ be provided with the adapted coordinates $\left(x^{\lambda}, y_{\Lambda}^{i}\right)$. There exists a local coordinate system $\left(z^{A}\right), A=1, \ldots$, codime, on $J^{k} Y$ such that $\mathfrak{E}$ is locally given (in the sense of item (i) of Theorem 1.1.1) by equations

$$
\begin{equation*}
\mathcal{E}^{A}\left(x^{\lambda}, y_{\Lambda}^{i}\right)=0, \quad A=1, \ldots, \text { codime } . \tag{2.3.1}
\end{equation*}
$$

Differential equations are often associated to differential operators. There are several equivalent definitions of differential operators.

Definition 2.3.2: Let $Y \rightarrow X$ and $E \rightarrow X$ be fibre bundles, which are assumed to have global sections. A $k$-order $E$-valued differential operator on a fibre bundle $Y \rightarrow X$ is defined as a section $\mathcal{E}$ of the pull-back bundle

$$
\begin{equation*}
\operatorname{pr}_{1}: E_{Y}^{k}=J^{k} \underset{X}{\times} E \rightarrow J^{k} Y \tag{2.3.2}
\end{equation*}
$$

Given bundle coordinates $\left(x^{\lambda}, y^{i}\right)$ on $Y$ and $\left(x^{\lambda}, \chi^{a}\right)$ on $E$, the pullback (2.3.2) is provided with coordinates $\left(x^{\lambda}, y_{\Sigma}^{j}, \chi^{a}\right), 0 \leq|\Sigma| \leq k$. With respect to these coordinates, a differential operator $\mathcal{E}$ seen as a closed imbedded submanifold $\mathcal{E} \subset E_{Y}^{k}$ is given by the equalities

$$
\begin{equation*}
\chi^{a}=\mathcal{E}^{a}\left(x^{\lambda}, y_{\Sigma}^{j}\right) \tag{2.3.3}
\end{equation*}
$$

There is obvious one-to-one correspondence between the sections $\mathcal{E}$ (2.3.3) of the fibre bundle (2.3.2) and the bundle morphisms

$$
\begin{align*}
& \Phi: J^{k} Y \underset{X}{\longrightarrow} E  \tag{2.3.4}\\
& \Phi=\operatorname{pr}_{2} \circ \mathcal{E} \Longleftrightarrow \mathcal{E}=\left(\operatorname{Id} J^{k} Y, \Phi\right) .
\end{align*}
$$

Therefore, we come to the following equivalent definition of differential operators on $Y \rightarrow X$.

Definition 2.3.3: Let $Y \rightarrow X$ and $E \rightarrow X$ be fibre bundles. A bundle morphism $J^{k} Y \rightarrow E$ over $X$ is called a $E$-valued $k$-order differential operator on $Y \rightarrow X$.

It is readily observed that the differential operator $\Phi$ (2.3.4) sends each section $s$ of $Y \rightarrow X$ onto the section $\Phi \circ J^{k} s$ of $E \rightarrow X$. The mapping

$$
\begin{aligned}
& \Delta_{\Phi}: \mathcal{S}(Y) \rightarrow \mathcal{S}(E), \\
& \Delta_{\Phi}: s \rightarrow \Phi \circ J^{k} s, \quad \chi^{a}(x)=\mathcal{E}^{a}\left(x^{\lambda}, \partial_{\Sigma} s^{j}(x)\right),
\end{aligned}
$$

is called the standard form of a differential operator.
Let $e$ be a global section of a fibre bundle $E \rightarrow X$, the kernel of a $E$-valued differential operator $\Phi$ is defined as the kernel

$$
\begin{equation*}
\operatorname{Ker}_{e} \Phi=\Phi^{-1}(e(X)) \tag{2.3.5}
\end{equation*}
$$

of the bundle morphism $\Phi$ (2.3.4). If it is a closed subbundle of the jet bundle $J^{k} Y \rightarrow X$, one says that $\operatorname{Ker}_{e} \Phi(2.3 .5)$ is a differential equation associated to the differential operator $\Phi$. By virtue of Theorem 1.1.10, this condition holds if $\Phi$ is a bundle morphism of constant rank.

If $E \rightarrow X$ is a vector bundle, by the kernel of a $E$-valued differential operator is usually meant its kernel with respect to the canonical zerovalued section $\hat{0}$ of $E \rightarrow X$.

In the framework of Lagrangian formalism, we deal with differential operators of the following type. Let

$$
F \rightarrow Y \rightarrow X, \quad E \rightarrow Y \rightarrow X
$$

be composite bundles where $E \rightarrow Y$ is a vector bundle. By a $k$-order differential operator on $F \rightarrow X$ taking its values into $E \rightarrow X$ is meant a bundle morphism

$$
\begin{equation*}
\Phi: J^{k} F \underset{Y}{\longrightarrow} E \tag{2.3.6}
\end{equation*}
$$

which certainly is a bundle morphism over $X$ in accordance with Definition 2.3.3. Its kernel $\operatorname{Ker} \Phi$ is defined as the inverse image of the canonical zero-valued section of $E \rightarrow Y$. In an equivalent way, the differential operator (2.3.6) is represented by a section $\mathcal{E}_{\Phi}$ of the vector bundle

$$
J^{k} F \underset{Y}{\times} E \rightarrow J^{k} F
$$

Given bundle coordinates $\left(x^{\lambda}, y^{i}, w^{r}\right)$ on $F$ and $\left(x^{\lambda}, y^{i}, c^{A}\right)$ on $E$ with respect to the fibre basis $\left\{e_{A}\right\}$ for $E \rightarrow Y$, this section reads

$$
\begin{equation*}
\mathcal{E}_{\Phi}=\mathcal{E}^{A}\left(x^{\lambda}, y_{\Lambda}^{i}, w_{\Lambda}^{r}\right) e_{A}, \quad 0 \leq|\Lambda| \leq k \tag{2.3.7}
\end{equation*}
$$

Then the differential operator (2.3.6) also is represented by a function

$$
\begin{equation*}
\mathcal{E}_{\Phi}=\mathcal{E}^{A}\left(x^{\lambda}, y_{\Lambda}^{i}, w_{\Lambda}^{r}\right) c_{A} \in C^{\infty}\left(F \underset{Y}{\times} E^{*}\right) \tag{2.3.8}
\end{equation*}
$$

on the product $F \times_{Y} E^{*}$, where $E^{*} \rightarrow Y$ is the dual of $E \rightarrow Y$ coordinated by $\left(x^{\lambda}, y^{i}, c_{A}\right)$.

If $F \rightarrow Y$ is a vector bundle, a differential operator $\Phi$ (2.3.6) on the composite bundle $F \rightarrow Y \rightarrow X$ is called linear if it is linear on the fibres of the vector bundle $J^{k} F \rightarrow J^{k} Y$. In this case, its representations (2.3.7) and (2.3.8) take the form

$$
\begin{array}{ll}
\mathcal{E}_{\Phi}=\sum_{0 \leq|\Xi| \leq k} \mathcal{E}_{r}^{A, \Xi}\left(x^{\lambda}, y_{\Lambda}^{i}\right) w_{\Xi}^{r} e_{A}, & 0 \leq|\Lambda| \leq k, \\
\mathcal{E}_{\Phi}=\sum_{0 \leq|\Xi| \leq k} \mathcal{E}_{r}^{A, \Xi}\left(x^{\lambda}, y_{\Lambda}^{i}\right) w_{\Xi}^{r} c_{A}, & 0 \leq|\Lambda| \leq k \tag{2.3.10}
\end{array}
$$

### 2.4 Infinite order jet formalism

The finite order jet manifolds $J^{k} Y$ of a fibre bundle $Y \rightarrow X$ form the inverse sequence

$$
\begin{equation*}
Y \stackrel{\pi}{\longleftarrow} J^{1} Y \longleftarrow \cdots J^{r-1} Y \stackrel{\pi_{r-1}^{r}}{\longleftarrow} J^{r} Y \longleftarrow \cdots \tag{2.4.1}
\end{equation*}
$$

where $\pi_{r-1}^{r}$ are affine bundles modelled over the vector bundles (2.2.4). Its inductive limit $J^{\infty} Y$ is defined as a minimal set such that there exist surjections

$$
\begin{equation*}
\pi^{\infty}: J^{\infty} Y \rightarrow X, \quad \pi_{0}^{\infty}: J^{\infty} Y \rightarrow Y, \quad \pi_{k}^{\infty}: J^{\infty} Y \rightarrow J^{k} Y \tag{2.4.2}
\end{equation*}
$$

obeying the relations $\pi_{r}^{\infty}=\pi_{r}^{k} \circ \pi_{k}^{\infty}$ for all admissible $k$ and $r<k$. A projective limit of the inverse system (2.4.1) always exists. It consists of those elements

$$
\left(\ldots, z_{r}, \ldots, z_{k}, \ldots\right), \quad z_{r} \in J^{r} Y, \quad z_{k} \in J^{k} Y
$$

of the Cartesian product $\prod_{k} J^{k} Y$ which satisfy the relations $z_{r}=\pi_{r}^{k}\left(z_{k}\right)$ for all $k>r$. One can think of elements of $J^{\infty} Y$ as being infinite order jets of sections of $Y \rightarrow X$ identified by their Taylor series at points of $X$.

The set $J^{\infty} Y$ is provided with the projective limit topology. This is the coarsest topology such that the surjections $\pi_{r}^{\infty}(2.4 .2)$ are continuous. Its base consists of inverse images of open subsets of $J^{r} Y, r=0, \ldots$, under the maps $\pi_{r}^{\infty}$. With this topology, $J^{\infty} Y$ is a paracompact Fréchet (complete metrizable) manifold modelled over a locally convex vector space of number series $\left\{a^{\lambda}, a^{i}, a_{\lambda}^{i}, \cdots\right\}$. It is called the infinite order jet manifold. One can show that the surjections $\pi_{r}^{\infty}$ are open maps admitting local sections, i.e., $J^{\infty} Y \rightarrow J^{r} Y$ are continuous bundles. A bundle coordinate atlas $\left\{U_{Y},\left(x^{\lambda}, y^{i}\right)\right\}$ of $Y \rightarrow X$ provides $J^{\infty} Y$ with the manifold coordinate atlas

$$
\begin{equation*}
\left\{\left(\pi_{0}^{\infty}\right)^{-1}\left(U_{Y}\right),\left(x^{\lambda}, y_{\Lambda}^{i}\right)\right\}_{0 \leq|\Lambda|}, \quad y_{\lambda+\Lambda}^{\prime i}=\frac{\partial x^{\mu}}{\partial x^{\prime \lambda}} d_{\mu} y_{\Lambda}^{\prime i} \tag{2.4.3}
\end{equation*}
$$

Theorem 2.4.1: A fibre bundle $Y$ is a strong deformation retract of the infinite order jet manifold $J^{\infty} Y$.

Corollary 2.4.2: There is an isomorphism

$$
\begin{equation*}
H^{*}\left(J^{\infty} Y ; \mathbb{R}\right)=H^{*}(Y ; \mathbb{R}) \tag{2.4.4}
\end{equation*}
$$

between cohomology of $J^{\infty} Y$ and $Y$ with coefficients in the sheaf $\mathbb{R}$.
The inverse sequence (2.4.1) of jet manifolds yields the direct sequence of DGAs $\mathcal{O}_{r}^{*}$ of exterior forms on finite order jet manifolds

$$
\begin{equation*}
\mathcal{O}^{*}(X) \xrightarrow{\pi^{*}} \mathcal{O}^{*}(Y) \xrightarrow{\pi_{0}^{1 *}} \mathcal{O}_{1}^{*} \longrightarrow \cdots \mathcal{O}_{r-1}^{*} \xrightarrow{\pi_{r-1}^{r}{ }^{*}} \mathcal{O}_{r}^{*} \longrightarrow \cdots, \tag{2.4.5}
\end{equation*}
$$

where $\pi_{r-1}^{r}{ }^{*}$ are the pull-back monomorphisms. Its direct limit

$$
\begin{equation*}
\mathcal{O}_{\infty}^{*} Y=\overrightarrow{\lim } \mathcal{O}_{r}^{*} \tag{2.4.6}
\end{equation*}
$$

exists and consists of all exterior forms on finite order jet manifolds modulo the pull-back identification. In accordance with Theorem 8.1.5, $\mathcal{O}_{\infty}^{*} Y$ is a DGA which inherits the operations of the exterior differential $d$ and exterior product $\wedge$ of exterior algebras $\mathcal{O}_{r}^{*}$. If there is no danger of confusion, we denote $\mathcal{O}_{\infty}^{*}=\mathcal{O}_{\infty}^{*} Y$.

Theorem 2.4.3: The cohomology $H^{*}\left(\mathcal{O}_{\infty}^{*}\right)$ of the de Rham complex

$$
\begin{equation*}
0 \longrightarrow \mathbb{R} \longrightarrow \mathcal{O}_{\infty}^{0} \xrightarrow{d} \mathcal{O}_{\infty}^{1} \xrightarrow{d} \cdots \tag{2.4.7}
\end{equation*}
$$

of the DGA $\mathcal{O}_{\infty}^{*}$ equals the de Rham cohomology $H_{\mathrm{DR}}^{*}(Y)$ of $Y$.
Corollary 2.4.4: Any closed form $\phi \in \mathcal{O}_{\infty}^{*}$ is decomposed into the sum $\phi=\sigma+d \xi$, where $\sigma$ is a closed form on $Y$.

One can think of elements of $\mathcal{O}_{\infty}^{*}$ as being differential forms on the infinite order jet manifold $J^{\infty} Y$ as follows. Let $\mathfrak{O}_{r}^{*}$ be a sheaf of germs of exterior forms on $J^{r} Y$ and $\overline{\mathfrak{D}}_{r}^{*}$ the canonical presheaf of local sections of $\mathfrak{O}_{r}^{*}$. Since $\pi_{r-1}^{r}$ are open maps, there is the direct sequence of presheaves

$$
\overline{\mathfrak{O}}_{0}^{*} \xrightarrow{\pi_{0}^{1 *}} \overline{\mathfrak{O}}_{1}^{*} \cdots \xrightarrow{\pi_{r-1}^{r}} \overline{\mathfrak{D}}_{r}^{*} \longrightarrow \cdots
$$

Its direct limit $\overline{\mathfrak{D}}_{\infty}^{*}$ is a presheaf of DGAs on $J^{\infty} Y$. Let $\mathfrak{Q}_{\infty}^{*}$ be the sheaf of DGAs of germs of $\overline{\mathfrak{D}}_{\infty}^{*}$ on $J^{\infty} Y$. The structure module

$$
\begin{equation*}
\mathcal{Q}_{\infty}^{*}=\Gamma\left(\mathfrak{Q}_{\infty}^{*}\right) \tag{2.4.8}
\end{equation*}
$$

of global sections of $\mathfrak{Q}_{\infty}^{*}$ is a DGA such that, given an element $\phi \in \mathcal{Q}_{\infty}^{*}$ and a point $z \in J^{\infty} Y$, there exist an open neighbourhood $U$ of $z$ and an exterior form $\phi^{(k)}$ on some finite order jet manifold $J^{k} Y$ so that

$$
\left.\phi\right|_{U}=\left.\pi_{k}^{\infty *} \phi^{(k)}\right|_{U}
$$

Therefore, one can regard $\mathcal{Q}_{\infty}^{*}$ as an algebra of locally exterior forms on finite order jet manifolds. There is a monomorphism $\mathcal{O}_{\infty}^{*} \rightarrow \mathcal{Q}_{\infty}^{*}$.

Theorem 2.4.5: The paracompact space $J^{\infty} Y$ admits a partition of unity by elements of the ring $\mathcal{Q}_{\infty}^{0}$.

Since elements of the DGA $\mathcal{Q}_{\infty}^{*}$ are locally exterior forms on finite order jet manifolds, the following Poincaré lemma holds.

Lemma 2.4.6: Given a closed element $\phi \in \mathcal{Q}_{\infty}^{*}$, there exists a neighbourhood $U$ of each point $z \in J^{\infty} Y$ such that $\left.\phi\right|_{U}$ is exact.

Theorem 2.4.7: The cohomology $H^{*}\left(\mathcal{Q}_{\infty}^{*}\right)$ of the de Rham complex

$$
\begin{equation*}
0 \longrightarrow \mathbb{R} \longrightarrow \mathcal{Q}_{\infty}^{0} \xrightarrow{d} \mathcal{Q}_{\infty}^{1} \xrightarrow{d} \cdots \tag{2.4.9}
\end{equation*}
$$

of the DGA $\mathcal{Q}_{\infty}^{*}$ equals the de Rham cohomology of a fibre bundle $Y$.
Due to a monomorphism $\mathcal{O}_{\infty}^{*} \rightarrow \mathcal{Q}_{\infty}^{*}$, one can restrict $\mathcal{O}_{\infty}^{*}$ to the coordinate chart (2.4.3) where horizontal forms $d x^{\lambda}$ and contact oneforms

$$
\theta_{\Lambda}^{i}=d y_{\Lambda}^{i}-y_{\lambda+\Lambda}^{i} d x^{\lambda}
$$

make up a local basis for the $\mathcal{O}_{\infty}^{0}$-algebra $\mathcal{O}_{\infty}^{*}$. Though $J^{\infty} Y$ is not a smooth manifold, elements of $\mathcal{O}_{\infty}^{*}$ are exterior forms on finite order jet manifolds and, therefore, their coordinate transformations are smooth. Moreover, there is the canonical decomposition

$$
\mathcal{O}_{\infty}^{*}=\oplus \mathcal{O}_{\infty}^{k, m}
$$

of $\mathcal{O}_{\infty}^{*}$ into $\mathcal{O}_{\infty}^{0}$-modules $\mathcal{O}_{\infty}^{k, m}$ of $k$-contact and $m$-horizontal forms together with the corresponding projectors

$$
h_{k}: \mathcal{O}_{\infty}^{*} \rightarrow \mathcal{O}_{\infty}^{k, *}, \quad h^{m}: \mathcal{O}_{\infty}^{*} \rightarrow \mathcal{O}_{\infty}^{*, m}
$$

Accordingly, the exterior differential on $\mathcal{O}_{\infty}^{*}$ is decomposed into the sum

$$
d=d_{V}+d_{H}
$$

of the vertical differential

$$
d_{V} \circ h^{m}=h^{m} \circ d \circ h^{m}, \quad d_{V}(\phi)=\theta_{\Lambda}^{i} \wedge \partial_{i}^{\Lambda} \phi, \quad \phi \in \mathcal{O}_{\infty}^{*}
$$

and the total differential

$$
d_{H} \circ h_{k}=h_{k} \circ d \circ h_{k}, \quad d_{H} \circ h_{0}=h_{0} \circ d, \quad d_{H}(\phi)=d x^{\lambda} \wedge d_{\lambda}(\phi)
$$

where

$$
\begin{equation*}
d_{\lambda}=\partial_{\lambda}+y_{\lambda}^{i} \partial_{i}+\sum_{0<|\Lambda|} y_{\lambda+\Lambda}^{i} \partial_{i}^{\Lambda} \tag{2.4.10}
\end{equation*}
$$

are the infinite order total derivatives. They obey the nilpotent conditions

$$
\begin{equation*}
d_{H} \circ d_{H}=0, \quad d_{V} \circ d_{V}=0, \quad d_{H} \circ d_{V}+d_{V} \circ d_{H}=0 \tag{2.4.11}
\end{equation*}
$$

and make $\mathcal{O}_{\infty}^{*, *}$ into a bicomplex.
Let us consider the $\mathcal{O}_{\infty}^{0}$-module $\mathfrak{o} \mathcal{O}_{\infty}^{0}$ of derivations of the real ring $\mathcal{O}_{\infty}^{0}$.

Theorem 2.4.8: The derivation module $\mathfrak{d} \mathcal{O}_{\infty}^{0}$ is isomorphic to the $\mathcal{O}_{\infty^{-}}^{0}$ dual $\left(\mathcal{O}_{\infty}^{1}\right)^{*}$ of the module of one-forms $\mathcal{O}_{\infty}^{1}$.

One can say something more. The DGA $\mathcal{O}_{\infty}^{*}$ is a minimal ChevalleyEilenberg differential calculus $\mathcal{O}^{*} \mathcal{A}$ over the real ring $\mathcal{A}=\mathcal{O}_{\infty}^{0}$ of smooth real functions on finite order jet manifolds of $Y \rightarrow X$. Let $\vartheta\rfloor \phi, \vartheta \in \mathfrak{d} \mathcal{O}_{\infty}^{0}$, $\phi \in \mathcal{O}_{\infty}^{1}$, denote the interior product. Extended to the DGA $\mathcal{O}_{\infty}^{*}$, the interior product $\rfloor$ obeys the rule

$$
\left.\vartheta\rfloor(\phi \wedge \sigma)=(\vartheta\rfloor \phi) \wedge \sigma+(-1)^{|\phi|} \phi \wedge(\vartheta\rfloor \sigma\right) .
$$

Restricted to a coordinate chart (2.4.3), $\mathcal{O}_{\infty}^{1}$ is a free $\mathcal{O}_{\infty}^{0}$-module generated by one-forms $d x^{\lambda}, \theta_{\Lambda}^{i}$. Since $\mathfrak{d} \mathcal{O}_{\infty}^{0}=\left(\mathcal{O}_{\infty}^{1}\right)^{*}$, any derivation of the real ring $\mathcal{O}_{\infty}^{0}$ takes the coordinate form

$$
\begin{align*}
& \vartheta=\vartheta^{\lambda} \partial_{\lambda}+\vartheta^{i} \partial_{i}+\sum_{0<|\Lambda|} \vartheta_{\Lambda}^{i} \partial_{i}^{\Lambda},  \tag{2.4.12}\\
& \left.\partial_{i}^{\Lambda}\left(y_{\Sigma}^{j}\right)=\partial_{i}^{\Lambda}\right\rfloor d y_{\Sigma}^{j}=\delta_{i}^{j} \delta_{\Sigma}^{\Lambda}, \\
& \vartheta^{\prime \lambda}=\frac{\partial x^{\prime \lambda}}{\partial x^{\mu}} \vartheta^{\mu}, \quad \vartheta^{\prime i}=\frac{\partial y^{\prime i}}{\partial y^{j}} \vartheta^{j}+\frac{\partial y^{\prime i}}{\partial x^{\mu}} \vartheta^{\mu}, \\
& \vartheta_{\Lambda}^{\prime i}=\sum_{|\Sigma| \leq|\Lambda|} \frac{\partial y_{\Lambda}^{\prime i}}{\partial y_{\Sigma}^{j}} \vartheta_{\Sigma}^{j}+\frac{\partial y_{\Lambda}^{\prime i}}{\partial x^{\mu}} \vartheta^{\mu} . \tag{2.4.13}
\end{align*}
$$

Any derivation $\vartheta$ (2.4.12) of the ring $\mathcal{O}_{\infty}^{0}$ yields a derivation (called the Lie derivative) $\mathbf{L}_{\vartheta}$ of the DGA $\mathcal{O}_{\infty}^{*}$ given by the relations

$$
\begin{aligned}
& \left.\left.\mathbf{L}_{\vartheta} \phi=\vartheta\right\rfloor d \phi+d(\vartheta\rfloor \phi\right) \\
& \mathbf{L}_{\vartheta}\left(\phi \wedge \phi^{\prime}\right)=\mathbf{L}_{\vartheta}(\phi) \wedge \phi^{\prime}+\phi \wedge \mathbf{L}_{\vartheta}\left(\phi^{\prime}\right)
\end{aligned}
$$

Remark 2.4.1: In particular, the total derivatives (2.4.10) are defined as the local derivations of $\mathcal{O}_{\infty}^{0}$ and the corresponding Lie derivatives

$$
d_{\lambda} \phi=\mathbf{L}_{d_{\lambda}} \phi
$$

of $\mathcal{O}_{\infty}^{*}$. Moreover, the $C^{\infty}(X)$-ring $\mathcal{O}_{\infty}^{0}$ possesses the canonical connection

$$
\begin{equation*}
\nabla=d x^{\lambda} \otimes d_{\lambda} \tag{2.4.14}
\end{equation*}
$$

in the sense of Definition 8.2.4.

## Chapter 3

## Connections on fibre bundles

There are several equivalent definitions of a connection on a fibre bundle. We start with the traditional notion of a connection as a splitting of the exact sequences (1.2.10) - (1.2.11), but then follow its definition as a global section of an affine jet bundle. In the case of vector bundles, there is an equivalent definition (8.6.3) of a linear connection on their structure modules.

### 3.1 Connections as tangent-valued forms

A connection on a fibre bundle $Y \rightarrow X$ is defined traditionally as a linear bundle monomorphism

$$
\begin{align*}
& \Gamma: Y \times T X \rightarrow T Y,  \tag{3.1.1}\\
& \Gamma: \dot{x}^{\lambda} \partial_{\lambda} \rightarrow \dot{x}^{\lambda}\left(\partial_{\lambda}+\Gamma_{\lambda}^{i} \partial_{i}\right),
\end{align*}
$$

over $Y$ which splits the exact sequence (1.2.10), i.e.,

$$
\pi_{T} \circ \Gamma=\operatorname{Id}(Y \underset{X}{\times} T X) .
$$

This is a definition of connections on fibred manifolds (see Remark 3.1.2). By virtue of Theorem 1.2.2, a connection always exists. The local functions $\Gamma_{\lambda}^{i}(y)$ in (3.1.1) are said to be components of the connection $\Gamma$ with respect to the bundle coordinates $\left(x^{\lambda}, y^{i}\right)$ on $Y \rightarrow X$.

The image of $Y \times T X$ by the connection $\Gamma$ defines the horizontal distribution $H Y \subset T Y$ which splits the tangent bundle $T Y$ as follows:

$$
\begin{align*}
& T Y=H Y \underset{Y}{\oplus} V Y  \tag{3.1.2}\\
& \dot{x}^{\lambda} \partial_{\lambda}+\dot{y}^{i} \partial_{i}=\dot{x}^{\lambda}\left(\partial_{\lambda}+\Gamma_{\lambda}^{i} \partial_{i}\right)+\left(\dot{y}^{i}-\dot{x}^{\lambda} \Gamma_{\lambda}^{i}\right) \partial_{i}
\end{align*}
$$

Its annihilator is locally generated by the one-forms $d y^{i}-\Gamma_{\lambda}^{i} d x^{\lambda}$.
Given the horizontal splitting (3.1.2), the surjection

$$
\begin{equation*}
\Gamma: T Y \underset{Y}{\rightarrow} V Y, \quad \dot{y}^{i} \circ \Gamma=\dot{y}^{i}-\Gamma_{\lambda}^{i} \dot{x}^{\lambda} \tag{3.1.3}
\end{equation*}
$$

defines a connection on $Y \rightarrow X$ in an equivalent way.
The linear morphism $\Gamma$ over $Y$ (3.1.1) yields uniquely the horizontal tangent-valued one-form

$$
\begin{equation*}
\Gamma=d x^{\lambda} \otimes\left(\partial_{\lambda}+\Gamma_{\lambda}^{i} \partial_{i}\right) \tag{3.1.4}
\end{equation*}
$$

on $Y$ which projects onto the canonical tangent-valued form $\theta_{X}$ (1.4.5) on $X$. With this form called the connection form, the morphism (3.1.1) reads

$$
\left.\Gamma: \partial_{\lambda} \rightarrow \partial_{\lambda}\right\rfloor \Gamma=\partial_{\lambda}+\Gamma_{\lambda}^{i} \partial_{i}
$$

Given a connection $\Gamma$ and the corresponding horizontal distribution (3.1.2), a vector field $u$ on a fibre bundle $Y \rightarrow X$ is called horizontal if it lives in $H Y$. A horizontal vector field takes the form

$$
\begin{equation*}
u=u^{\lambda}(y)\left(\partial_{\lambda}+\Gamma_{\lambda}^{i} \partial_{i}\right) \tag{3.1.5}
\end{equation*}
$$

In particular, let $\tau$ be a vector field on the base $X$. By means of the connection form $\Gamma$ (3.1.4), we obtain the projectable horizontal vector field

$$
\begin{equation*}
\Gamma \tau=\tau\rfloor \Gamma=\tau^{\lambda}\left(\partial_{\lambda}+\Gamma_{\lambda}^{i} \partial_{i}\right) \tag{3.1.6}
\end{equation*}
$$

on $Y$, called the horizontal lift of $\tau$ by means of a connection $\Gamma$. Conversely, any projectable horizontal vector field $u$ on $Y$ is the horizontal
lift $\Gamma \tau$ of its projection $\tau$ on $X$. Moreover, the horizontal distribution $H Y$ is generated by the horizontal lifts $\Gamma \tau$ (3.1.6) of vector fields $\tau$ on $X$. The horizontal lift

$$
\begin{equation*}
\mathcal{T}(X) \ni \tau \rightarrow \Gamma \tau \in \mathcal{T}(Y) \tag{3.1.7}
\end{equation*}
$$

is a $C^{\infty}(X)$-linear module morphism.
Given the splitting (3.1.1), the dual splitting of the exact sequence (1.2.11) is

$$
\begin{equation*}
\Gamma: V^{*} Y \rightarrow T^{*} Y, \quad \Gamma: \bar{d} y^{i} \rightarrow d y^{i}-\Gamma_{\lambda}^{i} d x^{\lambda} . \tag{3.1.8}
\end{equation*}
$$

Hence, a connection $\Gamma$ on $Y \rightarrow X$ is represented by the vertical-valued form

$$
\begin{equation*}
\Gamma=\left(d y^{i}-\Gamma_{\lambda}^{i} d x^{\lambda}\right) \otimes \partial_{i} \tag{3.1.9}
\end{equation*}
$$

such that the morphism (3.1.8) reads

$$
\left.\Gamma: \bar{d} y^{i} \rightarrow \Gamma\right\rfloor \bar{d} y^{i}=d y^{i}-\Gamma_{\lambda}^{i} d x^{\lambda}
$$

We call $\Gamma$ (3.1.9) the vertical connection form. The corresponding horizontal splitting of the cotangent bundle $T^{*} Y$ takes the form

$$
\begin{align*}
& T^{*} Y=T^{*} X \underset{Y}{\oplus} \Gamma\left(V^{*} Y\right),  \tag{3.1.10}\\
& \dot{x}_{\lambda} d x^{\lambda}+\dot{y}_{i} d y^{i}=\left(\dot{x}_{\lambda}+\dot{y}_{i} \Gamma_{\lambda}^{i}\right) d x^{\lambda}+\dot{y}_{i}\left(d y^{i}-\Gamma_{\lambda}^{i} d x^{\lambda}\right) .
\end{align*}
$$

Then we have the surjection

$$
\begin{equation*}
\Gamma=\operatorname{pr}_{1}: T^{*} Y \rightarrow T^{*} X, \quad \dot{x}_{\lambda} \circ \Gamma=\dot{x}_{\lambda}+\dot{y}_{i} \Gamma_{\lambda}^{i} \tag{3.1.11}
\end{equation*}
$$

which also defines a connection on a fibre bundle $Y \rightarrow X$.
Remark 3.1.1: Treating a connection as the vertical-valued form (3.1.9), we come to the following important construction. Given a fibre bundle $Y \rightarrow X$, let $f: X^{\prime} \rightarrow X$ be a morphism and $f^{*} Y \rightarrow X^{\prime}$ the pull-back of $Y$ by $f$. Any connection $\Gamma$ (3.1.9) on $Y \rightarrow X$ induces the pull-back connection

$$
\begin{equation*}
f^{*} \Gamma=\left(d y^{i}-\left(\Gamma \circ f_{Y}\right)_{\lambda}^{i} \frac{\partial f^{\lambda}}{\partial x^{\prime \mu}} d x^{\prime \mu}\right) \otimes \partial_{i} \tag{3.1.12}
\end{equation*}
$$

on $f^{*} Y \rightarrow X^{\prime}$ (see Remark 1.4.4).
Remark 3.1.2: Let $\pi: Y \rightarrow X$ be a fibred manifold. Any connection $\Gamma$ on $Y \rightarrow X$ yields a horizontal lift of a vector field on $X$ onto $Y$, but need not defines the similar lift of a path in $X$ into $Y$. Let

$$
\mathbb{R} \supset[,] \ni t \rightarrow x(t) \in X, \quad \mathbb{R} \ni t \rightarrow y(t) \in Y
$$

be smooth paths in $X$ and $Y$, respectively. Then $t \rightarrow y(t)$ is called a horizontal lift of $x(t)$ if

$$
\pi(y(t))=x(t), \quad \dot{y}(t) \in H_{y(t)} Y, \quad t \in \mathbb{R}
$$

where $H Y \subset T Y$ is the horizontal subbundle associated to the connection $\Gamma$. If, for each path $x(t)\left(t_{0} \leq t \leq t_{1}\right)$ and for any $y_{0} \in \pi^{-1}\left(x\left(t_{0}\right)\right)$, there exists a horizontal lift $y(t)\left(t_{0} \leq t \leq t_{1}\right)$ such that $y\left(t_{0}\right)=y_{0}$, then $\Gamma$ is called the Ehresmann connection. A fibred manifold is a fibre bundle iff it admits an Ehresmann connection.

### 3.2 Connections as jet bundle sections

Throughout the book, we follow the equivalent definition of connections on a fibre bundle $Y \rightarrow X$ as sections of the affine jet bundle $J^{1} Y \rightarrow Y$.

Let $Y \rightarrow X$ be a fibre bundle, and $J^{1} Y$ its first order jet manifold. Given the canonical morphisms (2.1.5) and (2.1.6), we have the corresponding morphisms

$$
\begin{align*}
& \left.\hat{\lambda}_{(1)}: J^{1} Y \underset{X}{\times} T X \ni \partial_{\lambda} \rightarrow d_{\lambda}=\partial_{\lambda}\right\rfloor \lambda_{(1)} \in J^{1} Y \underset{Y}{\times} T Y,  \tag{3.2.1}\\
& \left.\hat{\theta}_{(1)}: J^{1} Y \underset{Y}{\times} V^{*} Y \ni \bar{d} y^{i} \rightarrow \theta^{i}=\theta_{(1)}\right\rfloor d y^{i} \in J^{1} \underset{Y}{\times} T^{*} Y \tag{3.2.2}
\end{align*}
$$

(see Remark 1.2.2). These morphisms yield the canonical horizontal splittings of the pull-back bundles

$$
\begin{align*}
& J^{1} Y \underset{Y}{\times} T Y=\widehat{\lambda}_{(1)}(T X) \underset{J^{1} Y}{\oplus} V Y,  \tag{3.2.3}\\
& \dot{x}^{\lambda} \partial_{\lambda}+\dot{y}^{i} \partial_{i}=\dot{x}^{\lambda}\left(\partial_{\lambda}+y_{\lambda}^{i} \partial_{i}\right)+\left(\dot{y}^{i}-\dot{x}^{\lambda} y_{\lambda}^{i}\right) \partial_{i},
\end{align*}
$$

$$
\begin{align*}
& J^{1} Y \underset{Y}{\times} T^{*} Y=T^{*} X \underset{J^{1} Y}{\oplus} \hat{\theta}_{(1)}\left(V^{*} Y\right)  \tag{3.2.4}\\
& \dot{x}_{\lambda} d x^{\lambda}+\dot{y}_{i} d y^{i}=\left(\dot{x}_{\lambda}+\dot{y}_{i} y_{\lambda}^{i}\right) d x^{\lambda}+\dot{y}_{i}\left(d y^{i}-y_{\lambda}^{i} d x^{\lambda}\right) .
\end{align*}
$$

Let $\Gamma$ be a global section of $J^{1} Y \rightarrow Y$. Substituting the tangent-valued form

$$
\lambda_{(1)} \circ \Gamma=d x^{\lambda} \otimes\left(\partial_{\lambda}+\Gamma_{\lambda}^{i} \partial_{i}\right)
$$

in the canonical splitting (3.2.3), we obtain the familiar horizontal splitting (3.1.2) of $T Y$ by means of a connection $\Gamma$ on $Y \rightarrow X$. Accordingly, substitution of the tangent-valued form

$$
\theta_{(1)} \circ \Gamma=\left(d y^{i}-\Gamma_{\lambda}^{i} d x^{\lambda}\right) \otimes \partial_{i}
$$

in the canonical splitting (3.2.4) leads to the dual splitting (3.1.10) of $T^{*} Y$ by means of a connection $\Gamma$.

Theorem 3.2.1: There is one-to-one correspondence between the connections $\Gamma$ on a fibre bundle $Y \rightarrow X$ and the global sections

$$
\begin{equation*}
\Gamma: Y \rightarrow J^{1} Y, \quad\left(x^{\lambda}, y^{i}, y_{\lambda}^{i}\right) \circ \Gamma=\left(x^{\lambda}, y^{i}, \Gamma_{\lambda}^{i}\right) \tag{3.2.5}
\end{equation*}
$$

of the affine jet bundle $J^{1} Y \rightarrow Y$.
There are the following corollaries of this theorem.

- Since $J^{1} Y \rightarrow Y$ is affine, a connection on a fibre bundle $Y \rightarrow X$ exists in accordance with Theorem 1.1.4.
- Connections on a fibre bundle $Y \rightarrow X$ make up an affine space modelled over the vector space of soldering forms on $Y \rightarrow X$, i.e., sections of the vector bundle (2.1.4).
- Connection components possess the coordinate transformation law

$$
\Gamma_{\lambda}^{\prime i}=\frac{\partial x^{\mu}}{\partial x^{\prime \lambda}}\left(\partial_{\mu}+\Gamma_{\mu}^{j} \partial_{j}\right) y^{\prime i}
$$

- Every connection $\Gamma(3.2 .5)$ on a fibre bundle $Y \rightarrow X$ yields the first order differential operator

$$
\begin{align*}
& D_{\Gamma}: J^{1} Y \underset{Y}{\rightarrow} T^{*} X \underset{Y}{\otimes} V Y  \tag{3.2.6}\\
& D_{\Gamma}=\lambda_{(1)}-\Gamma \circ \pi_{0}^{1}=\left(y_{\lambda}^{i}-\Gamma_{\lambda}^{i}\right) d x^{\lambda} \otimes \partial_{i}
\end{align*}
$$

on $Y$ called the covariant differential relative to the connection $\Gamma$. If $s: X \rightarrow Y$ is a section, from (3.2.6) we obtain its covariant differential

$$
\begin{align*}
\nabla^{\Gamma} s & =D_{\Gamma} \circ J^{1} s: X \rightarrow T^{*} X \otimes V Y,  \tag{3.2.7}\\
\nabla^{\Gamma} s & =\left(\partial_{\lambda} s^{i}-\Gamma_{\lambda}^{i} \circ s\right) d x^{\lambda} \otimes \partial_{i},
\end{align*}
$$

and the covariant derivative

$$
\left.\nabla_{\tau}^{\Gamma}=\tau\right\rfloor \nabla^{\Gamma}
$$

along a vector field $\tau$ on $X$. A section $s$ is said to be an integral section of a connection $\Gamma$ if it belongs to the kernel of the covariant differential $D_{\Gamma}$ (3.2.6), i.e.,

$$
\begin{equation*}
\nabla^{\Gamma} s=0 \quad \text { or } \quad J^{1} s=\Gamma \circ s . \tag{3.2.8}
\end{equation*}
$$

Theorem 3.2.2: For any global section $s: X \rightarrow Y$, there always exists a connection $\Gamma$ such that $s$ is an integral section of $\Gamma$.

Treating connections as jet bundle sections, one comes to the following two constructions.
(i) Let $Y$ and $Y^{\prime}$ be fibre bundles over the same base $X$. Given a connection $\Gamma$ on $Y \rightarrow X$ and a connection $\Gamma^{\prime}$ on $Y^{\prime} \rightarrow X$, the bundle product $Y \times Y^{\prime}$ is provided with the product connection

$$
\begin{align*}
& \Gamma \times \Gamma^{\prime}: Y \underset{X}{Y \times} Y^{\prime} \rightarrow J^{1}\left(\underset{X}{Y \times Y^{\prime}}\right)=J^{1} Y \times \underset{X}{\times} J^{1} Y^{\prime}, \\
& \Gamma \times \Gamma^{\prime}=d x^{\lambda} \otimes\left(\partial_{\lambda}+\Gamma_{\lambda}^{i} \frac{\partial}{\partial y^{i}}+\Gamma_{\lambda}^{\prime j} \frac{\partial}{\partial y^{\prime j}}\right) . \tag{3.2.9}
\end{align*}
$$

(ii) Let $i_{Y}: Y \rightarrow Y^{\prime}$ be a subbundle of a fibre bundle $Y^{\prime} \rightarrow X$ and $\Gamma^{\prime}$ a connection on $Y^{\prime} \rightarrow X$. If there exists a connection $\Gamma$ on $Y \rightarrow X$ such that the diagram

is commutative, we say that $\Gamma^{\prime}$ is reducible to a connection $\Gamma$. The following conditions are equivalent:

- $\Gamma^{\prime}$ is reducible to $\Gamma$;
- $T i_{Y}(H Y)=\left.H Y^{\prime}\right|_{i_{Y}(Y)}$, where $H Y \subset T Y$ and $H Y^{\prime} \subset T Y^{\prime}$ are the horizontal subbundles determined by $\Gamma$ and $\Gamma^{\prime}$, respectively;
- for every vector field $\tau$ on $X$, the vector fields $\Gamma \tau$ and $\Gamma^{\prime} \tau$ are related as follows:

$$
\begin{equation*}
T i_{Y} \circ \Gamma \tau=\Gamma^{\prime} \tau \circ i_{Y} \tag{3.2.10}
\end{equation*}
$$

### 3.3 Curvature and torsion

Let $\Gamma$ be a connection on a fibre bundle $Y \rightarrow X$. Its curvature is defined as the Nijenhuis differential

$$
\begin{align*}
& R=\frac{1}{2} d_{\Gamma} \Gamma=\frac{1}{2}[\Gamma, \Gamma]_{\mathrm{FN}}: Y \rightarrow \stackrel{2}{\wedge} T^{*} X \otimes V Y,  \tag{3.3.1}\\
& R=\frac{1}{2} R_{\lambda \mu}^{i} d x^{\lambda} \wedge d x^{\mu} \otimes \partial_{i},  \tag{3.3.2}\\
& R_{\lambda \mu}^{i}=\partial_{\lambda} \Gamma_{\mu}^{i}-\partial_{\mu} \Gamma_{\lambda}^{i}+\Gamma_{\lambda}^{j} \partial_{j} \Gamma_{\mu}^{i}-\Gamma_{\mu}^{j} \partial_{j} \Gamma_{\lambda}^{i} .
\end{align*}
$$

This is a $V Y$-valued horizontal two-form on $Y$. Given vector fields $\tau$, $\tau^{\prime}$ on $X$ and their horizontal lifts $\Gamma \tau$ and $\Gamma \tau^{\prime}$ (3.1.6) on $Y$, we have the relation

$$
\begin{equation*}
R\left(\tau, \tau^{\prime}\right)=-\Gamma\left[\tau, \tau^{\prime}\right]+\left[\Gamma \tau, \Gamma \tau^{\prime}\right]=\tau^{\lambda} \tau^{\prime \mu} R_{\lambda \mu}^{i} \partial_{i} \tag{3.3.3}
\end{equation*}
$$

The curvature (3.3.1) obeys the identities

$$
\begin{align*}
& {[R, R]_{\mathrm{FN}}=0}  \tag{3.3.4}\\
& d_{\Gamma} R=[\Gamma, R]_{\mathrm{FN}}=0 \tag{3.3.5}
\end{align*}
$$

They result from the identity (1.4.9) and the graded Jacobi identity (1.4.10), respectively. The identity (3.3.5) is called the second Bianchi identity. It takes the coordinate form

$$
\begin{equation*}
\sum_{(\lambda \mu \nu)}\left(\partial_{\lambda} R_{\mu \nu}^{i}+\Gamma_{\lambda}^{j} \partial_{j} R_{\mu \nu}^{i}-\partial_{j} \Gamma_{\lambda}^{i} R_{\mu \nu}^{j}\right)=0 \tag{3.3.6}
\end{equation*}
$$

where the sum is cyclic over the indices $\lambda, \mu$ and $\nu$.
Given a soldering form $\sigma$, one defines the soldered curvature

$$
\begin{align*}
\rho & =\frac{1}{2} d_{\sigma} \sigma=\frac{1}{2}[\sigma, \sigma]_{\mathrm{FN}}: Y \rightarrow \stackrel{2}{\wedge} T^{*} X \otimes V Y,  \tag{3.3.7}\\
\rho & =\frac{1}{2} \rho_{\lambda \mu}^{i} d x^{\lambda} \wedge d x^{\mu} \otimes \partial_{i}, \quad \rho_{\lambda \mu}^{i}=\sigma_{\lambda}^{j} \partial_{j} \sigma_{\mu}^{i}-\sigma_{\mu}^{j} \partial_{j} \sigma_{\lambda}^{i} .
\end{align*}
$$

It fulfills the identities

$$
[\rho, \rho]_{\mathrm{FN}}=0, \quad d_{\sigma} \rho=[\sigma, \rho]_{\mathrm{FN}}=0
$$

similar to (3.3.4) - (3.3.5).
Given a connection $\Gamma$ and a soldering form $\sigma$, the torsion form of $\Gamma$ with respect to $\sigma$ is defined as

$$
\begin{align*}
& T=d_{\Gamma} \sigma=d_{\sigma} \Gamma: Y \rightarrow \stackrel{2}{\wedge} T^{*} X \otimes V Y \\
& T=\left(\partial_{\lambda} \sigma_{\mu}^{i}+\Gamma_{\lambda}^{j} \partial_{j} \sigma_{\mu}^{i}-\partial_{j} \Gamma_{\lambda}^{i} \sigma_{\mu}^{j}\right) d x^{\lambda} \wedge d x^{\mu} \otimes \partial_{i} \tag{3.3.8}
\end{align*}
$$

It obeys the first Bianchi identity

$$
\begin{equation*}
d_{\Gamma} T=d_{\Gamma}^{2} \sigma=[R, \sigma]_{\mathrm{FN}}=-d_{\sigma} R \tag{3.3.9}
\end{equation*}
$$

If $\Gamma^{\prime}=\Gamma+\sigma$, we have the relations

$$
\begin{align*}
& T^{\prime}=T+2 \rho  \tag{3.3.10}\\
& R^{\prime}=R+\rho+T \tag{3.3.11}
\end{align*}
$$

### 3.4 Linear and affine connections

A connection $\Gamma$ on a vector bundle $Y \rightarrow X$ is called the linear connection if the section

$$
\begin{equation*}
\Gamma: Y \rightarrow J^{1} Y, \quad \Gamma=d x^{\lambda} \otimes\left(\partial_{\lambda}+\Gamma_{\lambda}{ }_{j}{ }_{j}(x) y^{j} \partial_{i}\right) \tag{3.4.1}
\end{equation*}
$$

is a linear bundle morphism over $X$. Note that linear connections are principal connections, and they always exist (see Theorem 4.4.1).

The curvature $R$ (3.3.2) of a linear connection $\Gamma$ (3.4.1) reads

$$
\begin{align*}
& R=\frac{1}{2} R_{\lambda \mu}{ }^{i}{ }_{j}(x) y^{j} d x^{\lambda} \wedge d x^{\mu} \otimes \partial_{i}, \\
& R_{\lambda \mu}{ }^{i}{ }_{j}=\partial_{\lambda} \Gamma_{\mu}{ }^{i}{ }_{j}-\partial_{\mu} \Gamma_{\lambda}{ }^{i}{ }_{j}+\Gamma_{\lambda}{ }^{h}{ }_{j} \Gamma_{\mu}{ }^{i}{ }_{h}-\Gamma_{\mu}{ }^{h}{ }_{j} \Gamma_{\lambda}{ }^{i}{ }_{h} . \tag{3.4.2}
\end{align*}
$$

Due to the vertical splitting (1.2.8), we have the linear morphism

$$
\begin{equation*}
R: Y \ni y^{i} e_{i} \rightarrow \frac{1}{2} R_{\lambda \mu}{ }^{i}{ }_{j} y^{j} d x^{\lambda} \wedge d x^{\mu} \otimes e_{i} \in \mathcal{O}^{2}(X) \otimes Y \tag{3.4.3}
\end{equation*}
$$

There are the following standard constructions of new linear connections from the old ones.

- Let $Y \rightarrow X$ be a vector bundle, coordinated by $\left(x^{\lambda}, y^{i}\right)$, and $Y^{*} \rightarrow X$ its dual, coordinated by $\left(x^{\lambda}, y_{i}\right)$. Any linear connection $\Gamma$ (3.4.1) on a vector bundle $Y \rightarrow X$ defines the dual linear connection

$$
\begin{equation*}
\Gamma^{*}=d x^{\lambda} \otimes\left(\partial_{\lambda}-\Gamma_{\lambda}{ }^{j}{ }_{i}(x) y_{j} \partial^{i}\right) \tag{3.4.4}
\end{equation*}
$$

on $Y^{*} \rightarrow X$.

- Let $\Gamma$ and $\Gamma^{\prime}$ be linear connections on vector bundles $Y \rightarrow X$ and $Y^{\prime} \rightarrow X$, respectively. The direct sum connection $\Gamma \oplus \Gamma^{\prime}$ on their Whitney sum $Y \oplus Y^{\prime}$ is defined as the product connection (3.2.9).
- Let $Y$ coordinated by $\left(x^{\lambda}, y^{i}\right)$ and $Y^{\prime}$ coordinated by $\left(x^{\lambda}, y^{a}\right)$ be vector bundles over the same base $X$. Their tensor product $Y \otimes Y^{\prime}$ is endowed with the bundle coordinates $\left(x^{\lambda}, y^{i a}\right)$. Linear connections $\Gamma$ and $\Gamma^{\prime}$ on $Y \rightarrow X$ and $Y^{\prime} \rightarrow X$ define the linear tensor product connection

$$
\begin{equation*}
\Gamma \otimes \Gamma^{\prime}=d x^{\lambda} \otimes\left[\partial_{\lambda}+\left(\Gamma_{\lambda}{ }^{i}{ }_{j} y^{j a}+\Gamma_{\lambda}^{\prime}{ }^{a}{ }_{b} y^{i b}\right) \frac{\partial}{\partial y^{i a}}\right] \tag{3.4.5}
\end{equation*}
$$

on $Y \otimes Y^{\prime}$.
An important example of linear connections is a linear connection

$$
\begin{equation*}
\Gamma=d x^{\lambda} \otimes\left(\partial_{\lambda}+\Gamma_{\lambda}{ }_{\nu}{ }_{\nu} \dot{x}^{\nu} \dot{\partial}_{\mu}\right) \tag{3.4.6}
\end{equation*}
$$

on the tangent bundle $T X$ of a manifold $X$. We agree to call it a world connection on a manifold $X$. The dual world connection (3.4.4) on the cotangent bundle $T^{*} X$ is

$$
\begin{equation*}
\Gamma^{*}=d x^{\lambda} \otimes\left(\partial_{\lambda}-\Gamma_{\lambda}{ }_{\nu}^{\mu} \dot{x}_{\mu} \dot{\partial}^{\nu}\right) \tag{3.4.7}
\end{equation*}
$$

Then, using the construction of the tensor product connection (3.4.5), one can introduce the corresponding linear world connection on an arbitrary tensor bundle $T$ (1.2.5).

Remark 3.4.1: It should be emphasized that the expressions (3.4.6) and (3.4.7) for a world connection differ in a minus sign from those usually used in the physical literature.

The curvature of a world connection is defined as the curvature $R$ (3.4.2) of the connection $\Gamma$ (3.4.6) on the tangent bundle $T X$. It reads

$$
\begin{align*}
& R=\frac{1}{2} R_{\lambda \mu}{ }^{\alpha}{ }_{\beta} \dot{x}^{\beta} d x^{\lambda} \wedge d x^{\mu} \otimes \dot{\partial}_{\alpha},  \tag{3.4.8}\\
& R_{\lambda \mu}{ }^{\alpha}{ }_{\beta}=\partial_{\lambda} \Gamma_{\mu}{ }^{\alpha}{ }_{\beta}-\partial_{\mu} \Gamma_{\lambda}{ }^{\alpha}{ }_{\beta}+\Gamma_{\lambda}{ }^{\gamma}{ }_{\beta} \Gamma_{\mu}{ }^{\alpha}{ }_{\gamma}-\Gamma_{\mu}{ }^{\gamma}{ }_{\beta} \Gamma_{\lambda}{ }^{\alpha}{ }_{\gamma} .
\end{align*}
$$

By the torsion of a world connection is meant the torsion (3.3.8) of the connection $\Gamma$ (3.4.6) on the tangent bundle $T X$ with respect to the canonical soldering form $\theta_{J}$ (1.4.7):

$$
\begin{equation*}
T=\frac{1}{2} T_{\mu}{ }^{\nu}{ }_{\lambda} d x^{\lambda} \wedge d x^{\mu} \otimes \dot{\partial}_{\nu}, \quad T_{\mu}{ }^{\nu}{ }_{\lambda}=\Gamma_{\mu}{ }^{\nu}{ }_{\lambda}-\Gamma_{\lambda}{ }^{\nu}{ }_{\mu} . \tag{3.4.9}
\end{equation*}
$$

A world connection is said to be symmetric if its torsion (3.4.9) vanishes, i.e., $\Gamma_{\mu}{ }^{\nu}{ }_{\lambda}=\Gamma_{\lambda}{ }^{\nu}{ }_{\mu}$.

Remark 3.4.2: For any vector field $\tau$ on a manifold $X$, there exists a connection $\Gamma$ on the tangent bundle $T X \rightarrow X$ such that $\tau$ is an integral section of $\Gamma$, but this connection is not necessarily linear. If a vector field $\tau$ is non-vanishing at a point $x \in X$, then there exists a local symmetric world connection $\Gamma$ (3.4.6) around $x$ for which $\tau$ is an integral section

$$
\begin{equation*}
\partial_{\nu} \tau^{\alpha}=\Gamma_{\nu}{ }^{\alpha}{ }_{\beta} \tau^{\beta} . \tag{3.4.10}
\end{equation*}
$$

Then the canonical lift $\widetilde{\tau}$ (1.3.4) of $\tau$ onto $T X$ can be seen locally as the horizontal lift $\Gamma \tau$ (3.1.6) of $\tau$ by means of this connection.

Remark 3.4.3: Every manifold $X$ can be provided with a nondegenerate fibre metric

$$
g \in \vee^{2} \mathcal{O}^{1}(X), \quad g=g_{\lambda \mu} d x^{\lambda} \otimes d x^{\mu}
$$

in the tangent bundle $T X$, and with the corresponding metric

$$
g \in \vee^{2} \mathcal{T}^{1}(X), \quad g=g^{\lambda \mu} \partial_{\lambda} \otimes \partial_{\mu}
$$

in the cotangent bundle $T^{*} X$. We call it a world metric on $X$. For any world metric $g$, there exists a unique symmetric world connection $\Gamma$ (3.4.6) with the components

$$
\begin{equation*}
\Gamma_{\lambda}{ }^{\nu}{ }_{\mu}=\left\{\lambda^{\nu}{ }_{\mu}\right\}=-\frac{1}{2} g^{\nu \rho}\left(\partial_{\lambda} g_{\rho \mu}+\partial_{\mu} g_{\rho \lambda}-\partial_{\rho} g_{\lambda \mu}\right), \tag{3.4.11}
\end{equation*}
$$

called the Christoffel symbols, such that $g$ is an integral section of $\Gamma$, i.e.

$$
\partial_{\lambda} g^{\alpha \beta}=g^{\alpha \gamma}\left\{\lambda^{\beta}{ }_{\gamma}\right\}+g^{\beta \gamma}\left\{\lambda^{\alpha}{ }_{\gamma}\right\} .
$$

It is called the Levi-Civita connection associated to $g$.
Let $Y \rightarrow X$ be an affine bundle modelled over a vector bundle $\bar{Y} \rightarrow X$. A connection $\Gamma$ on $Y \rightarrow X$ is called an affine connection if the section $\Gamma: Y \rightarrow J^{1} Y(3.2 .5)$ is an affine bundle morphism over $X$. Associated to principal connections, affine connections always exist (see Theorem 4.4.1).

For any affine connection $\Gamma: Y \rightarrow J^{1} Y$, the corresponding linear derivative $\bar{\Gamma}: \bar{Y} \rightarrow J^{1} \bar{Y}(1.2 .14)$ defines a unique linear connection on the vector bundle $\bar{Y} \rightarrow X$. Since every vector bundle is an affine bundle, any linear connection on a vector bundle also is an affine connection.

With respect to affine bundle coordinates $\left(x^{\lambda}, y^{i}\right)$ on $Y$, an affine connection $\Gamma$ on $Y \rightarrow X$ reads

$$
\begin{equation*}
\Gamma_{\lambda}^{i}=\Gamma_{\lambda}{ }^{i}{ }_{j}(x) y^{j}+\sigma_{\lambda}^{i}(x) . \tag{3.4.12}
\end{equation*}
$$

The coordinate expression of the associated linear connection is

$$
\begin{equation*}
\bar{\Gamma}_{\lambda}^{i}=\Gamma_{\lambda}{ }^{i}{ }_{j}(x) \bar{y}^{j}, \tag{3.4.13}
\end{equation*}
$$

where $\left(x^{\lambda}, \bar{y}^{i}\right)$ are the associated linear bundle coordinates on $\bar{Y}$.
Affine connections on an affine bundle $Y \rightarrow X$ constitute an affine space modelled over the soldering forms on $Y \rightarrow X$. In view of the
vertical splitting (1.2.15), these soldering forms can be seen as global sections of the vector bundle

$$
T^{*} X \underset{X}{\otimes} \bar{Y} \rightarrow X
$$

If $Y \rightarrow X$ is a vector bundle, both the affine connection $\Gamma$ (3.4.12) and the associated linear connection $\bar{\Gamma}$ are connections on the same vector bundle $Y \rightarrow X$, and their difference is a basic soldering form on $Y$. Thus, every affine connection on a vector bundle $Y \rightarrow X$ is the sum of a linear connection and a basic soldering form on $Y \rightarrow X$.

Given an affine connection $\Gamma$ on a vector bundle $Y \rightarrow X$, let $R$ and $\bar{R}$ be the curvatures of a connection $\Gamma$ and the associated linear connection $\bar{\Gamma}$, respectively. It is readily observed that $R=\bar{R}+T$, where the $V Y$ valued two-form

$$
\begin{align*}
& T=d_{\Gamma} \sigma=d_{\sigma} \Gamma: X \rightarrow \stackrel{2}{\wedge} T^{*} X \underset{X}{\otimes} V Y  \tag{3.4.14}\\
& T=\frac{1}{2} T_{\lambda \mu}^{i} d x^{\lambda} \wedge d x^{\mu} \otimes \partial_{i}, \\
& T_{\lambda \mu}^{i}=\partial_{\lambda} \sigma_{\mu}^{i}-\partial_{\mu} \sigma_{\lambda}^{i}+\sigma_{\lambda}^{h} \Gamma_{\mu}{ }^{i}{ }_{h}-\sigma_{\mu}^{h} \Gamma_{\lambda}{ }^{i}{ }_{h}
\end{align*}
$$

is the torsion (3.3.8) of $\Gamma$ with respect to the basic soldering form $\sigma$.
In particular, let us consider the tangent bundle $T X$ of a manifold $X$. We have the canonical soldering form $\sigma=\theta_{J}=\theta_{X}$ (1.4.7) on $T X$. Given an arbitrary world connection $\Gamma$ (3.4.6) on $T X$, the corresponding affine connection

$$
\begin{equation*}
A=\Gamma+\theta_{X}, \quad A_{\lambda}^{\mu}=\Gamma_{\lambda}{ }^{\mu}{ }_{\nu} \dot{x}^{\nu}+\delta_{\lambda}^{\mu} \tag{3.4.15}
\end{equation*}
$$

on $T X$ is called the Cartan connection. Since the soldered curvature $\rho$ (3.3.7) of $\theta_{J}$ equals zero, the torsion (3.3.10) of the Cartan connection coincides with the torsion $T$ (3.4.9) of the world connection $\Gamma$, while its curvature (3.3.11) is the sum $R+T$ of the curvature and the torsion of $\Gamma$.

### 3.5 Flat connections

By a flat or curvature-free connection is meant a connection which satisfies the following equivalent conditions.

Theorem 3.5.1: Let $\Gamma$ be a connection on a fibre bundle $Y \rightarrow X$. The following assertions are equivalent.
(i) The curvature $R$ of a connection $\Gamma$ vanishes identically, i.e., $R \equiv 0$.
(ii) The horizontal lift (3.1.7) of vector fields on $X$ onto $Y$ is an $\mathbb{R}$ linear Lie algebra morphism (in accordance with the formula (3.3.3)).
(iii) The horizontal distribution is involutive.
(iv) There exists a local integral section for $\Gamma$ through any point $y \in$ $Y$.

By virtue of Theorem 1.3.3 and item (iii) of Theorem 3.5.1, a flat connection $\Gamma$ on a fibre bundle $Y \rightarrow X$ yields a horizontal foliation on $Y$, transversal to the fibration $Y \rightarrow X$. The leaf of this foliation through a point $y \in Y$ is defined locally by an integral section $s_{y}$ for the connection $\Gamma$ through $y$. Conversely, let a fibre bundle $Y \rightarrow X$ admit a transversal foliation such that, for each point $y \in Y$, the leaf of this foliation through $y$ is locally defined by a section $s_{y}$ of $Y \rightarrow X$ through $y$. Then the map

$$
\Gamma: Y \rightarrow J^{1} Y, \quad \Gamma(y)=j_{x}^{1} s_{y}, \quad \pi(y)=x
$$

introduces a flat connection on $Y \rightarrow X$. Thus, there is one-to-one correspondence between the flat connections and the transversal foliations of a fibre bundle $Y \rightarrow X$.

Given a transversal foliation on a fibre bundle $Y \rightarrow X$, there exists the associated atlas of bundle coordinates $\left(x^{\lambda}, y^{i}\right)$ of $Y$ such that every leaf of this foliation is locally generated by the equations $y^{i}=$ const., and the transition functions $y^{i} \rightarrow y^{\prime i}\left(y^{j}\right)$ are independent of the base coordinates $x^{\lambda}$. This is called the atlas of constant local trivializations. Two such atlases are said to be equivalent if their union also is an atlas of
constant local trivializations. They are associated to the same horizontal foliation. Thus, we come to the following assertion.

Theorem 3.5.2: There is one-to-one correspondence between the flat connections $\Gamma$ on a fibre bundle $Y \rightarrow X$ and the equivalence classes of atlases of constant local trivializations of $Y$ such that

$$
\Gamma=d x^{\lambda} \otimes \partial_{\lambda}
$$

relative to these atlases.
In particular, if $Y \rightarrow X$ is a trivial bundle, one associates to each its trivialization a flat connection represented by the global zero section $\hat{0}(Y)$ of $J^{1} Y \rightarrow Y$ with respect to this trivialization (see Remark 2.1.2).

### 3.6 Connections on composite bundles

Let $Y \rightarrow \Sigma \rightarrow X$ be a composite bundle (1.1.10). Let us consider the jet manifolds $J^{1} \Sigma, J_{\Sigma}^{1} Y$, and $J^{1} Y$ of the fibre bundles

$$
\Sigma \rightarrow X, \quad Y \rightarrow \Sigma, \quad Y \rightarrow X
$$

respectively. They are provided with the adapted coordinates

$$
\left(x^{\lambda}, \sigma^{m}, \sigma_{\lambda}^{m}\right), \quad\left(x^{\lambda}, \sigma^{m}, y^{i}, \widetilde{y}_{\lambda}^{i}, y_{m}^{i}\right), \quad\left(x^{\lambda}, \sigma^{m}, y^{i}, \sigma_{\lambda}^{m}, y_{\lambda}^{i}\right)
$$

One can show the following.
Theorem 3.6.1: There is the canonical map

$$
\begin{equation*}
\varrho: J^{1} \Sigma \underset{\Sigma}{\times} J_{\Sigma}^{1} Y \underset{Y}{\longrightarrow} J^{1} Y, \quad y_{\lambda}^{i} \circ \varrho=y_{m}^{i} \sigma_{\lambda}^{m}+\widetilde{y}_{\lambda}^{i} . \tag{3.6.1}
\end{equation*}
$$

Using the canonical map (3.6.1), we can get the relations between connections on the fibre bundles $Y \rightarrow X, Y \rightarrow \Sigma$ and $\Sigma \rightarrow X$. These connections are given by the corresponding connection forms

$$
\begin{align*}
& \gamma=d x^{\lambda} \otimes\left(\partial_{\lambda}+\gamma_{\lambda}^{m} \partial_{m}+\gamma_{\lambda}^{i} \partial_{i}\right)  \tag{3.6.2}\\
& A_{\Sigma}=d x^{\lambda} \otimes\left(\partial_{\lambda}+A_{\lambda}^{i} \partial_{i}\right)+d \sigma^{m} \otimes\left(\partial_{m}+A_{m}^{i} \partial_{i}\right)  \tag{3.6.3}\\
& \Gamma=d x^{\lambda} \otimes\left(\partial_{\lambda}+\Gamma_{\lambda}^{m} \partial_{m}\right) \tag{3.6.4}
\end{align*}
$$

A connection $\gamma$ (3.6.2) on the fibre bundle $Y \rightarrow X$ is called projectable onto a connection $\Gamma$ (3.6.4) on the fibre bundle $\Sigma \rightarrow X$ if, for any vector field $\tau$ on $X$, its horizontal lift $\gamma \tau$ on $Y$ by means of the connection $\gamma$ is a projectable vector field over the horizontal lift $\Gamma \tau$ of $\tau$ on $\Sigma$ by means of the connection $\Gamma$. This property holds iff $\gamma_{\lambda}^{m}=\Gamma_{\lambda}^{m}$, i.e., components $\gamma_{\lambda}^{m}$ of the connection $\gamma$ (3.6.2) must be independent of the fibre coordinates $y^{i}$.

A connection $A_{\Sigma}$ (3.6.3) on the fibre bundle $Y \rightarrow \Sigma$ and a connection $\Gamma$ (3.6.4) on the fibre bundle $\Sigma \rightarrow X$ define a connection on the composite bundle $Y \rightarrow X$ as the composition of bundle morphisms

$$
\gamma: Y \underset{X}{\times} T X \xrightarrow{(\mathrm{Id}, \Gamma)} Y \underset{\Sigma}{\times} T \Sigma \xrightarrow{A_{\Sigma}} T Y .
$$

It is called the composite connection. This composite connection reads

$$
\begin{equation*}
\gamma=d x^{\lambda} \otimes\left(\partial_{\lambda}+\Gamma_{\lambda}^{m} \partial_{m}+\left(A_{\lambda}^{i}+A_{m}^{i} \Gamma_{\lambda}^{m}\right) \partial_{i}\right) . \tag{3.6.5}
\end{equation*}
$$

It is projectable onto $\Gamma$. Moreover, this is a unique connection such that the horizontal lift $\gamma \tau$ on $Y$ of a vector field $\tau$ on $X$ by means of the composite connection $\gamma$ (3.6.5) coincides with the composition $A_{\Sigma}(\Gamma \tau)$ of horizontal lifts of $\tau$ on $\Sigma$ by means of the connection $\Gamma$ and then on $Y$ by means of the connection $A_{\Sigma}$. For the sake of brevity, let us write $\gamma=A_{\Sigma} \circ \Gamma$.

Given a composite bundle $Y$ (1.1.10), there are the exact sequences of vector bundles over $Y$ :

$$
\begin{align*}
& 0 \rightarrow V_{\Sigma} Y \longrightarrow V Y \rightarrow Y \stackrel{\times}{\Sigma} V \Sigma \rightarrow 0  \tag{3.6.6}\\
& 0 \rightarrow Y \underset{\Sigma}{\times V^{*} \Sigma \longrightarrow V^{*} Y \rightarrow V_{\Sigma}^{*} Y \rightarrow 0} \tag{3.6.7}
\end{align*}
$$

where $V_{\Sigma} Y$ and $V_{\Sigma}^{*} Y$ are the vertical tangent and the vertical cotangent bundles of $Y \rightarrow \Sigma$ which are coordinated by ( $x^{\lambda}, \sigma^{m}, y^{i}, \dot{y}^{i}$ ) and $\left(x^{\lambda}, \sigma^{m}, y^{i}, \dot{y}_{i}\right)$, respectively. Let us consider a splitting of these exact sequences

$$
\begin{equation*}
\left.B: V Y \ni \dot{y}^{i} \partial_{i}+\dot{\sigma}^{m} \partial_{m} \rightarrow\left(\dot{y}^{i} \partial_{i}+\dot{\sigma}^{m} \partial_{m}\right)\right\rfloor B= \tag{3.6.8}
\end{equation*}
$$

$$
\begin{align*}
& \quad\left(\dot{y}^{i}-\dot{\sigma}^{m} B_{m}^{i}\right) \partial_{i} \in V_{\Sigma} Y \\
& \left.B: V_{\Sigma}^{*} Y \ni \bar{d} y^{i} \rightarrow B\right\rfloor \bar{d} y^{i}=\bar{d} y^{i}-B_{m}^{i} \bar{d} \sigma^{m} \in V^{*} Y, \tag{3.6.9}
\end{align*}
$$

given by the form

$$
\begin{equation*}
B=\left(\bar{d} y^{i}-B_{m}^{i} \bar{d} \sigma^{m}\right) \otimes \partial_{i} \tag{3.6.10}
\end{equation*}
$$

Then the connection $\gamma$ (3.6.2) on $Y \rightarrow X$ and the splitting $B$ (3.6.8) define the connection

$$
\begin{align*}
& A_{\Sigma}=B \circ \gamma: T Y \rightarrow V Y \rightarrow V_{\Sigma} Y  \tag{3.6.11}\\
& A_{\Sigma}=d x^{\lambda} \otimes\left(\partial_{\lambda}+\left(\gamma_{\lambda}^{i}-B_{m}^{i} \gamma_{\lambda}^{m}\right) \partial_{i}\right)+d \sigma^{m} \otimes\left(\partial_{m}+B_{m}^{i} \partial_{i}\right)
\end{align*}
$$

on the fibre bundle $Y \rightarrow \Sigma$.
Conversely, every connection $A_{\Sigma}$ (3.6.3) on the fibre bundle $Y \rightarrow \Sigma$ yields the splitting

$$
\begin{equation*}
A_{\Sigma}: T Y \supset V Y \ni \dot{y}^{i} \partial_{i}+\dot{\sigma}^{m} \partial_{m} \rightarrow\left(\dot{y}^{i}-A_{m}^{i} \dot{\sigma}^{m}\right) \partial_{i} \tag{3.6.12}
\end{equation*}
$$

of the exact sequence (3.6.6). Using this splitting, one can construct a first order differential operator

$$
\begin{align*}
& \widetilde{D}: J^{1} Y \rightarrow T^{*} X \otimes_{Y}^{\otimes} V_{\Sigma} Y,  \tag{3.6.13}\\
& \widetilde{D}=d x^{\lambda} \otimes\left(y_{\lambda}^{i}-A_{\lambda}^{i}-A_{m}^{i} \sigma_{\lambda}^{m}\right) \partial_{i},
\end{align*}
$$

on the composite bundle $Y \rightarrow X$. It is called the vertical covariant differential. This operator also can be defined as the composition

$$
\widetilde{D}=\operatorname{pr}_{1} \circ D^{\gamma}: J^{1} Y \rightarrow T^{*} X \underset{Y}{\otimes} V Y \rightarrow T^{*} X \underset{Y}{\otimes} V Y_{\Sigma}
$$

where $D^{\gamma}$ is the covariant differential (3.2.6) relative to some composite connection $A_{\Sigma} \circ \Gamma$ (3.6.5), but $\widetilde{D}$ does not depend on the choice of a connection $\Gamma$ on the fibre bundle $\Sigma \rightarrow X$.

The vertical covariant differential (3.6.13) possesses the following important property. Let $h$ be a section of the fibre bundle $\Sigma \rightarrow X$, and let $Y^{h} \rightarrow X$ be the restriction (1.1.13) of the fibre bundle $Y \rightarrow \Sigma$ to $h(X) \subset \Sigma$. This is a subbundle

$$
i_{h}: Y^{h} \rightarrow Y
$$

of the fibre bundle $Y \rightarrow X$. Every connection $A_{\Sigma}$ (3.6.3) induces the pull-back connection

$$
\begin{equation*}
A_{h}=i_{h}^{*} A_{\Sigma}=d x^{\lambda} \otimes\left[\partial_{\lambda}+\left(\left(A_{m}^{i} \circ h\right) \partial_{\lambda} h^{m}+(A \circ h)_{\lambda}^{i}\right) \partial_{i}\right] \tag{3.6.14}
\end{equation*}
$$

on $Y^{h} \rightarrow X$. Then the restriction of $\widetilde{D}(3.6 .13)$ to

$$
J^{1} i_{h}\left(J^{1} Y^{h}\right) \subset J^{1} Y
$$

coincides with the familiar covariant differential $D^{A_{h}}(3.2 .6)$ on $Y^{h}$ relative to the pull-back connection $A_{h}$ (3.6.14).

Remark 3.6.1: Let $\Gamma: Y \rightarrow J^{1} Y$ be a connection on a fibre bundle $Y \rightarrow X$. In accordance with the canonical isomorphism $V J^{1} Y=J^{1} V Y$ (2.1.9), the vertical tangent map

$$
V \Gamma: V Y \rightarrow V J^{1} Y
$$

to $\Gamma$ defines the connection

$$
\begin{align*}
& V \Gamma: V Y \rightarrow J^{1} V Y \\
& V \Gamma=d x^{\lambda} \otimes\left(\partial_{\lambda}+\Gamma_{\lambda}^{i} \partial_{i}+\partial_{j} \Gamma_{\lambda}^{i} \dot{y}^{j} \dot{\partial}_{i}\right) \tag{3.6.15}
\end{align*}
$$

on the composite vertical tangent bundle

$$
V Y \rightarrow Y \rightarrow X
$$

This is called the vertical connection to $\Gamma$. Of course, the connection $V \Gamma$ projects onto $\Gamma$. Moreover, $V \Gamma$ is linear over $\Gamma$. Then the dual connection of $V \Gamma$ on the composite vertical cotangent bundle

$$
V^{*} Y \rightarrow Y \rightarrow X
$$

reads

$$
\begin{align*}
& V^{*} \Gamma: V^{*} Y \rightarrow J^{1} V^{*} Y, \\
& V^{*} \Gamma=d x^{\lambda} \otimes\left(\partial_{\lambda}+\Gamma_{\lambda}^{i} \partial_{i}-\partial_{j} \Gamma_{\lambda}^{i} \dot{y}_{i} \dot{\partial}^{j}\right) . \tag{3.6.16}
\end{align*}
$$

It is called the covertical connection to $\Gamma$. If $Y \rightarrow X$ is an affine bundle, the connection $V \Gamma$ (3.6.15) can be seen as the composite connection generated by the connection $\Gamma$ on $Y \rightarrow X$ and the linear connection

$$
\begin{equation*}
\tilde{\Gamma}=d x^{\lambda} \otimes\left(\partial_{\lambda}+\partial_{j} \Gamma_{\lambda}^{i} \dot{y}^{j} \dot{\partial}_{i}\right)+d y^{i} \otimes \partial_{i} \tag{3.6.17}
\end{equation*}
$$

on the vertical tangent bundle $V Y \rightarrow Y$.

## Chapter 4

## Geometry of principal bundles

Classical gauge theory is adequately formulated as Lagrangian field theory on principal and associated bundles where gauge potentials are identified with principal connections. The main ingredient in this formulation is the bundle of principal connections $C=J^{1} P / G$ whose sections are principal connections on a principal bundle $P$ with a structure group $G$.

### 4.1 Geometry of Lie groups

Let $G$ be a topological group which is not reduced to the unit 1. Let $V$ be a topological space. By a continuous action of $G$ on $V$ on the left is meant a continuous map

$$
\begin{equation*}
\zeta: G \times V \rightarrow V, \quad \zeta\left(g^{\prime} g, v\right)=\zeta\left(g^{\prime}, \zeta(g, v)\right) \tag{4.1.1}
\end{equation*}
$$

If there is no danger of confusion, we denote $\zeta(g, v)=g v$. One says that a group $G$ acts on $V$ on the right if the map (4.1.1) obeys the relations

$$
\zeta\left(g^{\prime} g, v\right)=\zeta\left(g, \zeta\left(g^{\prime}, v\right)\right)
$$

In this case, we agree to write $\zeta(g, v)=v g$.
Remark 4.1.1: $\quad$ Strictly speaking, by an action of a group $G$ on $V$ is meant a class of morphisms $\zeta$ (4.1.1) which differ from each other in
inner automorphisms of $G$, that is,

$$
\zeta^{\prime}(g, v)=\zeta\left(g^{\prime-1} g g^{\prime}, v\right)
$$

for some element $g^{\prime} \in G$.
An action of $G$ on $V$ is called:

- effective if there is no $g \neq \mathbf{1}$ such that $\zeta(g, v)=v$ for all $v \in V$,
- free if, for any two elements $v, v \in V$, there exists an element $g \in G$ such that $\zeta(g, v)=v^{\prime}$.
- transitive if there is no element $v \in V$ such that $\zeta(g, v)=v$ for all $g \in G$.
Unless otherwise stated, an action of a group is assumed to be effective. If an action $\zeta$ (4.1.1) of $G$ on $V$ is transitive, then $V$ is called the homogeneous space, homeomorphic to the quotient $V=G / H$ of $G$ with respect to some subgroup $H \subset G$. If an action $\zeta$ is both free and transitive, then $V$ is homeomorphic to the group space of $G$. For instance, this is the case of action of $G$ on itself by left $\left(\zeta=L_{G}\right)$ and right $\left(\zeta=R_{G}\right)$ multiplications.

Let $G$ be a connected real Lie group of finite dimension $\operatorname{dim} G>0$. A vector field $\xi$ on $G$ is called left-invariant if

$$
\xi(g)=T L_{g}(\xi(\mathbf{1})), \quad g \in G
$$

where $T L_{g}$ denotes the tangent morphism to the map

$$
L_{g}: G \rightarrow g G
$$

Accordingly, right-invariant vector fields $\xi$ on $G$ obey the condition

$$
\xi(g)=T R_{g}(\xi(\mathbf{1}))
$$

where $T R_{g}$ is the tangent morphism to the map

$$
T_{g}: G \rightarrow g G
$$

Let $\mathfrak{g}_{l}$ (resp. $\mathfrak{g}_{r}$ ) denote the Lie algebra of left-invariant (resp. rightinvariant) vector fields on $G$. They are called the left and right Lie
algebras of $G$, respectively. Every left-invariant vector field $\xi_{l}(g)$ (resp. a right-invariant vector field $\left.\xi_{r}(g)\right)$ can be associated to the element $v=\xi_{l}(\mathbf{1})$ (resp. $\left.v=\xi_{r}(\mathbf{1})\right)$ of the tangent space $T_{\mathbf{1}} G$ at the unit $\mathbf{1}$ of $G$. Accordingly, this tangent space is provided both with left and right Lie algebra structures. Given $v \in T_{1} G$, let $v_{l}(g)$ and $v_{r}(g)$ be the corresponding left-invariant and right-invariant vector fields on $G$, respectively. There is the relation

$$
v_{l}(g)=\left(T L_{g} \circ T R_{g}^{-1}\right)\left(v_{r}(g)\right)
$$

Let $\left\{\epsilon_{m}=\epsilon_{m}(\mathbf{1})\right\}$ (resp. $\left.\left\{\varepsilon_{m}=\varepsilon_{m}(\mathbf{1})\right\}\right)$ denote the basis for the left (resp. right) Lie algebra, and let $c_{m n}^{k}$ be the right structure constants:

$$
\left[\varepsilon_{m}, \varepsilon_{n}\right]=c_{m n}^{k} \varepsilon_{k}
$$

The map $g \rightarrow g^{-1}$ yields an isomorphism

$$
\mathfrak{g}_{l} \ni \epsilon_{m} \rightarrow \varepsilon_{m}=-\epsilon_{m} \in \mathfrak{g}_{r}
$$

of left and right Lie algebras.
The tangent bundle

$$
\begin{equation*}
\pi_{G}: T G \rightarrow G \tag{4.1.2}
\end{equation*}
$$

of a Lie group $G$ is trivial because of the isomorphisms

$$
\begin{aligned}
& \varrho_{l}: T G \ni q \rightarrow\left(g=\pi_{G}(q), T L_{g}^{-1}(q)\right) \in G \times \mathfrak{g}_{l} \\
& \varrho_{r}: T G \ni q \rightarrow\left(g=\pi_{G}(q), T R_{g}^{-1}(q)\right) \in G \times \mathfrak{g}_{r}
\end{aligned}
$$

Let $\zeta$ (4.1.1) be a smooth action of a Lie group $G$ on a smooth manifold $V$. Let us consider the tangent morphism

$$
\begin{equation*}
T \zeta: T G \times T V \rightarrow T V \tag{4.1.3}
\end{equation*}
$$

to this action. Given an element $g \in G$, the restriction of $T \zeta$ (4.1.3) to $(g, 0) \times T V$ is the tangent morphism $T \zeta_{g}$ to the map

$$
\zeta_{g}: g \times V \rightarrow V
$$

Therefore, the restriction

$$
\begin{equation*}
T \zeta_{G}: \widehat{0}(G) \times T V \rightarrow T V \tag{4.1.4}
\end{equation*}
$$

of the tangent morphism $T \zeta$ (4.1.3) to $\widehat{0}(G) \times T V$ (where $\hat{0}$ is the canonical zero section of $T G \rightarrow G)$ is called the tangent prolongation of a smooth action of $G$ on $V$.

In particular, the above mentioned morphisms

$$
T L_{g}=\left.T L_{G}\right|_{(g, 0) \times T G}, \quad T R_{g}=\left.T R_{G}\right|_{(g, 0) \times T G}
$$

are of this type. For instance, the morphism $T L_{G}\left(\right.$ resp. $\left.T R_{G}\right)(4.1 .4)$ defines the adjoint representation $g \rightarrow \operatorname{Ad}_{g}$ (resp. $g \rightarrow \operatorname{Ad}_{g^{-1}}$ ) of a group $G$ in its right Lie algebra $\mathfrak{g}_{r}$ (resp. left Lie algebra $\mathfrak{g}_{l}$ ) and the identity representation in its left (resp. right) one.

Restricting $T \zeta$ (4.1.3) to $T_{1} G \times \widehat{0}(V)$, one obtains a homomorphism of the right (resp. left) Lie algebra of $G$ to the Lie algebra $\mathcal{T}(V)$ of vector field on $V$ if $\zeta$ is a left (resp. right) action. We call this homomorphism a representation of the Lie algebra of $G$ in $V$. For instance, a vector field on a manifold $V$ associated to a local one-parameter group $G$ of diffeomorphisms of $V$ (see Section 1.3) is exactly an image of such a homomorphism of the one-dimensional Lie algebra of $G$ to $\mathcal{T}(V)$.

In particular, the adjoint representation $\mathrm{Ad}_{g}$ of a Lie group $G$ in its right Lie algebra $\mathfrak{g}_{r}$ yields the corresponding adjoint representation

$$
\begin{equation*}
\varepsilon^{\prime}: \varepsilon \rightarrow \operatorname{ad}_{\varepsilon^{\prime}}(\varepsilon)=\left[\varepsilon^{\prime}, \varepsilon\right], \quad \operatorname{ad}_{\varepsilon_{m}}\left(\varepsilon_{n}\right)=c_{m n}^{k} \varepsilon_{k}, \tag{4.1.5}
\end{equation*}
$$

of the right Lie algebra $\mathfrak{g}_{r}$ in itself. Accordingly, the adjoint representation of the left Lie algebra $\mathfrak{g}_{l}$ in itself reads

$$
\operatorname{ad}_{\epsilon_{m}}\left(\epsilon_{n}\right)=-c_{m n}^{k} \epsilon_{k},
$$

where $c_{m n}^{k}$ are the right structure constants (4.1.5).
Remark 4.1.2: Let $G$ be a matrix group, i.e., a subgroup of the algebra $M(V)$ of endomorphisms of some finite-dimensional vector space
$V$. Then its Lie algebras are Lie subalgebras of $M(V)$. In this case, the adjoint representation $\mathrm{Ad}_{g}$ of $G$ reads

$$
\begin{equation*}
\operatorname{Ad}_{g}(e)=g e g^{-1}, \quad e \in \mathfrak{g} \tag{4.1.6}
\end{equation*}
$$

An exterior form $\phi$ on a Lie group $G$ is said to be left-invariant (resp. right-invariant) if $\phi(\mathbf{1})=L_{g}^{*}(\phi(g))$ (resp. $\left.\phi(\mathbf{1})=R_{g}^{*}(\phi(g))\right)$. The exterior differential of a left-invariant (resp right-invariant) form is leftinvariant (resp. right-invariant). In particular, the left-invariant oneforms satisfy the Maurer-Cartan equation

$$
\begin{equation*}
d \phi\left(\epsilon, \epsilon^{\prime}\right)=-\frac{1}{2} \phi\left(\left[\epsilon, \epsilon^{\prime}\right]\right), \quad \epsilon, \epsilon^{\prime} \in \mathfrak{g}_{l} \tag{4.1.7}
\end{equation*}
$$

There is the canonical $\mathfrak{g}_{l}$-valued left-invariant one-form

$$
\begin{equation*}
\theta_{l}: T_{1} G \ni \epsilon \rightarrow \epsilon \in \mathfrak{g}_{l} \tag{4.1.8}
\end{equation*}
$$

on a Lie group $G$. The components $\theta_{l}^{m}$ of its decomposition $\theta_{l}=\theta_{l}^{m} \epsilon_{m}$ with respect to the basis for the left Lie algebra $\mathfrak{g}_{l}$ make up the basis for the space of left-invariant exterior one-forms on $G$ :

$$
\left.\epsilon_{m}\right\rfloor \theta_{l}^{n}=\delta_{m}^{n}
$$

The Maurer-Cartan equation (4.1.7), written with respect to this basis, reads

$$
d \theta_{l}^{m}=\frac{1}{2} c_{n k}^{m} \theta_{l}^{n} \wedge \theta_{l}^{k} .
$$

### 4.2 Bundles with structure groups

Principal bundles are particular bundles with a structure group. Since equivalence classes of these bundles are topological invariants (see Theorem 4.2.5), we consider continuous bundles with a structure topological group.

Let $G$ be a topological group. Let $\pi: Y \rightarrow X$ be a locally trivial continuous bundle (see Remark 1.1.1) whose typical fibre $V$ is provided with a certain left action (4.1.1) of a topological group $G$ (see Remark 4.1.1). Moreover, let $Y$ admit an atlas

$$
\begin{equation*}
\Psi=\left\{\left(U_{\alpha}, \psi_{\alpha}\right), \varrho_{\alpha \beta}\right\}, \quad \psi_{\alpha}=\varrho_{\alpha \beta} \psi_{\beta} \tag{4.2.1}
\end{equation*}
$$

whose transition functions $\varrho_{\alpha \beta}$ (1.1.3) factorize as

$$
\varrho_{\alpha \beta}: U_{\alpha} \cap U_{\beta} \times V \longrightarrow U_{\alpha} \cap U_{\beta} \times(G \times V) \xrightarrow{\operatorname{Id} \times \zeta} U_{\alpha} \cap U_{\beta} \times V \text { (4.2.2) }
$$

through local continuous $G$-valued functions

$$
\begin{equation*}
\varrho_{\alpha \beta}^{G}: U_{\alpha} \cap U_{\beta} \rightarrow G \tag{4.2.3}
\end{equation*}
$$

on $X$. This means that transition morphisms $\varrho_{\alpha \beta}(x)$ (1.1.6) are elements of $G$ acting on $V$. Transition functions (4.2.2) are called $G$-valued.

Provided with an atlas (4.2.1) with $G$-valued transition functions, a locally trivial continuous bundle $Y$ is called the bundle with a structure group $G$ or, in brief, a $G$-bundle. Two $G$-bundles $(Y, \Psi)$ and $\left(Y, \Psi^{\prime}\right)$ are called equivalent if their atlases $\Psi$ and $\Psi^{\prime}$ are equivalent. Atlases $\Psi$ and $\Psi^{\prime}$ with $G$-valued transition functions are said to be equivalent iff, given a common cover $\left\{U_{i}\right\}$ of $X$ for the union of these atlases, there exists a continuous $G$-valued function $g_{i}$ on each $U_{i}$ such that

$$
\begin{equation*}
\psi_{i}^{\prime}(x)=g_{i}(x) \psi_{i}(x), \quad x \in U_{i} \tag{4.2.4}
\end{equation*}
$$

Let $h(X, G, V)$ denote the set of equivalence classes of continuous bundles over $X$ with a structure group $G$ and a typical fibre $V$. In order to characterize this set, let us consider the presheaf $G_{\{U\}}^{0}$ of continuous $G$-valued functions on a topological space $X$. Let $G_{X}^{0}$ be the sheaf of germs of these functions generated by the canonical presheaf $G_{\{U\}}^{0}$, and let $H^{1}\left(X ; G_{X}^{0}\right)$ be the first cohomology of $X$ with coefficients in $G_{X}^{0}$ (see Remark 8.5.3). The group functions $\varrho_{\alpha \beta}^{G}(4.2 .3)$ obey the cocycle condition

$$
\varrho_{\alpha \beta}^{G} \varrho_{\beta \gamma}^{G}=\varrho_{\alpha \gamma}^{G}
$$

on overlaps $U_{\alpha} \cap U_{\beta} \cap U_{\gamma}$ (cf. (8.5.12)) and, consequently, they form a one-cocycle $\left\{\varrho_{\alpha \beta}^{G}\right\}$ of the presheaf $G_{\{U\}}^{0}$. This cocycle is a representative of some element of the first cohomology $H^{1}\left(X ; G_{X}^{0}\right)$ of $X$ with coefficients in the sheaf $G_{X}^{0}$.

Thus, any atlas of a $G$-bundle over $X$ defines an element of the cohomology set $H^{1}\left(X ; G_{X}^{0}\right)$. Moreover, it follows at once from the condition (4.2.4) that equivalent atlases define the same element of $H^{1}\left(X ; G_{X}^{0}\right)$. Thus, there is an injection

$$
\begin{equation*}
h(X, G, V) \rightarrow H^{1}\left(X ; G_{X}^{0}\right) \tag{4.2.5}
\end{equation*}
$$

of the set of equivalence classes of $G$-bundles over $X$ with a typical fibre $V$ to the first cohomology $H^{1}\left(X ; G_{X}^{0}\right)$ of $X$ with coefficients in the sheaf $G_{X}^{0}$. Moreover, the injection (4.2.5) is a bijection as follows.

Theorem 4.2.1: There is one-to-one correspondence between the equivalence classes of $G$-bundles over $X$ with a typical fibre $V$ and the elements of the cohomology set $H^{1}\left(X ; G_{X}^{0}\right)$.

The bijection (4.2.5) holds for $G$-bundles with any typical fibre $V$. Two $G$-bundles $(Y, \Psi)$ and $\left(Y^{\prime}, \Psi^{\prime}\right)$ over $X$ with different typical fibres are called associated if the cocycles of transition functions of their atlases $\Psi$ and $\Psi^{\prime}$ are representatives of the same element of the cohomology set $H^{1}\left(X ; G_{X}^{0}\right)$. Then Theorem 4.2.1 can be reformulated as follows.

Theorem 4.2.2: There is one-to-one correspondence between the classes of associated $G$-bundles over $X$ and the elements of the cohomology set $H^{1}\left(X ; G_{X}^{0}\right)$.

Let $f: X^{\prime} \rightarrow X$ be a continuous map. Every continuous $G$-bundle $Y \rightarrow X$ yields the pull-back bundle $f^{*} Y \rightarrow X^{\prime}(1.1 .8)$ with the same structure group $G$. Therefore, $f$ induces the map

$$
[f]: H^{1}\left(X ; G_{X}^{0}\right) \rightarrow H^{1}\left(X^{\prime} ; G_{X^{\prime}}^{0}\right)
$$

Theorem 4.2.3: Given a continuous $G$-bundle $Y$ over a paracompact base $X$, let $f_{1}$ and $f_{2}$ be two continuous maps of $X^{\prime}$ to $X$. If these
maps are homotopic, the pull-back $G$-bundles $f_{1}^{*} Y$ and $f_{2}^{*} Y$ over $X^{\prime}$ are equivalent.

Let us return to smooth fibre bundles. Let $G$, $\operatorname{dim} G>0$, be a real Lie group which acts on a smooth manifold $V$ on the left. A smooth fibre bundle $\pi: Y \rightarrow X$ is called a bundle with a structure group $G$ if it is a continuous $G$-bundle possessing a smooth atlas $\Psi$ (4.2.1) whose transition functions factorize as those (4.2.1) through smooth $G$-valued functions (4.2.3).

Example 4.2.1: Any vector (resp. affine) bundle of fibre dimension $\operatorname{dim} V=m$ is a bundle with a structure group which is the general linear $\operatorname{group} G L(m, \mathbb{R})($ resp. the general affine group $G A(m, \mathbb{R}))$.

Let $G_{X}^{\infty}$ be the sheaf of germs of smooth $G$-valued functions on $X$ and $H^{1}\left(X ; G_{X}^{\infty}\right)$ the first cohomology of a manifold $X$ with coefficients in the sheaf $G_{X}^{\infty}$. The following theorem is analogous to Theorem 4.2.2.

Theorem 4.2.4: There is one-to-one correspondence between the classes of associated smooth $G$-bundles over $X$ and the elements of the cohomology set $H^{1}\left(X ; G_{X}^{\infty}\right)$.

Since a smooth manifold is paracompact, one can show the following.
Theorem 4.2.5: There is a bijection

$$
\begin{equation*}
H^{1}\left(X ; G_{X}^{\infty}\right)=H^{1}\left(X ; G_{X}^{0}\right) \tag{4.2.6}
\end{equation*}
$$

where a Lie group $G$ is treated as a topological group.
The bijection (4.2.6) enables one to classify smooth $G$-bundles as the continuous ones by means of topological invariants.

### 4.3 Principal bundles

We restrict our consideration to smooth bundles with a structure Lie group of non-zero dimension.

Given a real Lie group $G$, let

$$
\begin{equation*}
\pi_{P}: P \rightarrow X \tag{4.3.1}
\end{equation*}
$$

be a $G$-bundle whose typical fibre is the group space of $G$, which a group $G$ acts on by left multiplications. It is called a principal bundle with a structure group $G$. Equivalently, a principal $G$-bundle is defined as a fibre bundle $P$ (4.3.1) which admits an action of $G$ on $P$ on the right by a fibrewise morphism

$$
\begin{align*}
& R_{G P}: G \underset{X}{\times} P \underset{X}{\longrightarrow} P  \tag{4.3.2}\\
& R_{g P}: p \rightarrow p g, \quad \pi_{P}(p)=\pi_{P}(p g), \quad p \in P
\end{align*}
$$

which is free and transitive on each fibre of $P$. As a consequence, the quotient of $P$ with respect to the action (4.3.2) of $G$ is diffeomorphic to a base $X$, i.e., $P / G=X$.

Remark 4.3.1: The definition of a continuous principal bundle is a repetition of that of a smooth one, but all morphisms are continuous.

A principal $G$-bundle $P$ is equipped with a bundle atlas

$$
\begin{equation*}
\Psi_{P}=\left\{\left(U_{\alpha}, \psi_{\alpha}^{P}\right), \varrho_{\alpha \beta}\right\} \tag{4.3.3}
\end{equation*}
$$

whose trivialization morphisms

$$
\psi_{\alpha}^{P}: \pi_{P}^{-1}\left(U_{\alpha}\right) \rightarrow U_{\alpha} \times G
$$

obey the condition

$$
\psi_{\alpha}^{P}(p g)=g \psi_{\alpha}^{P}(p), \quad g \in G
$$

Due to this property, every trivialization morphism $\psi_{\alpha}^{P}$ determines a unique local section $z_{\alpha}: U_{\alpha} \rightarrow P$ such that

$$
\left(\psi_{\alpha}^{P} \circ z_{\alpha}\right)(x)=1, \quad x \in U_{\alpha}
$$

The transformation law for $z_{\alpha}$ reads

$$
\begin{equation*}
z_{\beta}(x)=z_{\alpha}(x) \varrho_{\alpha \beta}(x), \quad x \in U_{\alpha} \cap U_{\beta} \tag{4.3.4}
\end{equation*}
$$

Conversely, the family

$$
\begin{equation*}
\left\{\left(U_{\alpha}, z_{\alpha}\right), \varrho_{\alpha \beta}\right\} \tag{4.3.5}
\end{equation*}
$$

of local sections of $P$ which obey the transformation law (4.3.4) uniquely determines a bundle atlas $\Psi_{P}$ of a principal bundle $P$.

Theorem 4.3.1: A principal bundle admits a global section iff it is trivial.

Example 4.3.2: Let $H$ be a closed subgroup of a real Lie group $G$. Then $H$ is a Lie group. Let $G / H$ be the quotient of $G$ with respect to an action of $H$ on $G$ by right multiplications. Then

$$
\begin{equation*}
\pi_{G H}: G \rightarrow G / H \tag{4.3.6}
\end{equation*}
$$

is a principal $H$-bundle. If $H$ is a maximal compact subgroup of $G$, then $G / H$ is diffeomorphic to $\mathbb{R}^{m}$ and, by virtue of Theorem 1.1.7, $G \rightarrow G / H$ is a trivial bundle, i.e., $G$ is diffeomorphic to the product $\mathbb{R}^{m} \times H$.

Remark 4.3.3: The pull-back $f^{*} P(1.1 .8)$ of a principal bundle also is a principal bundle with the same structure group.

Remark 4.3.4: Let $P \rightarrow X$ and $P^{\prime} \rightarrow X^{\prime}$ be principal $G$ - and $G^{\prime}$ bundles, respectively. A bundle morphism $\Phi: P \rightarrow P^{\prime}$ is a morphism of principal bundles if there exists a Lie group homomorphism $\gamma: G \rightarrow G^{\prime}$ such that

$$
\Phi(p g)=\Phi(p) \gamma(g)
$$

In particular, equivalent principal bundles are isomorphic.
Any class of associated smooth bundles on $X$ with a structure Lie group $G$ contains a principal bundle. In other words, any smooth bundle with a structure Lie group $G$ is associated with some principal bundle.

Let us consider the tangent morphism

$$
\begin{equation*}
T R_{G P}:\left(G \times \mathfrak{g}_{l}\right) \underset{X}{\times} T P \underset{X}{\longrightarrow} T P \tag{4.3.7}
\end{equation*}
$$

to the right action $R_{G P}(4.3 .2)$ of $G$ on $P$. Its restriction to $T_{1} G \times \underset{X}{\times} T P$ provides a homomorphism

$$
\begin{equation*}
\mathfrak{g}_{l} \ni \epsilon \rightarrow \xi_{\epsilon} \in \mathcal{T}(P) \tag{4.3.8}
\end{equation*}
$$

of the left Lie algebra $\mathfrak{g}_{l}$ of $G$ to the Lie algebra $\mathcal{T}(P)$ of vector fields on $P$. Vector fields $\xi_{\epsilon}(4.3 .8)$ are obviously vertical. They are called fundamental vector fields. Given a basis $\left\{\epsilon_{r}\right\}$ for $\mathfrak{g}_{l}$, the corresponding fundamental vector fields $\xi_{r}=\xi_{\epsilon_{r}}$ form a family of $m=\operatorname{dim} \mathfrak{g}_{l}$ nowhere vanishing and linearly independent sections of the vertical tangent bundle $V P$ of $P \rightarrow X$. Consequently, this bundle is trivial

$$
\begin{equation*}
V P=P \times \mathfrak{g}_{l} \tag{4.3.9}
\end{equation*}
$$

by virtue of Theorem 1.2.1.
Restricting the tangent morphism $T R_{G P}$ (4.3.7) to

$$
\begin{equation*}
T R_{G P}: \widehat{0}(G) \underset{X}{\times} T P \underset{X}{\longrightarrow} T P \tag{4.3.10}
\end{equation*}
$$

we obtain the tangent prolongation of the structure group action $R_{G P}$ (4.3.2). If there is no danger of confusion, it is simply called the action of $G$ on $T P$. Since the action of $G(4.3 .2)$ on $P$ is fibrewise, its action (4.3.10) is restricted to the vertical tangent bundle $V P$ of $P$.

Taking the quotient of the tangent bundle $T P \rightarrow P$ and the vertical tangent bundle $V P$ of $P$ by $G(4.3 .10)$, we obtain the vector bundles

$$
\begin{equation*}
T_{G} P=T P / G, \quad V_{G} P=V P / G \tag{4.3.11}
\end{equation*}
$$

over $X$. Sections of $T_{G} P \rightarrow X$ are $G$-invariant vector fields on $P$. Accordingly, sections of $V_{G} P \rightarrow X$ are $G$-invariant vertical vector fields on $P$. Hence, a typical fibre of $V_{G} P \rightarrow X$ is the right Lie algebra $\mathfrak{g}_{r}$ of $G$ subject to the adjoint representation of a structure group $G$. Therefore, $V_{G} P(4.3 .11)$ is called the Lie algebra bundle.

Given a bundle atlas $\Psi_{P}$ (4.3.3) of $P$, there is the corresponding atlas

$$
\begin{equation*}
\Psi=\left\{\left(U_{\alpha}, \psi_{\alpha}\right), \operatorname{Ad}_{\varrho_{\alpha \beta}}\right\} \tag{4.3.12}
\end{equation*}
$$

of a Lie algebra fibre bundle $V_{G} P$, which is provided with bundle coordinates $\left(U_{\alpha} ; x^{\mu}, \chi^{m}\right)$ with respect to the fibre frames

$$
\left\{e_{m}=\psi_{\alpha}^{-1}(x)\left(\varepsilon_{m}\right)\right\}
$$

where $\left\{\varepsilon_{m}\right\}$ is a basis for the Lie algebra $\mathfrak{g}_{r}$. These coordinates obey the transformation rule

$$
\begin{equation*}
\varrho\left(\chi^{m}\right) \varepsilon_{m}=\chi^{m} \mathrm{Ad}_{\varrho^{-1}}\left(\varepsilon_{m}\right) \tag{4.3.13}
\end{equation*}
$$

A glance at this transformation rule shows that $V_{G} P$ is a bundle with a structure group $G$. Moreover, it is associated with a principal $G$-bundle $P$ (see Example 4.7.2).

Accordingly, the vector bundle $T_{G} P(4.3 .11)$ is endowed with bundle coordinates $\left(x^{\mu}, \dot{x}^{\mu}, \chi^{m}\right)$ with respect to the fibre frames $\left\{\partial_{\mu}, e_{m}\right\}$. Their transformation rule is

$$
\begin{equation*}
\varrho\left(\chi^{m}\right) \varepsilon_{m}=\chi^{m} \operatorname{Ad}_{\varrho^{-1}}\left(\varepsilon_{m}\right)+\dot{x}^{\mu} R_{\mu}^{m} \varepsilon_{m} . \tag{4.3.14}
\end{equation*}
$$

If $G$ is a matrix group (see Remark 4.1.2), this transformation rule reads

$$
\begin{equation*}
\varrho\left(\chi^{m}\right) \varepsilon_{m}=\chi^{m} \varrho^{-1} \varepsilon_{m} \varrho-\dot{x}^{\mu} \partial_{\mu}\left(\varrho^{-1}\right) \varrho . \tag{4.3.15}
\end{equation*}
$$

Since the second term in the right-hand sides of expressions (4.3.14) (4.3.15) depend on derivatives of a $G$-valued function $\varrho$ on $X$, the vector bundle $T_{G} P$ (4.3.11) fails to be a $G$-bundle.

The Lie bracket of $G$-invariant vector fields on $P$ goes to the quotient by $G$ and defines the Lie bracket of sections of the vector bundle $T_{G} P \rightarrow$ $X$. This bracket reads

$$
\begin{align*}
& \xi=\xi^{\lambda} \partial_{\lambda}+\xi^{p} e_{p}, \quad \eta=\eta^{\mu} \partial_{\mu}+\eta^{q} e_{q}  \tag{4.3.16}\\
& {[\xi, \eta]=\left(\xi^{\mu} \partial_{\mu} \eta^{\lambda}-\eta^{\mu} \partial_{\mu} \xi^{\lambda}\right) \partial_{\lambda}+}  \tag{4.3.17}\\
& \quad\left(\xi^{\lambda} \partial_{\lambda} \eta^{r}-\eta^{\lambda} \partial_{\lambda} \xi^{r}+c_{p q}^{r} \xi^{p} \eta^{q}\right) e_{r} .
\end{align*}
$$

Putting $\xi^{\lambda}=0$ and $\eta^{\mu}=0$ in the formulas (4.3.16) - (4.3.17), we obtain the Lie bracket

$$
\begin{equation*}
[\xi, \eta]=c_{p q}^{r} \xi^{p} \eta^{q} e_{r} \tag{4.3.18}
\end{equation*}
$$

of sections of the Lie algebra bundle $V_{G} P$. A glance at the expression (4.3.18) shows that sections of $V_{G} P$ form a finite-dimensional Lie $C^{\infty}(X)$ algebra, called the gauge algebra. Therefore, $V_{G} P$ also is called the gauge algebra bundle.

### 4.4 Principal connections

Principal connections on a principal bundle $P$ (4.3.1) are connections on $P$ which are equivariant with respect to the right action (4.3.2) of a structure group $G$ on $P$. In order to describe them, we follow the definition of connections on a fibre bundle $Y \rightarrow X$ as global sections of the affine jet bundle $J^{1} Y \rightarrow X$ (Theorem 3.2.1).

Let $J^{1} P$ be the first order jet manifold of a principal $G$-bundle $P \rightarrow X$ (4.3.1). Then connections on a principal bundle $P \rightarrow X$ are global sections

$$
\begin{equation*}
A: P \rightarrow J^{1} P \tag{4.4.1}
\end{equation*}
$$

of the affine jet bundle $J^{1} P \rightarrow P$ modelled over the vector bundle

$$
T^{*} X \underset{P}{\otimes} V P=\left(T^{*} X \underset{P}{\otimes} \mathfrak{g}_{l}\right)
$$

In order to define principal connections on $P \rightarrow X$, let us consider the jet prolongation

$$
J^{1} R_{G}: J^{1}(X \times G) \underset{X}{\times} J^{1} P \rightarrow J^{1} P
$$

of the morphism $R_{G P}$ (4.3.2). Restricting this morphism to

$$
J^{1} R_{G}: \widehat{0}(G) \underset{X}{\times} J^{1} P \rightarrow J^{1} P
$$

we obtain the jet prolongation of the structure group action $R_{G P}$ (4.3.2) called, simply, the action of $G$ on $J^{1} P$. It reads

$$
\begin{equation*}
G \ni g: j_{x}^{1} p \rightarrow\left(j_{x}^{1} p\right) g=j_{x}^{1}(p g) . \tag{4.4.2}
\end{equation*}
$$

Taking the quotient of the affine jet bundle $J^{1} P$ by $G$ (4.4.2), we obtain the affine bundle

$$
\begin{equation*}
C=J^{1} P / G \rightarrow X \tag{4.4.3}
\end{equation*}
$$

modelled over the vector bundle

$$
\bar{C}=T^{*} X \underset{X}{\otimes} V_{G} P \rightarrow X
$$

Hence, there is the vertical splitting

$$
V C=C \underset{X}{\otimes} \bar{C}
$$

of the vertical tangent bundle $V C$ of $C \rightarrow X$.
Remark 4.4.1: A glance at the expression (4.4.2) shows that the fibre bundle $J^{1} P \rightarrow C$ is a principal bundle with the structure group $G$. It is canonically isomorphic to the pull-back

$$
\begin{equation*}
J^{1} P=P_{C}=C \underset{X}{\times} P \rightarrow C \tag{4.4.4}
\end{equation*}
$$

Taking the quotient with respect to the action of a structure group $G$, one can reduce the canonical imbedding (2.1.5) (where $Y=P$ ) to the bundle monomorphism

$$
\begin{align*}
& \lambda_{C}: C \underset{X}{\longrightarrow} T^{*} X \underset{X}{\otimes} T_{G} P, \\
& \lambda_{C}: d x^{\mu} \otimes\left(\partial_{\mu}+\chi_{\mu}^{m} e_{m}\right) . \tag{4.4.5}
\end{align*}
$$

It follows that, given atlases $\Psi_{P}(4.3 .3)$ of $P$ and $\Psi(4.3 .12)$ of $T_{G} P$, the bundle of principal connections $C$ is provided with bundle coordinates $\left(x^{\lambda}, a_{\mu}^{m}\right)$ possessing the transformation rule

$$
\begin{equation*}
\varrho\left(a_{\mu}^{m}\right) \varepsilon_{m}=\left(a_{\nu}^{m} \operatorname{Ad}_{\varrho^{-1}}\left(\varepsilon_{m}\right)+R_{\nu}^{m} \varepsilon_{m}\right) \frac{\partial x^{\nu}}{\partial x^{\prime \mu}} \tag{4.4.6}
\end{equation*}
$$

If $G$ is a matrix group, this transformation rule reads

$$
\begin{equation*}
\varrho\left(a_{\mu}^{m}\right) \varepsilon_{m}=\left(a_{\nu}^{m} \varrho^{-1}\left(\varepsilon_{m}\right) \varrho-\partial_{\mu}\left(\varrho^{-1}\right) \varrho\right) \frac{\partial x^{\nu}}{\partial x^{\prime \mu}} \tag{4.4.7}
\end{equation*}
$$

A glance at this expression shows that the bundle of principal connections $C$ as like as the vector bundle $T_{G} P(4.3 .11)$ fails to be a bundle with a structure group $G$.

As was mentioned above, a connection $A$ (4.4.1) on a principal bundle $P \rightarrow X$ is called a principal connection if it is equivariant under the action (4.4.2) of a structure group $G$, i.e.,

$$
\begin{equation*}
A(p g)=A(p) g \quad g \in G . \tag{4.4.8}
\end{equation*}
$$

There is obvious one-to-one correspondence between the principal connections on a principal $G$-bundle $P$ and global sections

$$
\begin{equation*}
A: X \rightarrow C \tag{4.4.9}
\end{equation*}
$$

of the bundle $C \rightarrow X$ (4.4.3), called the bundle of principal connections.

Theorem 4.4.1: Since the bundle of principal connections $C \rightarrow X$ is affine, principal connections on a principal bundle always exist.

Due to the bundle monomorphism (4.4.5), any principal connection $A$ (4.4.9) is represented by a $T_{G} P$-valued form

$$
\begin{equation*}
A=d x^{\lambda} \otimes\left(\partial_{\lambda}+A_{\lambda}^{q} e_{q}\right) \tag{4.4.10}
\end{equation*}
$$

Taking the quotient with respect to the action of a structure group $G$, one can reduce the exact sequence (1.2.10) (where $Y=P$ ) to the exact sequence

$$
\begin{equation*}
0 \rightarrow V_{G} P \underset{X}{\longrightarrow} T_{G} P \longrightarrow T X \rightarrow 0 \tag{4.4.11}
\end{equation*}
$$

A principal connection $A(4.4 .10)$ defines a splitting of this exact sequence.

Remark 4.4.2: A principal connection $A$ (4.4.1) on a principal bundle $P \rightarrow X$ can be represented by the vertical-valued form $A$ (3.1.9) on $P$ which is a $\mathfrak{g l}$-valued form due to the trivialization (4.3.9). It is the
familiar $\mathfrak{g l}$-valued connection form on a principal bundle $P$. Given a local bundle splitting $\left(U_{\alpha}, z_{\alpha}\right)$ of $P$, this form reads

$$
\begin{equation*}
\bar{A}=\psi_{\alpha}^{*}\left(\theta_{l}-\bar{A}_{\lambda}^{q} d x^{\lambda} \otimes \epsilon_{q}\right), \tag{4.4.12}
\end{equation*}
$$

where $\theta_{l}$ is the canonical $\mathfrak{g l}$-valued one-form (4.1.8) on $G$ and $A_{\lambda}^{p}$ are local functions on $P$ such that

$$
\bar{A}_{\lambda}^{q}(p g) \epsilon_{q}=\bar{A}_{\lambda}^{q}(p) \operatorname{Ad}_{g^{-1}}\left(\epsilon_{q}\right) .
$$

The pull-back $z_{\alpha}^{*} \bar{A}$ of the connection form $\bar{A}$ (4.4.12) onto $U_{\alpha}$ is the well-known local connection one-form

$$
\begin{equation*}
A_{\alpha}=-A_{\lambda}^{q} d x^{\lambda} \otimes \epsilon_{q}=A_{\lambda}^{q} d x^{\lambda} \otimes \varepsilon_{q}, \tag{4.4.13}
\end{equation*}
$$

where $A_{\lambda}^{q}=\bar{A}_{\lambda}^{q} \circ z_{\alpha}$ are local functions on $X$. It is readily observed that the coefficients $A_{\lambda}^{q}$ of this form are exactly the coefficients of the form (4.4.10).

There are both pull-back and push-forward operations of principal connections.

Theorem 4.4.2: Let $P$ be a principal bundle and $f^{*} P(1.1 .8)$ the pullback principal bundle with the same structure group. Let $f_{P}$ be the canonical morphism (1.1.9) of $f^{*} P$ to $P$. If $A$ is a principal connection on $P$, then the pull-back connection $f^{*} A(3.1 .12)$ on $f^{*} P$ is a principal connection.

Theorem 4.4.3: Let $P^{\prime} \rightarrow X$ and $P \rightarrow X$ be principle bundles with structure groups $G^{\prime}$ and $G$, respectively. Let $\Phi: P^{\prime} \rightarrow P$ be a principal bundle morphism over $X$ with the corresponding homomorphism $G^{\prime} \rightarrow$ $G$ (see Remark 4.3.4). For every principal connection $A^{\prime}$ on $P^{\prime}$, there exists a unique principal connection $A$ on $P$ such that $T \Phi$ sends the horizontal subspaces of $T P^{\prime} A^{\prime}$ onto the horizontal subspaces of $T P$ with respect to $A$.

Let $P \rightarrow X$ be a principal $G$-bundle. The Frölicher-Nijenhuis bracket (1.4.8) on the space $\mathcal{O}^{*}(P) \otimes \mathcal{T}(P)$ of tangent-valued forms on $P$ is
compatible with the right action $R_{G P}$ (4.3.2). Therefore, it induces the Frölicher-Nijenhuis bracket on the space $\mathcal{O}^{*}(X) \otimes T_{G} P(X)$ of $T_{G} P$-valued forms on $X$, where $T_{G} P(X)$ is the vector space of sections of the vector bundle $T_{G} P \rightarrow X$. Note that, as it follows from the exact sequence (4.4.11), there is an epimorphism

$$
T_{G} P(X) \rightarrow \mathcal{T}(X)
$$

Let $A \in \mathcal{O}^{1}(X) \otimes T_{G} P(X)$ be a principal connection (4.4.10). The associated Nijenhuis differential is

$$
\begin{align*}
& d_{A}: \mathcal{O}^{r}(X) \otimes T_{G} P(X) \rightarrow \mathcal{O}^{r+1}(X) \otimes V_{G} P(X) \\
& d_{A} \phi=[A, \phi]_{\mathrm{FN}}, \quad \phi \in \mathcal{O}^{r}(X) \otimes T_{G} P(X) \tag{4.4.14}
\end{align*}
$$

The strength of a principal connection $A(4.4 .10)$ is defined as the $V_{G} P-$ valued two-form

$$
\begin{equation*}
F_{A}=\frac{1}{2} d_{A} A=\frac{1}{2}[A, A]_{\mathrm{FN}} \in \mathcal{O}^{2}(X) \otimes V_{G} P(X) \tag{4.4.15}
\end{equation*}
$$

Its coordinated expression

$$
\begin{align*}
F_{A}= & \frac{1}{2} F_{\lambda \mu}^{r} d x^{\lambda} \wedge d x^{\mu} \otimes e_{r}, \\
F_{\lambda \mu}^{r}= & {\left[\partial_{\lambda}+A_{\lambda}^{p} e_{p}, \partial_{\mu}+A_{\mu}^{q} e_{q}\right]^{r}=}  \tag{4.4.16}\\
& \partial_{\lambda} A_{\mu}^{r}-\partial_{\mu} A_{\lambda}^{r}+c_{p q}^{r} A_{\lambda}^{p} A_{\mu}^{q},
\end{align*}
$$

results from the bracket (4.3.17).
Remark 4.4.3: It should be emphasized that the strength $F_{A}$ (4.4.15) is not the standard curvature (3.3.1) of a principal connection because $A$ (4.4.10) is not a tangent-valued form. The curvature of a principal connection $A$ (4.4.1) on $P$ is the $V P$-valued two-form $R$ (3.3.1) on $P$, which is brought into the $\mathfrak{g}_{l}$-valued form owing to the canonical isomorphism (4.3.9).

Remark 4.4.4: Given a principal connection $A$ (4.4.9), let $\Phi_{C}$ be a vertical principal automorphism of the bundle of principal connections $C$.

The connection $A^{\prime}=\Phi_{C} \circ A$ is called conjugate to a principal connection $A$. The strength forms (4.4.15) of conjugate principal connections $A$ and $A^{\prime}$ coincide with each other, i.e., $F_{A}=F_{A^{\prime}}$.

### 4.5 Canonical principal connection

Given a principal $G$-bundle $P \rightarrow X$ and its jet manifold $J^{1} P$, let us consider the canonical morphism $\theta_{(1)}(2.1 .5)$ where $Y=P$. By virtue of Remark 1.2.2, this morphism defines the morphism

$$
\theta: J^{1} P \underset{P}{\times} T P \rightarrow V P
$$

Taking its quotient with respect to $G$, we obtain the morphism

$$
\begin{align*}
& C \underset{X}{\times} T_{G} P \xrightarrow{\theta} V_{G} P,  \tag{4.5.1}\\
& \theta\left(\partial_{\lambda}\right)=-a_{\lambda}^{p} e_{p}, \quad \theta\left(e_{p}\right)=e_{p} .
\end{align*}
$$

Consequently, the exact sequence (4.4.11) admits the canonical splitting over $C$.

In view of this fact, let us consider the pull-back principal $G$-bundle $P_{C}$ (4.4.4). Since

$$
\begin{equation*}
V_{G}(C \underset{X}{\times} P)=C \underset{X}{\times} V_{G} P, \quad T_{G}(C \underset{X}{\times} P)=T C \underset{X}{\times} T_{G} P, \tag{4.5.2}
\end{equation*}
$$

the exact sequence (4.4.11) for the principal bundle $P_{C}$ reads

$$
\begin{equation*}
0 \rightarrow C \underset{X}{\times} V_{G} P \underset{C}{\longrightarrow} T C \underset{X}{\times} T_{G} P \longrightarrow T C \rightarrow 0 . \tag{4.5.3}
\end{equation*}
$$

The morphism (4.5.1) yields the horizontal splitting (3.1.3):

$$
T C \underset{X}{\times} T_{G} P \longrightarrow C \underset{X}{\times} T_{G} P \longrightarrow C \underset{X}{\times} V_{G} P,
$$

of the exact sequence (4.5.3). Thus, it defines the principal connection

$$
\begin{align*}
& \mathcal{A}: T C \rightarrow T C \underset{X}{\times} T_{G} P, \\
& \mathcal{A}=d x^{\lambda} \otimes\left(\partial_{\lambda}+a_{\lambda}^{p} e_{p}\right)+d a_{\lambda}^{r} \otimes \partial_{r}^{\lambda},  \tag{4.5.4}\\
& \mathcal{A} \in \mathcal{O}^{1}(C) \otimes T_{G}\left(C_{X}^{\times} P\right)(X),
\end{align*}
$$

on the principal bundle

$$
\begin{equation*}
P_{C}=C \underset{X}{\times} P \rightarrow C \tag{4.5.5}
\end{equation*}
$$

It follows that the principal bundle $P_{C}$ admits the canonical principal connection (4.5.4).

Following the expression (4.4.15), let us define the strength

$$
\begin{align*}
& F_{\mathcal{A}}=\frac{1}{2} d_{\mathcal{A}} \mathcal{A}=\frac{1}{2}[\mathcal{A}, \mathcal{A}] \in \mathcal{O}^{2}(C) \otimes V_{G} P(X) \\
& F_{\mathcal{A}}=\left(d a_{\mu}^{r} \wedge d x^{\mu}+\frac{1}{2} c_{p q}^{r} a_{\lambda}^{p} a_{\mu}^{q} d x^{\lambda} \wedge d x^{\mu}\right) \otimes e_{r} \tag{4.5.6}
\end{align*}
$$

of the canonical principal connection $\mathcal{A}$ (4.5.4). It is called the canonical strength because, given a principal connection $A$ (4.4.9) on a principal bundle $P \rightarrow X$, the pull-back

$$
\begin{equation*}
A^{*} F_{\mathcal{A}}=F_{A} \tag{4.5.7}
\end{equation*}
$$

is the strength (4.4.16) of $A$.
With the $V_{G} P$-valued two-form $F_{\mathcal{A}}(4.5 .6)$ on $C$, let us define the $V_{G} P$-valued horizontal two-form

$$
\begin{align*}
& \mathcal{F}=h_{0}\left(F_{\mathcal{A}}\right)=\frac{1}{2} \mathcal{F}_{\lambda \mu}^{r} d x^{\lambda} \wedge d x^{\mu} \otimes \varepsilon_{r} \\
& \mathcal{F}_{\lambda \mu}^{r}=a_{\lambda \mu}^{r}-a_{\mu \lambda}^{r}+c_{p q}^{r} a_{\lambda}^{p} a_{\mu}^{q} \tag{4.5.8}
\end{align*}
$$

on $J^{1} C$. It is called the strength form. For each principal connection $A$ (4.4.9) on $P$, the pull-back

$$
\begin{equation*}
J^{1} A^{*} \mathcal{F}=F_{A} \tag{4.5.9}
\end{equation*}
$$

is the strength (4.4.16) of $A$.
The strength form (4.5.8) yields an affine surjection

$$
\begin{equation*}
\mathcal{F} / 2: J^{1} C \underset{C}{\longrightarrow} C \underset{X}{\times}\left(\stackrel{2}{\wedge} T^{*} X \otimes V_{G} P\right) \tag{4.5.10}
\end{equation*}
$$

over $C$ of the affine jet bundle $J^{1} C \rightarrow C$ to the vector (and, consequently, affine) bundle

$$
C \underset{X}{\times}\left(2^{2} T^{*} X \otimes V_{G} P\right) \rightarrow C .
$$

By virtue of Theorem 1.1.10, its kernel $C_{+}=\operatorname{Ker} \mathcal{F} / 2$ is an affine subbundle of $J^{1} C \rightarrow C$. Thus, we have the canonical splitting of the affine jet bundle

$$
\begin{align*}
J^{1} C & =C_{+} \underset{C}{\oplus} C_{-}=C_{+} \underset{C}{\oplus}\left(C \underset{X}{\times} \stackrel{2}{\wedge} T^{*} X \otimes V_{G} P\right)  \tag{4.5.11}\\
a_{\lambda \mu}^{r}= & \frac{1}{2}\left(\mathcal{F}_{\lambda \mu}^{r}+\mathcal{S}_{\lambda \mu}^{r}\right)=\frac{1}{2}\left(a_{\lambda \mu}^{r}+a_{\mu \lambda}^{r}-c_{p q}^{r} a_{\lambda}^{p} a_{\mu}^{q}\right)+  \tag{4.5.12}\\
& \frac{1}{2}\left(a_{\lambda \mu}^{r}-a_{\mu \lambda}^{r}+c_{p q}^{r} a_{\lambda}^{p} a_{\mu}^{q}\right) .
\end{align*}
$$

The jet manifold $J^{1} C$ plays a role of the configuration space of classical gauge theory on principal bundles.

### 4.6 Gauge transformations

In classical gauge theory, gauge transformations are defined as principal automorphisms of a principal bundle $P$. In accordance with Remark 4.3.4, an automorphism $\Phi_{P}$ of a principal $G$-bundle $P$ is called principal if it is equivariant under the right action (4.3.2) of a structure group $G$ on $P$, i.e.,

$$
\begin{equation*}
\Phi_{P}(p g)=\Phi_{P}(p) g, \quad g \in G, \quad p \in P \tag{4.6.1}
\end{equation*}
$$

In particular, every vertical principal automorphism of a principal bundle $P$ is represented as

$$
\begin{equation*}
\Phi_{P}(p)=p f(p), \quad p \in P, \tag{4.6.2}
\end{equation*}
$$

where $f$ is a $G$-valued equivariant function on $P$, i.e.,

$$
\begin{equation*}
f(p g)=g^{-1} f(p) g, \quad g \in G . \tag{4.6.3}
\end{equation*}
$$

There is one-to-one correspondence between the equivariant functions $f$ (4.6.3) and the global sections $s$ of the associated group bundle

$$
\begin{equation*}
\pi_{P^{G}}: P^{G} \rightarrow X \tag{4.6.4}
\end{equation*}
$$

whose fibres are groups isomorphic to $G$ and whose typical fibre is the group $G$ which acts on itself by the adjoint representation. This one-toone correspondence is defined by the relation

$$
\begin{equation*}
s\left(\pi_{P}(p)\right) p=p f(p), \quad p \in P \tag{4.6.5}
\end{equation*}
$$

(see Example 4.7.3). The group of vertical principal automorphisms of a principal $G$-bundle is called the gauge group. It is isomorphic to the group $P^{G}(X)$ of global sections of the group bundle (4.6.4). Its unit element is the canonical global section $\hat{1}$ of $P^{G} \rightarrow X$ whose values are unit elements of fibres of $P^{G}$.

Remark 4.6.1: Note that transition functions of atlases of a principle bundle $P$ also are represented by local sections of the associated group bundle $P^{G}$ (4.6.4).

Let us consider (local) one-parameter groups of principal automorphisms of $P$. Their infinitesimal generators are $G$-invariant projectable vector fields $\xi$ on $P$, and vice versa. We call $\xi$ the principal vector fields or the infinitesimal gauge transformations. They are represented by sections $\xi(4.3 .16)$ of the vector bundle $T_{G} P$ (4.3.11). Principal vector fields constitute a real Lie algebra $T_{G} P(X)$ with respect to the Lie bracket (4.3.17). Vertical principal vector fields are the sections

$$
\begin{equation*}
\xi=\xi^{p} e_{p} \tag{4.6.6}
\end{equation*}
$$

of the gauge algebra bundle $V_{G} P \rightarrow X$ (4.3.11). They form a finitedimensional Lie $C^{\infty}(X)$-algebra $\mathcal{G}(X)=V_{G} P(X)$ (4.3.18) that has been called the gauge algebra.

Any (local) one-parameter group of principal automorphism $\Phi_{P}$ (4.6.1) of a principal bundle $P$ admits the jet prolongation $J^{1} \Phi_{P}(2.1 .7)$ to a one-parameter group of $G$-equivariant automorphism of the jet manifold $J^{1} P$ which, in turn, yields a one-parameter group of principal automorphisms $\Phi_{C}$ of the bundle of principal connections $C$ (4.4.3). Its infinitesimal generator is a vector field on $C$, called the principal vector field on
$C$ and regarded as an infinitesimal gauge transformation of $C$. Thus, any principal vector field $\xi(4.3 .16)$ on $P$ yields a principal vector field $u_{\xi}$ on $C$, which can be defined as follows.

Using the morphism (4.5.1), we obtain the morphism

$$
\xi\rfloor \theta: C \underset{X}{\longrightarrow} V_{G} P
$$

which is a section of of the Lie algebra bundle

$$
V_{G}(C \underset{X}{\times} P) \rightarrow C
$$

in accordance with the first formula (4.5.2). Then the equation

$$
\left.\left.u_{\xi}\right\rfloor F_{\mathcal{A}}=d_{\mathcal{A}}(\xi\rfloor \theta\right)
$$

uniquely determines a desired vector field $u_{\xi}$ on $C$. A direct computation leads to

$$
\begin{equation*}
u_{\xi}=\xi^{\mu} \partial_{\mu}+\left(\partial_{\mu} \xi^{r}+c_{p q}^{r} a_{\mu}^{p} \xi^{q}-a_{\nu}^{r} \partial_{\mu} \xi^{\nu}\right) \partial_{r}^{\mu} . \tag{4.6.7}
\end{equation*}
$$

In particular, if $\xi$ is a vertical principal field (4.6.6), we obtain the vertical vector field

$$
\begin{equation*}
u_{\xi}=\left(\partial_{\mu} \xi^{r}+c_{p q}^{r} a_{\mu}^{p} \xi^{q}\right) \partial_{r}^{\mu} \tag{4.6.8}
\end{equation*}
$$

Remark 4.6.2: The jet prolongation (2.1.8) of the vector field $u_{\xi}$ (4.6.7) onto $J^{1} C$ reads

$$
\begin{gather*}
J^{1} u_{\xi}=u_{\xi}+\left(\partial_{\lambda \mu} \xi^{r}+c_{p q}^{r} a_{\mu}^{p} \partial_{\lambda} \xi^{q}+c_{p q}^{r} a_{\lambda \mu}^{p} \xi^{q}-\right.  \tag{4.6.9}\\
\left.a_{\nu}^{r} \partial_{\lambda \mu} \xi^{\nu}-a_{\lambda \nu}^{r} \partial_{\mu} \xi^{\nu}-a_{\nu \mu}^{r} \partial_{\lambda} \xi^{\nu}\right) \partial_{r}^{\lambda \mu} .
\end{gather*}
$$

Example 4.6.3: Let $A(4.4 .10)$ be a principal connection on $P$. For any vector field $\tau$ on $X$, this connection yields a section

$$
\tau\rfloor A=\tau^{\lambda} \partial_{\lambda}+A_{\lambda}^{p} \tau^{\lambda} e_{p}
$$

of the vector bundle $T_{G} P \rightarrow X$. It, in turn, defines a principal vector field (4.6.7) on the bundle of principal connection $C$ which reads

$$
\begin{align*}
& \tau_{A}=\tau^{\lambda} \partial_{\lambda}+\left(\partial_{\mu}\left(A_{\nu}^{r} \tau^{\nu}\right)+c_{p q}^{r} a_{\mu}^{p} A_{\nu}^{q} \tau^{\nu}-a_{\nu}^{r} \partial_{\mu} \tau^{\nu}\right) \partial_{r}^{\mu}  \tag{4.6.10}\\
& \xi^{\lambda}=\tau^{\lambda}, \quad \xi^{p}=A_{\nu}^{p} \tau^{\nu}
\end{align*}
$$

It is readily justified that the monomorphism

$$
\begin{equation*}
T_{G} P(X) \ni \xi \rightarrow u_{\xi} \in \mathcal{T}(C) \tag{4.6.11}
\end{equation*}
$$

obeys the equality

$$
\begin{equation*}
u_{[\xi, \eta]}=\left[u_{\xi}, u_{\eta}\right] \tag{4.6.12}
\end{equation*}
$$

i.e., it is a monomorphism of the real Lie algebra $T_{G} P(X)$ to the real Lie algebra $\mathcal{T}(C)$. In particular, the image of the gauge algebra $\mathcal{G}(X)$ in $\mathcal{T}(C)$ also is a real Lie algebra, but not the $C^{\infty}(X)$-one because

$$
u_{f \xi} \neq f u_{\xi}, \quad f \in C^{\infty}(X)
$$

Remark 4.6.4: A glance at the expression (4.6.7) shows that the monomorphism (4.6.11) is a linear first order differential operator which sends sections of the pull-back bundle

$$
C \underset{X}{\times} T_{G} P \rightarrow C
$$

onto sections of the tangent bundle $T C \rightarrow C$. Refereing to Definition 7.2.10, we therefore can treat principal vector fields (4.6.7) as infinitesimal gauge transformations depending on gauge parameters $\xi \in T_{G} P(X)$.

### 4.7 Geometry of associated bundles

Given a principal $G$-bundle $P(4.3 .1)$, any associated $G$-bundle over $X$ with a typical fibre $V$ is equivalent to the following one.

Let us consider the quotient

$$
\begin{equation*}
Y=(P \times V) / G \tag{4.7.1}
\end{equation*}
$$

of the product $P \times V$ by identification of elements $(p, v)$ and $\left(p g, g^{-1} v\right)$ for all $g \in G$. Let $[p]$ denote the restriction of the canonical surjection

$$
\begin{equation*}
P \times V \rightarrow(P \times V) / G \tag{4.7.2}
\end{equation*}
$$

to the subset $\{p\} \times V$ so that

$$
[p](v)=[p g]\left(g^{-1} v\right)
$$

Then the map

$$
Y \ni[p](V) \rightarrow \pi_{P}(p) \in X
$$

makes the quotient $Y$ (4.7.1) into a fibre bundle over $X$. This is a smooth $G$-bundle with the typical fibre $V$ which is associated with the principal $G$-bundle $P$. For short, we call it the $P$-associated bundle.

Remark 4.7.1: The tangent morphism to the morphism (4.7.2) and the jet prolongation of the morphism (4.7.2) lead to the bundle isomorphisms

$$
\begin{align*}
& T Y=(T P \times T V) / G  \tag{4.7.3}\\
& J^{1} Y=\left(J^{1} P \times V\right) / G \tag{4.7.4}
\end{align*}
$$

The peculiarity of the $P$-associated bundle $Y(4.7 .1)$ is the following.
(i) Every bundle atlas $\Psi_{P}=\left\{\left(U_{\alpha}, z_{\alpha}\right)\right\}$ (4.3.5) of $P$ defines a unique associated bundle atlas

$$
\begin{equation*}
\Psi=\left\{\left(U_{\alpha}, \psi_{\alpha}(x)=\left[z_{\alpha}(x)\right]^{-1}\right)\right\} \tag{4.7.5}
\end{equation*}
$$

of the quotient $Y$ (4.7.1).
Example 4.7.2: Because of the splitting (4.3.9), the Lie algebra bundle

$$
V_{G} P=\left(P \times \mathfrak{g}_{l}\right) / G
$$

by definition, is of the form (4.7.1). Therefore, it is a $P$-associated bundle.

Example 4.7.3: The group bundle $\bar{P}(4.6 .4)$ is defined as the quotient

$$
\begin{equation*}
P^{G}=(P \times G) / G, \tag{4.7.6}
\end{equation*}
$$

where the group $G$ which acts on itself by the adjoint representation. There is the following fibre-to-fibre action of the group bundle $P^{G}$ on any $P$-associated bundle $Y$ (4.7.1):

$$
\begin{aligned}
& P^{G} \times Y \underset{X}{\longrightarrow} Y \\
& ((p, g) / G,(p, v) / G) \rightarrow(p, g v) / G, \quad g \in G, \quad v \in V
\end{aligned}
$$

For instance, the action of $P^{G}$ on $P$ in the formula (4.6.5) is of this type.
(ii) Any principal automorphism $\Phi_{P}$ (4.6.1) of $P$ yields a unique principal automorphism

$$
\begin{equation*}
\Phi_{Y}:(p, v) / G \rightarrow\left(\Phi_{P}(p), v\right) / G, \quad p \in P, \quad v \in V \tag{4.7.7}
\end{equation*}
$$

of the $P$-associated bundle $Y$ (4.7.1). For the sake of brevity, we agree to write

$$
\Phi_{Y}:(P \times V) / G \rightarrow\left(\Phi_{P}(P) \times V\right) / G
$$

Accordingly, any (local) one-parameter group of principal automorphisms of $P$ induces a (local) one-parameter group of automorphisms of the $P$ associated bundle $Y$ (4.7.1). Passing to infinitesimal generators of these groups, we obtain that any principal vector field $\xi(4.3 .16)$ yields a vector field $v_{\xi}$ on $Y$ regarded as an infinitesimal gauge transformation of $Y$. Owing to the bundle isomorphism (4.7.3), we have

$$
\begin{align*}
& v_{\xi}: X \rightarrow(\xi(P) \times T V) / G \subset T Y \\
& v_{\xi}=\xi^{\lambda} \partial_{\lambda}+\xi^{p} I_{p}^{i} \partial_{i} \tag{4.7.8}
\end{align*}
$$

where $\left\{I_{p}\right\}$ is a representation of the Lie algebra $\mathfrak{g}_{r}$ of $G$ in $V$.
(iii) Any principal connection on $P \rightarrow X$ defines a unique connection on the $P$-associated fibre bundle $Y(4.7 .1)$ as follows. Given a principal connection $A(4.4 .8)$ on $P$ and the corresponding horizontal distribution $H P \subset T P$, the tangent map to the canonical morphism (4.7.2) defines the horizontal splitting of the tangent bundle $T Y$ of $Y(4.7 .1)$ and the corresponding connection on $Y \rightarrow X$. Owing to the bundle isomorphism (4.7.4), we have

$$
\begin{align*}
& A:(P \times V) / G \rightarrow(A(P) \times V) / G \subset J^{1} Y \\
& A=d x^{\lambda} \otimes\left(\partial_{\lambda}+A_{\lambda}^{p} I_{p}^{i} \partial_{i}\right) \tag{4.7.9}
\end{align*}
$$

where $\left\{I_{p}\right\}$ is a representation of the Lie algebra $\mathfrak{g}_{r}$ of $G$ in $V$. The connection $A$ (4.7.9) on $Y$ is called the associated principal connection or, simply, a principal connection on $Y \rightarrow X$. The curvature (3.3.2) of this connection takes the form

$$
\begin{equation*}
R=\frac{1}{2} F_{\lambda \mu}^{p} I_{p}^{i} d x^{\lambda} \wedge d x^{\mu} \otimes \partial_{i} \tag{4.7.10}
\end{equation*}
$$

Example 4.7.4: A principal connection $A$ on $P$ yields the associated connection (4.7.9) on the associated Lie algebra bundle $V_{G} P$ which reads

$$
\begin{equation*}
A=d x^{\lambda} \otimes\left(\partial_{\lambda}-c_{p q}^{m} A_{\lambda}^{p} \xi^{q} e_{m}\right) \tag{4.7.11}
\end{equation*}
$$

Remark 4.7.5: If an associated principal connection $A$ is linear, one can define its strength

$$
\begin{equation*}
F_{A}=\frac{1}{2} F_{\lambda \mu}^{p} I_{p} d x^{\lambda} \wedge d x^{\mu} \tag{4.7.12}
\end{equation*}
$$

where $I_{p}$ are matrices of a representation of the Lie algebra $\mathfrak{g}_{r}$ in fibres of $Y$ with respect to the fibre bases $\left\{e_{i}(x)\right\}$. They coincide with the matrices of a representation of $\mathfrak{g}_{r}$ in the typical fibre $V$ of $Y$ with respect to its fixed basis $\left\{e_{i}\right\}$ (see the relation (1.2.1)). It follows that $G$-valued transition functions act on $I_{p}$ by the adjoint representation. Note that, because of the canonical splitting (1.2.8), one can identify $e_{i}(x)=\partial_{i}$.

In view of the above mentioned properties, the $P$-associated bundle $Y(4.7 .1)$ is called canonically associated to a principal bundle $P$. Unless otherwise stated, only canonically associated bundles are considered, and we simply call $Y$ (4.7.1) an associated bundle.

### 4.8 Reduced structure

Let $H$ and $G$ be Lie groups and $\phi: H \rightarrow G$ a Lie group homomorphism. If $P_{H} \rightarrow X$ is a principal $H$-bundle, there always exists a principal $G$ bundle $P_{G} \rightarrow X$ together with the principal bundle morphism

$$
\begin{equation*}
\Phi: P_{H} \xrightarrow[X]{\longrightarrow} P_{G} \tag{4.8.1}
\end{equation*}
$$

over $X$ (see Remark 4.3.4). It is the $P_{H}$-associated bundle

$$
P_{G}=\left(P_{H} \times G\right) / H
$$

with the typical fibre $G$ on which $H$ acts on the left by the rule $h(g)=$ $\phi(h) g$, while $G$ acts on $P_{G}$ as

$$
G \ni g^{\prime}:(p, g) / H \rightarrow\left(p, g g^{\prime}\right) / H
$$

Conversely, if $P_{G} \rightarrow X$ is a principal $G$-bundle, a problem is to find a principal $H$-bundle $P_{H} \rightarrow X$ together with a principal bundle morphism (4.8.1). If $H \rightarrow G$ is a closed subgroup, we have the structure group reduction. If $H \rightarrow G$ is a group epimorphism, one says that $P_{G}$ lifts to $P_{H}$.

Here, we restrict our consideration to the reduction problem. In this case, the bundle monomorphism (4.8.1) is called a reduced $H$-structure.

Let $P$ (4.3.1) be a principal $G$-bundle, and let $H, \operatorname{dim} H>0$, be a closed (and, consequently, Lie) subgroup of $G$. Then we have the composite bundle

$$
\begin{equation*}
P \rightarrow P / H \rightarrow X \tag{4.8.2}
\end{equation*}
$$

where

$$
\begin{equation*}
P_{\Sigma}=P \xrightarrow{\pi_{P \Sigma}} P / H \tag{4.8.3}
\end{equation*}
$$

is a principal bundle with a structure group $H$ and

$$
\begin{equation*}
\Sigma=P / H \xrightarrow{\pi_{\Sigma X}} X \tag{4.8.4}
\end{equation*}
$$

is a $P$-associated bundle with the typical fibre $G / H$ on which the structure group $G$ acts on the left (see Example 4.3.2).

One says that a structure Lie group $G$ of a principal bundle $P$ is reduced to its closed subgroup $H$ if the following equivalent conditions hold.

- A principal bundle $P$ admits a bundle atlas $\Psi_{P}$ (4.3.3) with $H$ valued transition functions $\varrho_{\alpha \beta}$.
- There exists a principal reduced subbundle $P_{H}$ of $P$ with a structure group $H$.

Theorem 4.8.1: There is one-to-one correspondence

$$
\begin{equation*}
P^{h}=\pi_{P \Sigma}^{-1}(h(X)) \tag{4.8.5}
\end{equation*}
$$

between the reduced principal $H$-subbundles $i_{h}: P^{h} \rightarrow P$ of $P$ and the global sections $h$ of the quotient bundle $P / H \rightarrow X$ (4.8.4).

Corollary 4.8.2: A glance at the formula (4.8.5) shows that the reduced principal $H$-bundle $P^{h}$ is the restriction $h^{*} P_{\Sigma}$ (1.1.13) of the principal $H$-bundle $P_{\Sigma}$ (4.8.3) to $h(X) \subset \Sigma$.

In general, there is topological obstruction to reduction of a structure group of a principal bundle to its subgroup.

Theorem 4.8.3: In accordance with Theorem 1.1.4, the structure group $G$ of a principal bundle $P$ is always reducible to its closed subgroup $H$, if the quotient $G / H$ is diffeomorphic to a Euclidean space $\mathbb{R}^{m}$.

In particular, this is the case of a maximal compact subgroup $H$ of a Lie group $G$. Then the following is a corollary of Theorem 4.8.3.

Theorem 4.8.4: A structure group $G$ of a principal bundle is always reducible to its maximal compact subgroup $H$.

Example 4.8.1: For instance, this is the case of $G=G L(n, \mathbb{C}), H=$ $U(n)$ and $G=G L(n, \mathbb{R}), H=O(n)$.

Example 4.8.2: Any affine bundle admits an atlas with linear transition functions. In accordance with Theorem 4.8.3, its structure group $G A(m, \mathbb{R})$ is always reducible to the linear subgroup $G L(m, \mathbb{R})$ because

$$
G A(m, \mathbb{R}) / G L(m, \mathbb{R})=\mathbb{R}^{m}
$$

Different principal $H$-subbundles $P^{h}$ and $P^{h^{\prime}}$ of a principal $G$-bundle $P$ are not isomorphic to each other in general.

Theorem 4.8.5: Let a structure Lie group $G$ of a principal bundle be reducible to its closed subgroup $H$.
(i) Every vertical principal automorphism $\Phi$ of $P$ sends a reduced principal $H$-subbundle $P^{h}$ of $P$ onto an isomorphic principal $H$-subbundle $P^{h^{\prime}}$.
(ii) Conversely, let two reduced subbundles $P^{h}$ and $P^{h^{\prime}}$ of a principal bundle $P \rightarrow X$ be isomorphic to each other, and let $\Phi: P^{h} \rightarrow P^{h^{\prime}}$ be their isomorphism over $X$. Then $\Phi$ is extended to a vertical principal automorphism of $P$.

Theorem 4.8.6: If the quotient $G / H$ is homeomorphic to a Euclidean space $\mathbb{R}^{m}$, all principal $H$-subbundles of a principal $G$-bundle $P$ are isomorphic to each other.

There are the following properties of principal connections compatible with a reduced structure.

Theorem 4.8.7: Since principal connections are equivariant, every principal connection $A_{h}$ on a reduced principal $H$-subbundle $P^{h}$ of a principal $G$-bundle $P$ gives rise to a principal connection on $P$.

Theorem 4.8.8: A principal connection $A$ on a principal $G$-bundle $P$ is reducible to a principal connection on a reduced principal $H$-subbundle $P^{h}$ of $P$ iff the corresponding global section $h$ of the $P$-associated fibre bundle $P / H \rightarrow X$ is an integral section of the associated principal connection $A$ on $P / H \rightarrow X$.

Theorem 4.8.9: Let the Lie algebra $\mathfrak{g}_{l}$ of $G$ be the direct sum

$$
\begin{equation*}
\mathfrak{g}_{l}=\mathfrak{h}_{l} \oplus \mathfrak{m} \tag{4.8.6}
\end{equation*}
$$

of the Lie algebra $\mathfrak{h}_{l}$ of $H$ and a subspace $\mathfrak{m}$ such that $\operatorname{Ad}_{g}(\mathfrak{m}) \subset \mathfrak{m}, g \in H$ (e.g., $H$ is a Cartan subgroup of $G$ ). Let $\bar{A}$ be a $\mathfrak{g} l$-valued connection form (4.4.12) on $P$. Then, the pull-back of the $\mathfrak{h}_{l}$-valued component of $\bar{A}$ onto a reduced principal $H$-subbundle $P^{h}$ is a $\mathfrak{h}_{l}$-valued connection form of a principal connection $\bar{A}_{h}$ on $P^{h}$.

The following is a corollary of Theorem 4.4.2.
Theorem 4.8.10: Given the composite bundle (4.8.2), let $A_{\Sigma}$ be a principal connection on the principal $H$-bundle $P \rightarrow \Sigma$ (4.8.3). Then, for any reduced principal $H$-bundle $i_{h}: P^{h} \rightarrow P$, the pull-back connection $i_{h}^{*} A_{\Sigma}$ (3.6.14) is a principal connection on $P^{h}$.

## Chapter 5

## Geometry of natural bundles

Classical gravitation theory is formulated as field theory on natural bundles, exemplified by tensor bundles. Therefore, we agree to call connections on these bundles the world connections.

### 5.1 Natural bundles

Let $\pi: Y \rightarrow X$ be a smooth fibre bundle coordinated by $\left(x^{\lambda}, y^{i}\right)$. Any automorphism $(\Phi, f)$ of $Y$, by definition, is projected as

$$
\pi \circ \Phi=f \circ \pi
$$

onto a diffeomorphism $f$ of its base $X$. The converse is not true. A diffeomorphism of $X$ need not give rise to an automorphism of $Y$, unless $Y \rightarrow X$ is a trivial bundle.

Given a one-parameter group $\left(\Phi_{t}, f_{t}\right)$ of automorphisms of $Y$, its infinitesimal generator is a projectable vector field

$$
u=u^{\lambda}\left(x^{\mu}\right) \partial_{\lambda}+u^{i}\left(x^{\mu}, y^{j}\right) \partial_{i}
$$

on $Y$. This vector field is projected as

$$
\tau \circ \pi=T \pi \circ u
$$

onto a vector field $\tau=u^{\lambda} \partial_{\lambda}$ on $X$. Its flow is the one-parameter group $\left(f_{t}\right)$ of diffeomorphisms of $X$ which are projections of autmorphisms
( $\Phi_{t}, f_{t}$ ) of $Y$. Conversely, let $\tau=\tau^{\lambda} \partial_{\lambda}$ be a vector field on $X$. There is a problem of constructing its lift to a projectable vector field

$$
u=\tau^{\lambda} \partial_{\lambda}+u^{i} \partial_{i}
$$

on $Y$ projected onto $\tau$. Such a lift always exists, but it need not be canonical. Given a connection $\Gamma$ on $Y$, any vector field $\tau$ on $X$ gives rise to the horizontal vector field $\Gamma \tau$ (3.1.6) on $Y$. This horizontal lift $\tau \rightarrow \Gamma \tau$ yields a monomorphism of the $C^{\infty}(X)$-module $\mathcal{T}(X)$ of vector fields on $X$ to the $C^{\infty}(Y)$-module of vector fields on $Y$, but this monomorphisms is not a Lie algebra morphism, unless $\Gamma$ is a flat connection.

Let us address the category of natural bundles $T \rightarrow X$ which admit the functorial lift $\widetilde{\tau}$ onto $T$ of any vector field $\tau$ on $X$ such that $\tau \rightarrow \bar{\tau}$ is a monomorphism

$$
\mathcal{T}(X) \rightarrow \mathcal{T}(T), \quad\left[\tilde{\tau}, \widetilde{\tau}^{\prime}\right]=\left[\widetilde{\tau, \tau^{\prime}}\right]
$$

of the real Lie algebra $\mathcal{T}(X)$ of vector fields on $X$ to the real Lie algebra $\mathcal{T}(Y)$ of vector fields on $T$. One treats the functorial lift $\widetilde{\tau}$ as an infinitesimal general covariant transformation, i.e., an infinitesimal generator of a local one-parameter group of general covariant transformations of $T$.

Remark 5.1.1: It should be emphasized that, in general, there exist diffeomorphisms of $X$ which do not belong to any one-parameter group of diffeomorphisms of $X$. In a general setting, one therefore considers a monomorphism $f \rightarrow \tilde{f}$ of the group of diffeomorphisms of $X$ to the group of bundle automorphisms of a natural bundle $T \rightarrow X$. Automorphisms $\tilde{f}$ are called general covariant transformations of $T$. No vertical automorphism of $T$, unless it is the identity morphism, is a general covariant transformation.

Natural bundles are exemplified by tensor bundles (1.2.5). For instance, the tangent and cotangent bundles $T X$ and $T^{*} X$ of $X$ are natural bundles. Given a vector field $\tau$ on $X$, its functorial (or canonical) lift onto the tensor bundle $T(1.2 .5)$ is given by the formula (1.3.2). In
particular, let us recall the functorial lift (1.3.4) and (1.3.5) of $\tau$ onto the tangent bundle $T X$ and the cotangent bundle $T^{*} X$, respectively.

Remark 5.1.2: Any diffeomorphism $f$ of $X$ gives rise to the tangent automorphisms $\tilde{f}=T f$ of $T X$ which is a general covariant transformation of $T X$ as a natural bundle. Accordingly, the general covariant transformation of the cotangent bundle $T^{*} X$ over a diffeomorphism $f$ of its base $X$ reads

$$
\dot{x}_{\mu}^{\prime}=\frac{\partial x^{\nu}}{\partial x^{\prime \mu}} \dot{x}_{\nu}
$$

Tensor bundles over a manifold $X$ have the structure group

$$
\begin{equation*}
G L_{n}=G L^{+}(n, \mathbb{R}) \tag{5.1.1}
\end{equation*}
$$

The associated principal bundle is the fibre bundle

$$
\pi_{L X}: L X \rightarrow X
$$

of oriented linear frames in the tangent spaces to a manifold $X$. It is called the linear frame bundle. Its (local) sections are termed frame fields.

Given holonomic frames $\left\{\partial_{\mu}\right\}$ in the tangent bundle $T X$ associated with the holonomic atlas $\Psi_{T}(1.2 .4)$, every element $\left\{H_{a}\right\}$ of the linear frame bundle $L X$ takes the form $H_{a}=H_{a}^{\mu} \partial_{\mu}$, where $H_{a}^{\mu}$ is a matrix of the natural representation of the group $G L_{n}$ in $\mathbb{R}^{n}$. These matrices constitute the bundle coordinates

$$
\left(x^{\lambda}, H_{a}^{\mu}\right), \quad H_{a}^{\prime \mu}=\frac{\partial x^{\prime \mu}}{\partial x^{\lambda}} H_{a}^{\lambda}
$$

on $L X$ associated to its holonomic atlas

$$
\begin{equation*}
\Psi_{T}=\left\{\left(U_{\iota}, z_{\iota}=\left\{\partial_{\mu}\right\}\right)\right\} \tag{5.1.2}
\end{equation*}
$$

given by the local frame fields $z_{\iota}=\left\{\partial_{\mu}\right\}$. With respect to these coordinates, the right action (4.3.2) of $G L_{n}$ on $L X$ reads

$$
R_{g P}: H_{a}^{\mu} \rightarrow H_{b}^{\mu} g_{a}^{b}, \quad g \in G L_{n}
$$

The frame bundle $L X$ admits the canonical $\mathbb{R}^{n}$-valued one-form

$$
\begin{equation*}
\theta_{L X}=H_{\mu}^{a} d x^{\mu} \otimes t_{a} \tag{5.1.3}
\end{equation*}
$$

where $\left\{t_{a}\right\}$ is a fixed basis for $\mathbb{R}^{n}$ and $H_{\mu}^{a}$ is the inverse matrix of $H_{a}^{\mu}$.
The frame bundle $L X \rightarrow X$ belongs to the category of natural bundles. Any diffeomorphism $f$ of $X$ gives rise to the principal automorphism

$$
\begin{equation*}
\tilde{f}:\left(x^{\lambda}, H_{a}^{\lambda}\right) \rightarrow\left(f^{\lambda}(x), \partial_{\mu} f^{\lambda} H_{a}^{\mu}\right) \tag{5.1.4}
\end{equation*}
$$

of $L X$ which is its general covariant transformation (or a holonomic automorphism). For instance, the associated automorphism of $T X$ is the tangent morphism $T f$ to $f$.

Given a (local) one-parameter group of diffeomorphisms of $X$ and its infinitesimal generator $\tau$, their lift (5.1.4) results in the functorial lift

$$
\begin{equation*}
\tilde{\tau}=\tau^{\mu} \partial_{\mu}+\partial_{\nu} \tau^{\alpha} H_{a}^{\nu} \frac{\partial}{\partial H_{a}^{\alpha}} \tag{5.1.5}
\end{equation*}
$$

of a vector field $\tau$ on $X$ onto $L X$ defined by the condition

$$
\mathbf{L}_{\tilde{\tau}} \theta_{L X}=0 .
$$

Every $L X$-associated bundle $Y \rightarrow X$ admits a lift of any diffeomorphism $f$ of its base to the principal automorphism $\tilde{f}_{Y}(4.7 .7)$ of $Y$ associated with the principal automorphism $\tilde{f}$ (5.1.4) of the liner frame bundle $L X$. Thus, all bundles associated with the linear frame bundle $L X$ are natural bundles. However, there are natural bundles which are not associated with $L X$.

Remark 5.1.3: In a more general setting, higher order natural bundles and gauge natural bundles are considered. Note that the linear frame bundle $L X$ over a manifold $X$ is the set of first order jets of local diffeomorphisms of the vector space $\mathbb{R}^{n}$ to $X, n=\operatorname{dim} X$, at the origin of $\mathbb{R}^{n}$. Accordingly, one considers $r$-order frame bundles $L^{r} X$ of $r$-order jets of local diffeomorphisms of $\mathbb{R}^{n}$ to $X$. Furthermore, given a principal bundle
$P \rightarrow X$ with a structure group $G$, the $r$-order jet bundle $J^{1} P \rightarrow X$ of its sections fails to be a principal bundle. However, the product

$$
W^{r} P=L^{r} X \times J^{r} P
$$

is a principal bundle with the structure group $W_{n}^{r} G$ which is a semidirect product of the group $G_{n}^{r}$ of invertible $r$-order jets of maps $\mathbb{R}^{n}$ to itself at its origin (e.g., $G_{n}^{1}=G L(n, \mathbb{R})$ ) and the group $T_{n}^{r} G$ of $r$-order jets of morphisms $\mathbb{R}^{n} \rightarrow G$ at the origin of $\mathbb{R}^{n}$. Moreover, if $Y \rightarrow X$ is a $P$-associated bundle, the jet bundle $J^{r} Y \rightarrow X$ is a vector bundle associated with the principal bundle $W^{r} P$. It exemplifies gauge natural bundles which can be described as fibre bundles associated with principal bundles $W^{r} P$. Natural bundles are gauge natural bundles for a trivial group $G=1$. The bundle of principal connections $C$ (4.4.3) is a first order gauge natural bundle.

### 5.2 Linear world connections

Since the tangent bundle $T X$ is associated with the linear frame bundle $L X$, every world connection (3.4.6):

$$
\begin{equation*}
\Gamma=d x^{\lambda} \otimes\left(\partial_{\lambda}+\Gamma_{\lambda}{ }^{\mu}{ }_{\nu} \dot{x}^{\nu} \dot{\partial}_{\mu}\right) \tag{5.2.1}
\end{equation*}
$$

on a manifold $X$ is associated with a principal connection on $L X$. We agree to call $\Gamma$ (5.2.1) the linear world connection in order to distinct it from an affine world connection in Section 5.3.

Being principal connections on the linear frame bundle $L X$, linear world connections are represented by sections of the quotient bundle

$$
\begin{equation*}
C_{\mathrm{W}}=J^{1} L X / G L_{n} \tag{5.2.2}
\end{equation*}
$$

called the bundle of world connections. With respect to the holonomic atlas $\Psi_{T}$ (5.1.2), this bundle is provided with the coordinates

$$
\left(x^{\lambda}, k_{\lambda}{ }^{\nu}{ }_{\alpha}\right), \quad k_{\lambda}^{\prime \nu}{ }_{\alpha}=\left[\frac{\partial x^{\prime \nu}}{\partial x^{\gamma}} \frac{\partial x^{\beta}}{\partial x^{\prime \alpha}} k_{\mu}{ }^{\gamma}{ }_{\beta}+\frac{\partial x^{\beta}}{\partial x^{\prime \alpha}} \frac{\partial^{2} x^{\prime \nu}}{\partial x^{\mu} \partial x^{\beta}}\right] \frac{\partial x^{\mu}}{\partial x^{\prime \lambda}},
$$

so that, for any section $\Gamma$ of $C_{\mathrm{W}} \rightarrow X$,

$$
k_{\lambda}{ }^{\nu}{ }_{\alpha} \circ \Gamma=\Gamma_{\lambda}{ }^{\nu}{ }_{\alpha}
$$

are components of the linear world connection $\Gamma$ (5.2.1).
Though the bundle of world connections $C_{\mathrm{W}} \rightarrow X(5.2 .2)$ is not $L X$ associated, it is a natural bundle. It admits the lift

$$
\tilde{f}_{C}: J^{1} L X / G L_{n} \rightarrow J^{1} \tilde{f}\left(J^{1} L X\right) / G L_{n}
$$

of any diffeomorphism $f$ of its base $X$ and, consequently, the functorial lift

$$
\tilde{\tau}_{C}=\tau^{\mu} \partial_{\mu}+\left[\partial_{\nu} \tau^{\alpha} k_{\mu}{ }^{\nu}{ }_{\beta}-\partial_{\beta} \tau^{\nu} k_{\mu}{ }^{\alpha}{ }_{\nu}-\partial_{\mu} \tau^{\nu} k_{\nu}{ }^{\alpha}{ }_{\beta}+\partial_{\mu \beta} \tau^{\alpha}\right] \frac{\partial}{\partial k_{\mu}{ }^{\alpha}{ }_{\beta}}(5.2 .3)
$$

of any vector field $\tau$ on $X$.
The first order jet manifold $J^{1} C_{\mathrm{W}}$ of the bundle of world connections admits the canonical splitting (4.5.11). In order to obtain its coordinate expression, let us consider the strength (4.7.12) of the linear world connection $\Gamma$ (5.2.1). It reads

$$
F_{\Gamma}=\frac{1}{2} F_{\lambda \mu}{ }^{b}{ }_{a} I_{b}{ }^{a} d x^{\lambda} \wedge d x^{\mu}=\frac{1}{2} R_{\mu \nu}{ }^{\alpha}{ }_{\beta} d x^{\lambda} \wedge d x^{\mu},
$$

where

$$
\left(I_{b}{ }^{a}\right)^{\alpha}{ }_{\beta}=H_{b}^{\alpha} H_{\beta}^{a}
$$

are generators of the group $G L_{n}(5.1 .1)$ in fibres of $T X$ with respect to the holonomic frames, and

$$
\begin{equation*}
R_{\lambda \mu}{ }^{\alpha}{ }_{\beta}=\partial_{\lambda} \Gamma_{\mu}{ }^{\alpha}{ }_{\beta}-\partial_{\mu} \Gamma_{\lambda}{ }^{\alpha}{ }_{\beta}+\Gamma_{\lambda}{ }^{\gamma}{ }_{\beta} \Gamma_{\mu}{ }^{\alpha}{ }_{\gamma}-\Gamma_{\mu}{ }^{\gamma}{ }_{\beta} \Gamma_{\lambda}{ }^{\alpha}{ }_{\gamma} \tag{5.2.4}
\end{equation*}
$$

are components if the curvature (3.4.8) of a linear world connection $\Gamma$. Accordingly, the above mentioned canonical splitting (4.5.11) of $J^{1} C_{\mathrm{W}}$ can be written in the form

$$
\begin{align*}
& k_{\lambda \mu}{ }^{\alpha}{ }_{\beta}=\frac{1}{2}\left(\mathcal{R}_{\lambda \mu}{ }^{\alpha}{ }_{\beta}+\mathcal{S}_{\lambda \mu}{ }^{\alpha}{ }_{\beta}\right)=  \tag{5.2.5}\\
& \quad \frac{1}{2}\left(k_{\lambda \mu}{ }^{\alpha}{ }_{\beta}-k_{\mu \lambda}{ }^{\alpha}{ }_{\beta}+k_{\lambda}{ }^{\gamma}{ }_{\beta} k_{\mu}{ }^{\alpha}{ }_{\gamma}-k_{\mu}{ }^{\gamma}{ }_{\beta} k_{\lambda}{ }^{\alpha}{ }_{\gamma}\right)+ \\
& \quad \frac{1}{2}\left(k_{\lambda \mu}{ }^{\alpha}{ }_{\beta}+k_{\mu \lambda}{ }^{\alpha}{ }_{\beta}-k_{\lambda}{ }^{\gamma}{ }_{\beta} k_{\mu}{ }^{\alpha}{ }_{\gamma}+k_{\mu}{ }^{\gamma}{ }_{\beta} k_{\lambda}{ }^{\alpha}{ }_{\gamma}\right) .
\end{align*}
$$

It is readily observed that, if $\Gamma$ is a section of $C_{\mathrm{W}} \rightarrow X$, then

$$
\mathcal{R}_{\lambda \mu}{ }^{\alpha}{ }_{\beta} \circ J^{1} \Gamma=R_{\lambda \mu}{ }^{\alpha}{ }_{\beta} .
$$

Because of the canonical vertical splitting (1.4.11) of the vertical tangent bundle VTX of $T X$, the curvature form (3.4.8) of a linear world connection $\Gamma$ can be represented by the tangent-valued two-form

$$
\begin{equation*}
R=\frac{1}{2} R_{\lambda \mu}{ }^{\alpha}{ }_{\beta} \dot{x}^{\beta} d x^{\lambda} \wedge d x^{\mu} \otimes \partial_{\alpha} \tag{5.2.6}
\end{equation*}
$$

on $T X$. Due to this representation, the Ricci tensor

$$
\begin{equation*}
R_{c}=\frac{1}{2} R_{\lambda \mu}{ }^{\lambda}{ }_{\beta} d x^{\mu} \otimes d x^{\beta} \tag{5.2.7}
\end{equation*}
$$

of a linear world connection $\Gamma$ is defined.
Owing to the above mentioned vertical splitting (1.4.11) of $V T X$, the torsion form $T$ (3.4.9) of $\Gamma$ can be written as the tangent-valued two-form

$$
\begin{align*}
& T=\frac{1}{2} T_{\mu}{ }^{\nu}{ }_{\lambda} d x^{\lambda} \wedge d x^{\mu} \otimes \partial_{\nu}  \tag{5.2.8}\\
& T_{\mu}{ }^{\nu}{ }_{\lambda}=\Gamma_{\mu}{ }^{\nu}{ }_{\lambda}-\Gamma_{\lambda}{ }^{\nu}{ }_{\mu}
\end{align*}
$$

on $X$. The soldering torsion form

$$
\begin{equation*}
T=T_{\mu}{ }^{\nu}{ }_{\lambda} \dot{x}^{\lambda} d x^{\mu} \otimes \dot{\partial}_{\nu} \tag{5.2.9}
\end{equation*}
$$

on $T X$ is also defined. Then one can show the following.

- Given a linear world connection $\Gamma(5.2 .1)$ and its soldering torsion form $T$ (5.2.9), the sum $\Gamma+c T, c \in \mathbb{R}$, is a linear world connection.
- Every linear world connection $\Gamma$ defines a unique symmetric world connection

$$
\begin{equation*}
\Gamma^{\prime}=\Gamma-\frac{1}{2} T \tag{5.2.10}
\end{equation*}
$$

- If $\Gamma$ and $\Gamma^{\prime}$ are linear world connections, then
$c \Gamma+(1-c) \Gamma^{\prime}$
is so for any $c \in \mathbb{R}$.

A manifold $X$ is said to be flat if it admits a flat linear world connection $\Gamma$. By virtue of Theorem 3.5.2, there exists an atlas of local constant trivializations of $T X$ such that

$$
\Gamma=d x^{\lambda} \otimes \partial_{\lambda}
$$

relative to this atlas. As a consequence, the curvature form $R$ (5.2.6) of this connection equals zero. However, such an atlas is not holonomic in general. Relative to this atlas, the canonical soldering form (1.4.7) on $T X$ reads

$$
\theta_{J}=H_{\mu}^{a} d x^{\mu} \dot{\partial}_{a}
$$

and the torsion form $T$ (3.4.9) of $\Gamma$ defined as the Nijenhuis differential $d_{\Gamma} \theta_{J}$ (3.3.8) need not vanish.

A manifold $X$ is called parallelizable if the tangent bundle $T X \rightarrow X$ is trivial. By virtue of Theorem 3.5.2, a parallelizable manifold is flat. Conversely, a flat manifold is parallelizable if it is simply connected.

Every linear world connection $\Gamma$ (5.2.1) yields the horizontal lift

$$
\begin{equation*}
\Gamma \tau=\tau^{\lambda}\left(\partial_{\lambda}+\Gamma_{\lambda}{ }^{\beta}{ }_{\alpha} \dot{x}^{\alpha} \dot{\partial}_{\beta}\right) \tag{5.2.11}
\end{equation*}
$$

of a vector field $\tau$ on $X$ onto the tangent bundle $T X$. A vector field $\tau$ on $X$ is said to be parallel relative to a connection $\Gamma$ if it is an integral section of $\Gamma$. Its integral curve is called the autoparallel of a world connection $\Gamma$.

Remark 5.2.1: By virtue of Theorem 3.2.2, any vector field on $X$ is an integral section of some linear world connection. If $\tau(x) \neq 0$ at a point $x \in X$, there exists a coordinate system $\left(q^{i}\right)$ on some neighbourhood $U$ of $x$ such that $\tau^{i}(x)=$ const. on $U$. Then $\tau$ on $U$ is an integral section of the local symmetric linear world connection

$$
\begin{equation*}
\Gamma_{\tau}(x)=d q^{i} \otimes \partial_{i}, \quad x \in U \tag{5.2.12}
\end{equation*}
$$

on $U$. The functorial lift $\widetilde{\tau}$ (1.3.4) can be obtained at each point $x \in X$ as the horizontal lift of $\tau$ by means of the local symmetric connection (5.2.12).

The horizontal lift of a vector field $\tau$ on $X$ onto the linear frame bundle $L X$ by means of a world connection $K$ reads

$$
\begin{equation*}
\Gamma \tau=\tau^{\lambda}\left(\partial_{\lambda}+\Gamma_{\lambda}{ }^{\nu}{ }_{\alpha} H_{a}^{\alpha} \frac{\partial}{\partial H_{a}^{\nu}}\right) . \tag{5.2.13}
\end{equation*}
$$

It is called standard if the morphism

$$
u\rfloor \theta_{L X}: L X \rightarrow \mathbb{R}^{n}
$$

is constant on $L X$. It is readily observed that every standard horizontal vector field on $L X$ takes the form

$$
\begin{equation*}
u_{v}=H_{b}^{\lambda} v^{b}\left(\partial_{\lambda}+\Gamma_{\lambda}{ }^{\nu}{ }_{\alpha} H_{a}^{\alpha} \frac{\partial}{\partial H_{a}^{\nu}}\right) \tag{5.2.14}
\end{equation*}
$$

where $v=v^{b} t_{b} \in \mathbb{R}^{n}$. A glance at this expression shows that a standard horizontal vector field is not projectable.

Since $T X$ is an $L X$-associated fibre bundle, we have the canonical morphism

$$
L X \times \mathbb{R}^{n} \rightarrow T X, \quad\left(H_{a}^{\mu}, v^{a}\right) \rightarrow \dot{x}^{\mu}=H_{a}^{\mu} v^{a} .
$$

The tangent map to this morphism sends every standard horizontal vector field (5.2.14) on $L X$ to the horizontal vector field

$$
\begin{equation*}
u=\dot{x}^{\lambda}\left(\partial_{\lambda}+\Gamma_{\lambda}{ }^{\nu}{ }_{\alpha} \dot{x}^{\alpha} \dot{\partial}_{\nu}\right) \tag{5.2.15}
\end{equation*}
$$

on $T X$. Such a vector field on $T X$ is called holonomic. Given holonomic coordinates ( $x^{\mu}, \dot{x}^{\mu}, \dot{\mathrm{x}}^{\mu}, \ddot{x}^{\mu}$ ) on the double tangent bundle TTX, the holonomic vector field (5.2.15) defines the second order dynamic equation

$$
\begin{equation*}
\ddot{x}^{\nu}=\Gamma_{\lambda}{ }^{\nu}{ }_{\alpha} \dot{x}^{\lambda} \dot{x}^{\alpha} \tag{5.2.16}
\end{equation*}
$$

on $X$ which is called the geodesic equation with respect to a linear world connection $\Gamma$. Solutions of the geodesic equation (5.2.16), called the geodesics of $\Gamma$, are the projection of integral curves of the vector field (5.2.15) in $T X$ onto $X$. Moreover, one can show the following.

Theorem 5.2.1: The projection of an integral curve of any standard horizontal vector field (5.2.14) on $L X$ onto $X$ is a geodesic in $X$. Conversely, any geodesic in $X$ is of this type.

It is readily observed that, if linear world connections $\Gamma$ and $\Gamma^{\prime}$ differ from each other only in the torsion, they define the same holonomic vector field (5.2.15) and the same geodesic equation (5.2.16).

Let $\tau$ be an integral vector field of a linear world connection $\Gamma$, i.e., $\nabla_{\mu}^{\Gamma} \tau=0$. Consequently, it obeys the equation $\tau^{\mu} \nabla_{\mu}^{\Gamma} \tau=0$. Any autoparallel of a linear world connection $\Gamma$ is its geodesic and, conversely, a geodesic of $\Gamma$ is an autoparallel of its symmetric part (5.2.10).

### 5.3 Affine world connections

The tangent bundle $T X$ of a manifold $X$ as like as any vector bundle possesses a natural structure of an affine bundle (see Section 1.2). Therefore, one can consider affine connections on $T X$, called affine world connections. Here we study them as principal connections.

Let $Y \rightarrow X$ be an affine bundle with an $k$-dimensional typical fibre $V$. It is associated with a principal bundle $A Y$ of affine frames in $Y$, whose structure group is the general affine group $G A(k, \mathbb{R})$. Then any affine connection on $Y \rightarrow X$ can be seen as an associated with a principal connection on $A Y \rightarrow X$. These connections are represented by global sections of the affine bundle

$$
J^{1} P / G A(k, \mathbb{R}) \rightarrow X
$$

They always exist.
As was mentioned in Section 1.3.5, every affine connection $\Gamma$ (3.4.12) on $Y \rightarrow X$ defines a unique associated linear connection $\bar{\Gamma}$ (3.4.13) on the underlying vector bundle $\bar{Y} \rightarrow X$. This connection $\bar{\Gamma}$ is associated with a linear principal connection on the principal bundle $L \bar{Y}$ of linear
frames in $\bar{Y}$ whose structure group is the general linear group $G L(k, \mathbb{R})$. We have the exact sequence of groups

$$
\begin{equation*}
0 \rightarrow T_{k} \rightarrow G A(k, \mathbb{R}) \rightarrow G L(k, \mathbb{R}) \rightarrow \mathbf{1} \tag{5.3.1}
\end{equation*}
$$

where $T_{k}$ is the group of translations in $\mathbb{R}^{k}$. It is readily observed that there is the corresponding principal bundle morphism $A Y \rightarrow L \bar{Y}$ over $X$, and the principal connection $\bar{\Gamma}$ on $L \bar{Y}$ is the image of the principal connection $\Gamma$ on $A Y \rightarrow X$ under this morphism in accordance with Theorem 4.4.3.

The exact sequence (5.3.1) admits a splitting

$$
G L(k, \mathbb{R}) \rightarrow G A(k, \mathbb{R})
$$

but this splitting is not canonical. It depends on the morphism

$$
V \ni v \rightarrow v-v_{0} \in \bar{V}
$$

i.e., on the choice of an origin $v_{0}$ of the affine space $V$. Given $v_{0}$, the image of the corresponding monomorphism

$$
G L(k, \mathbb{R}) \rightarrow G A(k, \mathbb{R})
$$

is a stabilizer

$$
G\left(v_{0}\right) \subset G A(k, \mathbb{R})
$$

of $v_{0}$. Different subgroups $G\left(v_{0}\right)$ and $G\left(v_{0}^{\prime}\right)$ are related to each other as follows:

$$
G\left(v_{0}^{\prime}\right)=T\left(v_{0}^{\prime}-v_{0}\right) G\left(v_{0}\right) T^{-1}\left(v_{0}^{\prime}-v_{0}\right)
$$

where $T\left(v_{0}^{\prime}-v_{0}\right)$ is the translation along the vector $\left(v_{0}^{\prime}-v_{0}\right) \in \bar{V}$.
Remark 5.3.1: The well-known morphism of a $k$-dimensional affine space $V$ onto a hypersurface $\bar{y}^{k+1}=1$ in $\mathbb{R}^{k+1}$ and the corresponding representation of elements of $G A(k, \mathbb{R})$ by particular $(k+1) \times(k+1)$ matrices also fail to be canonical. They depend on a point $v_{0} \in V$ sent to vector $(0, \ldots, 0,1) \in \mathbb{R}^{k+1}$.

One can say something more if $Y \rightarrow X$ is a vector bundle provided with the natural structure of an affine bundle whose origin is the canonical zero section $\hat{0}$. In this case, we have the canonical splitting of the exact sequence (5.3.1) such that $G L(k, \mathbb{R})$ is a subgroup of $G A(k, \mathbb{R})$ and $G A(k, \mathbb{R})$ is the semidirect product of $G L(k, \mathbb{R})$ and the group $T(k, \mathbb{R})$ of translations in $\mathbb{R}^{k}$. Given a $G A(k, \mathbb{R})$-principal bundle $A Y \rightarrow X$, its affine structure group $G A(k, \mathbb{R})$ is always reducible to the linear subgroup since the quotient $G A(k, \mathbb{R}) / G L(k, \mathbb{R})$ is a vector space $\mathbb{R}^{k}$ provided with the natural affine structure (see Example 4.8.2). The corresponding quotient bundle is isomorphic to the vector bundle $Y \rightarrow X$. There is the canonical injection of the linear frame bundle $L Y \rightarrow A Y$ onto the reduced $G L(k, \mathbb{R})$-principal subbundle of $A Y$ which corresponds to the zero section $\hat{0}$ of $Y \rightarrow X$. In this case, every principal connection on the linear frame bundle $L Y$ gives rise to a principal connection on the affine frame bundle in accordance with Theorem 4.8.7. This is equivalent to the fact that any affine connection $\Gamma$ on a vector bundle $Y \rightarrow X$ defines a linear connection $\bar{\Gamma}$ on $Y \rightarrow X$ and that every linear connection on $Y \rightarrow X$ can be seen as an affine one. Then any affine connection $\Gamma$ on the vector bundle $Y \rightarrow X$ is represented by the sum of the associated linear connection $\bar{\Gamma}$ and a basic soldering form $\sigma$ on $Y \rightarrow X$. Due to the vertical splitting (1.2.8), this soldering form is represented by a global section of the tensor product $T^{*} X \otimes Y$.

Let now $Y \rightarrow X$ be the tangent bundle $T X \rightarrow X$ considered as an affine bundle. Then the relationship between affine and linear world connections on $T X$ is the repetition of that we have said in the case of an arbitrary vector bundle $Y \rightarrow X$. In particular, any affine world connection

$$
\begin{equation*}
\Gamma=d x^{\lambda} \otimes\left(\partial_{\lambda}+\Gamma_{\lambda}{ }^{\alpha}{ }_{\mu}(x) \dot{x}^{\mu}+\sigma_{\lambda}^{\alpha}(x)\right) \partial_{\alpha} \tag{5.3.2}
\end{equation*}
$$

on $T X \rightarrow X$ is represented by the sum of the associated linear world
connection

$$
\begin{equation*}
\bar{\Gamma}=\Gamma_{\lambda}{ }^{\alpha}{ }_{\mu}(x) \dot{x}^{\mu} d x^{\lambda} \otimes \partial_{\alpha} \tag{5.3.3}
\end{equation*}
$$

on $T X \rightarrow X$ and a basic soldering form

$$
\begin{equation*}
\sigma=\sigma_{\lambda}^{\alpha}(x) d x^{\lambda} \otimes \partial_{\alpha} \tag{5.3.4}
\end{equation*}
$$

on $Y \rightarrow X$, which is the $(1,1)$-tensor field on $X$. For instance, if $\sigma=\theta_{X}$ (1.4.5), we have the Cartan connection (3.4.15).

It is readily observed that the soldered curvature (3.3.7) of any soldering form (5.3.4) equals zero. Then we obtain from (3.3.10) that the torsion (3.4.14) of the affine connection $\Gamma$ (5.3.2) with respect to $\sigma$ (5.3.4) coincides with that of the associated linear connection $\bar{\Gamma}$ (5.3.3) and reads

$$
\begin{align*}
& T=\frac{1}{2} T_{\lambda \mu}^{i} d x^{\mu} \wedge d x^{\lambda} \otimes \partial_{i} \\
& T_{\lambda}{ }^{\lambda}{ }_{\mu}=\Gamma_{\lambda}{ }^{\alpha}{ }_{\nu} \sigma_{\mu}^{\nu}-\Gamma_{\mu}{ }^{\alpha}{ }_{\nu} \sigma_{\lambda}^{\nu} \tag{5.3.5}
\end{align*}
$$

The relation between the curvatures of an affine world connection $\Gamma$ (5.3.2) and the associated linear connection $\bar{\Gamma}$ (5.3.3) is given by the general expression (3.3.11) where $\rho=0$ and $T$ is (5.3.5).

## Chapter 6

## Geometry of graded manifolds

In classical field theory, there are different descriptions of odd fields on graded manifolds and supermanifolds. Both graded manifolds and supermanifolds are phrased in terms of sheaves of graded commutative algebras. However, graded manifolds are characterized by sheaves on smooth manifolds, while supermanifolds are constructed by gluing of sheaves on supervector spaces. Treating odd fields on a manifold $X$, we follow the Serre-Swan theorem generalized to graded manifolds (Theorem 6.3.2). It states that, if a Grassmann algebra is an exterior algebra of some projective $C^{\infty}(X)$-module of finite rank, it is isomorphic to the algebra of graded functions on a graded manifold whose body is $X$. By virtue of this theorem, odd fields on an arbitrary manifold $X$ are described as generating elements of the structure ring of a graded manifold whose body is $X[9,24]$.

### 6.1 Grassmann-graded algebraic calculus

Throughout the book, by the Grassmann gradation is meant $\mathbb{Z}_{2}$-gradation. Hereafter, the symbol [.] stands for the Grassmann parity. In the literature, a $\mathbb{Z}_{2}$-graded structure is simply called the graded structure if there is no danger of confusion. Let us summarize the relevant notions of the Grassmann-graded algebraic calculus.

An algebra $\mathcal{A}$ is called graded if it is endowed with a grading automorphism $\gamma$ such that $\gamma^{2}=\mathrm{Id}$. A graded algebra falls into the direct $\operatorname{sum} \mathcal{A}=\mathcal{A}_{0} \oplus \mathcal{A}_{1}$ of $\mathbb{Z}$-modules $\mathcal{A}_{0}$ and $\mathcal{A}_{1}$ of even and odd elements such that

$$
\begin{aligned}
\gamma(a) & =(-1)^{i} a, \quad a \in \mathcal{A}_{i}, \quad i=0,1, \\
{\left[a a^{\prime}\right] } & =\left([a]+\left[a^{\prime}\right]\right) \bmod 2, \quad a \in \mathcal{A}_{[a]}, \quad a^{\prime} \in \mathcal{A}_{\left[a^{\prime}\right]}
\end{aligned}
$$

One calls $\mathcal{A}_{0}$ and $\mathcal{A}_{1}$ the even and odd parts of $\mathcal{A}$, respectively. The even part $\mathcal{A}_{0}$ is a subalgebra of $\mathcal{A}$ and the odd one $\mathcal{A}_{1}$ is an $\mathcal{A}_{0}$-module. If $\mathcal{A}$ is a graded ring, then $[\mathbf{1}]=0$.

A graded algebra $\mathcal{A}$ is called graded commutative if

$$
a a^{\prime}=(-1)^{[a]\left[a^{\prime}\right]} a^{\prime} a
$$

where $a$ and $a^{\prime}$ are graded-homogeneous elements of $\mathcal{A}$.
Given a graded algebra $\mathcal{A}$, a left graded $\mathcal{A}$-module $Q$ is defined as a left $\mathcal{A}$-module provided with the grading automorphism $\gamma$ such that

$$
\begin{aligned}
& \gamma(a q)=\gamma(a) \gamma(q), \quad a \in \mathcal{A}, \quad q \in Q \\
& {[a q]=([a]+[q]) \bmod 2}
\end{aligned}
$$

A graded module $Q$ is split into the direct sum $Q=Q_{0} \oplus Q_{1}$ of two $\mathcal{A}_{0}$-modules $Q_{0}$ and $Q_{1}$ of even and odd elements.

If $\mathcal{K}$ is a graded commutative ring, a graded $\mathcal{K}$-module can be provided with a graded $\mathcal{K}$-bimodule structure by letting

$$
q a=(-1)^{[a][q]} a q, \quad a \in \mathcal{K}, \quad q \in Q
$$

A graded module is called free if it has a basis generated by gradedhomogeneous elements. This basis is said to be of type $(n, m)$ if it contains $n$ even and $m$ odd elements.

In particular, by a real graded vector space $B=B_{0} \oplus B_{1}$ is meant a graded $\mathbb{R}$-module. A real graded vector space is said to be $(n, m)$ dimensional if $B_{0}=\mathbb{R}^{n}$ and $B_{1}=\mathbb{R}^{m}$.

Given a graded commutative ring $\mathcal{K}$, the following are standard constructions of new graded modules from old ones.

- The direct sum of graded modules and a graded factor module are defined just as those of modules over a commutative ring.
- The tensor product $P \otimes Q$ of graded $\mathcal{K}$-modules $P$ and $Q$ is an additive group generated by elements $p \otimes q, p \in P, q \in Q$, obeying the relations

$$
\begin{aligned}
& \left(p+p^{\prime}\right) \otimes q=p \otimes q+p^{\prime} \otimes q \\
& p \otimes\left(q+q^{\prime}\right)=p \otimes q+p \otimes q^{\prime} \\
& a p \otimes q=(-1)^{[p][a]} p a \otimes q=(-1)^{[p][a]} p \otimes a q, \quad a \in \mathcal{K}
\end{aligned}
$$

In particular, the tensor algebra $\otimes P$ of a graded $\mathcal{K}$-module $P$ is defined as that (8.1.5) of a module over a commutative ring. Its quotient $\wedge P$ with respect to the ideal generated by elements

$$
p \otimes p^{\prime}+(-1)^{[p]\left[p^{\prime}\right]} p^{\prime} \otimes p, \quad p, p^{\prime} \in P
$$

is the bigraded exterior algebra of a graded module $P$ with respect to the graded exterior product

$$
p \wedge p^{\prime}=-(-1)^{[p]\left[p^{\prime}\right]} p^{\prime} \wedge p
$$

- A morphism $\Phi: P \rightarrow Q$ of graded $\mathcal{K}$-modules seen as additive groups is said to be even graded morphism (resp. odd graded morphism) if $\Phi$ preserves (resp. change) the Grassmann parity of all gradedhomogeneous elements of $P$ and obeys the relations

$$
\Phi(a p)=(-1)^{[\Phi][a]} a \Phi(p), \quad p \in P, \quad a \in \mathcal{K}
$$

A morphism $\Phi: P \rightarrow Q$ of graded $\mathcal{K}$-modules as additive groups is called a graded $\mathcal{K}$-module morphism if it is represented by a sum of even and odd graded morphisms. The set $\operatorname{Hom}_{\mathcal{K}}(P, Q)$ of graded morphisms of a graded $\mathcal{K}$-module $P$ to a graded $\mathcal{K}$-module $Q$ is naturally a graded $\mathcal{K}$-module. The graded $\mathcal{K}$-module $P^{*}=\operatorname{Hom}_{\mathcal{K}}(P, \mathcal{K})$ is called the dual of a graded $\mathcal{K}$-module $P$.

A graded commutative $\mathcal{K}$-ring $\mathcal{A}$ is a graded commutative ring which also is a graded $\mathcal{K}$-module. A real graded commutative ring is said to be of rank $N$ if it is a free algebra generated by the unit 1 and $N$ odd elements. A graded commutative Banach ring $\mathcal{A}$ is a real graded commutative ring which is a real Banach algebra whose norm obeys the condition

$$
\left\|a_{0}+a_{1}\right\|=\left\|a_{0}\right\|+\left\|a_{1}\right\|, \quad a_{0} \in \mathcal{A}_{0}, \quad a_{1} \in \mathcal{A}_{1} .
$$

Let $V$ be a real vector space, and let $\Lambda=\wedge V$ be its exterior algebra endowed with the Grassmann gradation

$$
\begin{equation*}
\Lambda=\Lambda_{0} \oplus \Lambda_{1}, \quad \Lambda_{0}=\mathbb{R} \bigoplus_{k=1}^{\wedge} \wedge V, \quad \Lambda_{1}=\bigoplus_{k=1}^{2 k-1} \wedge^{2 k} V \tag{6.1.1}
\end{equation*}
$$

It is a real graded commutative ring, called the Grassmann algebra. A Grassmann algebra, seen as an additive group, admits the decomposition

$$
\begin{equation*}
\Lambda=\mathbb{R} \oplus R=\mathbb{R} \oplus R_{0} \oplus R_{1}=\mathbb{R} \oplus\left(\Lambda_{1}\right)^{2} \oplus \Lambda_{1}, \tag{6.1.2}
\end{equation*}
$$

where $R$ is the ideal of nilpotents of $\Lambda$. The corresponding projections $\sigma: \Lambda \rightarrow \mathbb{R}$ and $s: \Lambda \rightarrow R$ are called the body and soul maps, respectively.

Hereafter, we restrict our consideration to Grassmann algebras of finite rank. Given a basis $\left\{c^{i}\right\}$ for the vector space $V$, the elements of the Grassmann algebra $\Lambda$ (6.1.1) take the form

$$
\begin{equation*}
a=\sum_{k=0,1, \ldots\left(i_{1} \cdots i_{k}\right)} a_{i_{1} \cdots i_{k}} c^{i_{1}} \cdots c^{i_{k}}, \tag{6.1.3}
\end{equation*}
$$

where the second sum runs through all the tuples $\left(i_{1} \cdots i_{k}\right)$ such that no two of them are permutations of each other. The Grassmann algebra $\Lambda$ becomes a graded commutative Banach ring with respect to the norm

$$
\|a\|=\sum_{k=0} \sum_{\left(i_{1} \cdots i_{k}\right)}\left|a_{i_{1} \cdots i_{k}}\right| .
$$

Let $B$ be a graded vector space. Given a Grassmann algebra $\Lambda$, it can be brought into a graded $\Lambda$-module

$$
\Lambda B=\Lambda B_{0} \oplus \Lambda B_{1}=\left(\Lambda_{0} \otimes B_{0} \oplus \Lambda_{1} \otimes B_{1}\right) \oplus\left(\Lambda_{1} \otimes B_{0} \oplus \Lambda_{0} \otimes B_{1}\right),
$$

called a superspace. The superspace

$$
\begin{equation*}
B^{n \mid m}=\left[\left(\stackrel{n}{\oplus} \Lambda_{0}\right) \oplus\left(\stackrel{m}{\oplus} \Lambda_{1}\right)\right] \oplus\left[\left(\stackrel{n}{\oplus} \Lambda_{1}\right) \oplus\left(\stackrel{m}{\oplus} \Lambda_{0}\right)\right] \tag{6.1.4}
\end{equation*}
$$

is said to be $(n, m)$-dimensional. The graded $\Lambda_{0}$-module

$$
B^{n, m}=\left(\stackrel{n}{\oplus} \Lambda_{0}\right) \oplus\left(\stackrel{m}{\oplus} \Lambda_{1}\right)
$$

is called an $(n, m)$-dimensional supervector space. Whenever referring to a topology on a supervector space $B^{n, m}$, we will mean the Euclidean topology on a $2^{N-1}[n+m]$-dimensional real vector space.

Let $\mathcal{K}$ be a graded commutative ring. A graded commutative (nonassociative) $\mathcal{K}$-algebra $\mathfrak{g}$ is called a Lie $\mathcal{K}$-superalgebra if its product [., .], called the Lie superbracket, obeys the relations

$$
\begin{aligned}
& {\left[\varepsilon, \varepsilon^{\prime}\right]=-(-1)^{[\varepsilon]\left[\varepsilon^{\prime}\right]}\left[\varepsilon^{\prime}, \varepsilon\right]} \\
& (-1)^{[\varepsilon]\left[\varepsilon^{\prime \prime}\right]}\left[\varepsilon,\left[\varepsilon^{\prime}, \varepsilon^{\prime \prime}\right]\right]+(-1)^{\left[\varepsilon^{\prime}\right][\varepsilon]}\left[\varepsilon^{\prime},\left[\varepsilon^{\prime \prime}, \varepsilon\right]\right]+(-1)^{\left[\varepsilon^{\prime \prime}\right]\left[\varepsilon^{\prime}\right]}\left[\varepsilon^{\prime \prime},\left[\varepsilon, \varepsilon^{\prime}\right]\right]=0 .
\end{aligned}
$$

The even part $\mathfrak{g}_{0}$ of a Lie $\mathcal{K}$-superalgebra $\mathfrak{g}$ is a Lie $\mathcal{K}_{0}$-algebra. A graded $\mathcal{K}$-module $P$ is called a $\mathfrak{g}$-module if it is provided with a $\mathcal{K}$-bilinear map

$$
\begin{aligned}
& \mathfrak{g} \times P \ni(\varepsilon, p) \rightarrow \varepsilon p \in P, \quad[\varepsilon p]=([\varepsilon]+[p]) \bmod 2, \\
& {\left[\varepsilon, \varepsilon^{\prime}\right] p=\left(\varepsilon \circ \varepsilon^{\prime}-(-1)^{[\varepsilon]\left[\varepsilon^{\prime}\right]} \varepsilon^{\prime} \circ \varepsilon\right) p .}
\end{aligned}
$$

### 6.2 Grassmann-graded differential calculus

Linear differential operators on graded modules over a graded commutative ring are defined similarly to those in commutative geometry (Section 8.2).

Let $\mathcal{K}$ be a graded commutative ring and $\mathcal{A}$ a graded commutative $\mathcal{K}$-ring. Let $P$ and $Q$ be graded $\mathcal{A}$-modules. The graded $\mathcal{K}$-module $\operatorname{Hom}_{\mathcal{K}}(P, Q)$ of graded $\mathcal{K}$-module homomorphisms $\Phi: P \rightarrow Q$ can be endowed with the two graded $\mathcal{A}$-module structures

$$
\begin{equation*}
(a \Phi)(p)=a \Phi(p), \quad(\Phi \bullet a)(p)=\Phi(a p), \quad a \in \mathcal{A}, \quad p \in P \tag{6.2.1}
\end{equation*}
$$

called $\mathcal{A}$ - and $\mathcal{A}^{\bullet}$-module structures, respectively. Let us put

$$
\begin{equation*}
\delta_{a} \Phi=a \Phi-(-1)^{[a][\Phi]} \Phi \bullet a, \quad a \in \mathcal{A} \tag{6.2.2}
\end{equation*}
$$

An element $\Delta \in \operatorname{Hom}_{\mathcal{K}}(P, Q)$ is said to be a $Q$-valued graded differential operator of order $s$ on $P$ if

$$
\delta_{a_{0}} \circ \cdots \circ \delta_{a_{s}} \Delta=0
$$

for any tuple of $s+1$ elements $a_{0}, \ldots, a_{s}$ of $\mathcal{A}$. The set $\operatorname{Diff}_{s}(P, Q)$ of these operators inherits the graded module structures (6.2.1).

In particular, zero order graded differential operators obey the condition

$$
\delta_{a} \Delta(p)=a \Delta(p)-(-1)^{[a][\Delta]} \Delta(a p)=0, \quad a \in \mathcal{A}, \quad p \in P
$$

i.e., they coincide with graded $\mathcal{A}$-module morphisms $P \rightarrow Q$. A first order graded differential operator $\Delta$ satisfies the relation

$$
\begin{aligned}
& \delta_{a} \circ \delta_{b} \Delta(p)=a b \Delta(p)-(-1)^{([b]+[\Delta])[a]} b \Delta(a p)-(-1)^{[b][\Delta]} a \Delta(b p)+ \\
& \quad(-1)^{[b][\Delta]+([\Delta]+[b])[a]}=0, \quad a, b \in \mathcal{A}, \quad p \in P .
\end{aligned}
$$

For instance, let $P=\mathcal{A}$. Any zero order $Q$-valued graded differential operator $\Delta$ on $\mathcal{A}$ is defined by its value $\Delta(\mathbf{1})$. Then there is a graded $\mathcal{A}$-module isomorphism

$$
\operatorname{Diff}_{0}(\mathcal{A}, Q)=Q, \quad Q \ni q \rightarrow \Delta_{q} \in \operatorname{Diff}_{0}(\mathcal{A}, Q)
$$

where $\Delta_{q}$ is given by the equality $\Delta_{q}(\mathbf{1})=q$. A first order $Q$-valued graded differential operator $\Delta$ on $\mathcal{A}$ fulfils the condition

$$
\Delta(a b)=\Delta(a) b+(-1)^{[a][\Delta]} a \Delta(b)-(-1)^{[[b]+[a])[\Delta]} a b \Delta(\mathbf{1}), \quad a, b \in \mathcal{A}
$$

It is called a $Q$-valued graded derivation of $\mathcal{A}$ if $\Delta(\mathbf{1})=0$, i.e., the Grassmann-graded Leibniz rule

$$
\begin{equation*}
\Delta(a b)=\Delta(a) b+(-1)^{[a][\Delta]} a \Delta(b), \quad a, b \in \mathcal{A} \tag{6.2.3}
\end{equation*}
$$

holds. One obtains at once that any first order graded differential operator on $\mathcal{A}$ falls into the sum

$$
\Delta(a)=\Delta(\mathbf{1}) a+[\Delta(a)-\Delta(\mathbf{1}) a]
$$

of a zero order graded differential operator $\Delta(\mathbf{1}) a$ and a graded derivation $\Delta(a)-\Delta(\mathbf{1}) a$. If $\partial$ is a graded derivation of $\mathcal{A}$, then $a \partial$ is so for any $a \in \mathcal{A}$. Hence, graded derivations of $\mathcal{A}$ constitute a graded $\mathcal{A}$-module $\mathfrak{o}(\mathcal{A}, Q)$, called the graded derivation module.

If $Q=\mathcal{A}$, the graded derivation module $\mathfrak{d} \mathcal{A}$ also is a Lie superalgebra over the graded commutative ring $\mathcal{K}$ with respect to the superbracket

$$
\begin{equation*}
\left[u, u^{\prime}\right]=u \circ u^{\prime}-(-1)^{[u]\left[u^{\prime}\right]} u^{\prime} \circ u, \quad u, u^{\prime} \in \mathcal{A} . \tag{6.2.4}
\end{equation*}
$$

We have the graded $\mathcal{A}$-module decomposition

$$
\begin{equation*}
\operatorname{Diff}_{1}(\mathcal{A})=\mathcal{A} \oplus \mathcal{O} \mathcal{A} \tag{6.2.5}
\end{equation*}
$$

Since $\mathfrak{o A}$ is a Lie $\mathcal{K}$-superalgebra, let us consider the ChevalleyEilenberg complex $C^{*}[\mathfrak{A} ; \mathcal{A}]$ where the graded commutative ring $\mathcal{A}$ is a regarded as a $\mathfrak{O} \mathcal{A}$-module. It is the complex

$$
\begin{equation*}
0 \rightarrow \mathcal{A} \xrightarrow{d} C^{1}[\mathfrak{o} \mathcal{A} ; \mathcal{A}] \xrightarrow{d} \cdots C^{k}[\mathfrak{o} \mathcal{A} ; \mathcal{A}] \xrightarrow{d} \cdots \tag{6.2.6}
\end{equation*}
$$

where

$$
C^{k}[\mathcal{O} ; \mathcal{A}]=\operatorname{Hom}_{\mathcal{K}}(\stackrel{k}{\wedge} \mathcal{O} \mathcal{A}, \mathcal{A})
$$

are $\mathfrak{D} \mathcal{A}$-modules of $\mathcal{K}$-linear graded morphisms of the graded exterior products ${ }^{k} \mathfrak{d} \mathcal{A}$ of the $\mathcal{K}$-module $\mathfrak{d} \mathcal{A}$ to $\mathcal{A}$. Let us bring homogeneous elements of ${ }_{\wedge}^{k} \mathcal{O} \mathcal{A}$ into the form

$$
\varepsilon_{1} \wedge \cdots \varepsilon_{r} \wedge \epsilon_{r+1} \wedge \cdots \wedge \epsilon_{k}, \quad \varepsilon_{i} \in \mathfrak{o} \mathcal{A}_{0}, \quad \epsilon_{j} \in \mathfrak{o} \mathcal{A}_{1}
$$

Then the even coboundary operator $d$ of the complex (6.2.6) is given by the expression

$$
\begin{equation*}
d c\left(\varepsilon_{1} \wedge \cdots \wedge \varepsilon_{r} \wedge \epsilon_{1} \wedge \cdots \wedge \epsilon_{s}\right)= \tag{6.2.7}
\end{equation*}
$$

$$
\begin{aligned}
& \sum_{i=1}^{r}(-1)^{i-1} \varepsilon_{i} c\left(\varepsilon_{1} \wedge \cdots \widehat{\varepsilon}_{i} \cdots \wedge \varepsilon_{r} \wedge \epsilon_{1} \wedge \cdots \epsilon_{s}\right)+ \\
& \sum_{j=1}^{s}(-1)^{r} \varepsilon_{i} c\left(\varepsilon_{1} \wedge \cdots \wedge \varepsilon_{r} \wedge \epsilon_{1} \wedge \cdots \widehat{\epsilon}_{j} \cdots \wedge \epsilon_{s}\right)+ \\
& \sum_{1 \leq i<j \leq r}(-1)^{i+j} c\left(\left[\varepsilon_{i}, \varepsilon_{j}\right] \wedge \varepsilon_{1} \wedge \cdots \widehat{\varepsilon}_{i} \cdots \widehat{\varepsilon}_{j} \cdots \wedge \varepsilon_{r} \wedge \epsilon_{1} \wedge \cdots \wedge \epsilon_{s}\right)+ \\
& \sum_{1 \leq i<j \leq s} c\left(\left[\epsilon_{i}, \epsilon_{j}\right] \wedge \varepsilon_{1} \wedge \cdots \wedge \varepsilon_{r} \wedge \epsilon_{1} \wedge \cdots \widehat{\epsilon}_{i} \cdots \widehat{\epsilon}_{j} \cdots \wedge \epsilon_{s}\right)+ \\
& \sum_{1 \leq i<r, 1 \leq j \leq s}(-1)^{i+r+1} c\left(\left[\varepsilon_{i}, \epsilon_{j}\right] \wedge \varepsilon_{1} \cdots \widehat{\varepsilon}_{i} \cdots \wedge \varepsilon_{r} \wedge \epsilon_{1} \cdots \widehat{\epsilon}_{j} \cdots \wedge \epsilon_{s}\right)
\end{aligned}
$$

where the caret ^ denotes omission. This operator is called the graded Chevalley-Eilenberg coboundary operator.

Let us consider the extended Chevalley-Eilenberg complex

$$
0 \rightarrow \mathcal{K} \xrightarrow{\text { in }} C^{*}[\mathfrak{d} \mathcal{A} ; \mathcal{A}] .
$$

This complex contains a subcomplex $\mathcal{O}^{*}[\mathfrak{d} \mathcal{A}]$ of $\mathcal{A}$-linear graded morphisms. The $\mathbb{N}$-graded module $\mathcal{O}^{*}[\mathfrak{d} \mathcal{A}]$ is provided with the structure of a bigraded $\mathcal{A}$-algebra with respect to the graded exterior product

$$
\begin{align*}
& \phi \wedge \phi^{\prime}\left(u_{1}, \ldots, u_{r+s}\right)=  \tag{6.2.8}\\
& \quad \sum_{i_{1}<\cdots<i_{r} ; j_{1}<\cdots<j_{s}} \operatorname{Sgn}_{1 \cdots r+s}^{i_{1} \cdots i_{r} j_{1} \cdots j_{s}} \phi\left(u_{i_{1}}, \ldots, u_{i_{r}}\right) \phi^{\prime}\left(u_{j_{1}}, \ldots, u_{j_{s}}\right), \\
& \phi \in \mathcal{O}^{r}[\mathfrak{d}], \quad \phi^{\prime} \in \mathcal{O}^{s}[\mathfrak{o} \mathcal{A}], \quad u_{k} \in \mathfrak{o} \mathcal{A},
\end{align*}
$$

where $u_{1}, \ldots, u_{r+s}$ are graded-homogeneous elements of $\mathfrak{o} \mathcal{A}$ and

$$
u_{1} \wedge \cdots \wedge u_{r+s}=\operatorname{Sgn}_{1 \cdots r+s}^{i_{1} \cdots i_{r} j_{1} \cdots j_{s}} u_{i_{1}} \wedge \cdots \wedge u_{i_{r}} \wedge u_{j_{1}} \wedge \cdots \wedge u_{j_{s}}
$$

The graded Chevalley-Eilenberg coboundary operator $d$ (6.2.7) and the graded exterior product $\wedge(6.2 .8)$ bring $\mathcal{O}^{*}[\mathfrak{o} \mathcal{A}]$ into a differential bigraded algebra (henceforth DBGA) whose elements obey the relations

$$
\begin{align*}
& \phi \wedge \phi^{\prime}=(-1)^{|\phi|\left|\phi^{\prime}\right|+[\phi]\left[\phi^{\prime}\right]} \phi^{\prime} \wedge \phi  \tag{6.2.9}\\
& d\left(\phi \wedge \phi^{\prime}\right)=d \phi \wedge \phi^{\prime}+(-1)^{|\phi|} \phi \wedge d \phi^{\prime} . \tag{6.2.10}
\end{align*}
$$

It is called the graded Chevalley-Eilenberg differential calculus over a graded commutative $\mathcal{K}$-ring $\mathcal{A}$. In particular, we have

$$
\begin{equation*}
\mathcal{O}^{1}[\mathfrak{d} \mathcal{A}]=\operatorname{Hom}_{\mathcal{A}}(\mathfrak{d} \mathcal{A}, \mathcal{A})=\mathfrak{o} \mathcal{A}^{*} \tag{6.2.11}
\end{equation*}
$$

One can extend this duality relation to the graded interior product of $u \in \mathcal{O} \mathcal{A}$ with any element $\phi \in \mathcal{O}^{*}[\mathfrak{O} \mathcal{A}]$ by the rules

$$
\begin{align*}
& u\rfloor(b d a)=(-1)^{[u][b]} u(a), \quad a, b \in \mathcal{A} \\
& \left.\left.u\rfloor\left(\phi \wedge \phi^{\prime}\right)=(u\rfloor \phi\right) \wedge \phi^{\prime}+(-1)^{|\phi|+[\phi][u]} \phi \wedge(u\rfloor \phi^{\prime}\right) . \tag{6.2.12}
\end{align*}
$$

As a consequence, any graded derivation $u \in \mathfrak{d} \mathcal{A}$ of $\mathcal{A}$ yields a derivation

$$
\begin{align*}
& \left.\left.\mathbf{L}_{u} \phi=u\right\rfloor d \phi+d(u\rfloor \phi\right), \quad \phi \in \mathcal{O}^{*}, \quad u \in \mathfrak{o} \mathcal{A}  \tag{6.2.13}\\
& \mathbf{L}_{u}\left(\phi \wedge \phi^{\prime}\right)=\mathbf{L}_{u}(\phi) \wedge \phi^{\prime}+(-1)^{[u][\phi]} \phi \wedge \mathbf{L}_{u}\left(\phi^{\prime}\right)
\end{align*}
$$

called the graded Lie derivative of the DBGA $\mathcal{O}^{*}[\mathfrak{O} \mathcal{A}]$.
The minimal graded Chevalley-Eilenberg differential calculus $\mathcal{O}^{*} \mathcal{A} \subset$ $\mathcal{O}^{*}[\mathfrak{O} \mathcal{A}]$ over a graded commutative ring $\mathcal{A}$ consists of the monomials

$$
a_{0} d a_{1} \wedge \cdots \wedge d a_{k}, \quad a_{i} \in \mathcal{A}
$$

The corresponding complex

$$
\begin{equation*}
0 \rightarrow \mathcal{K} \longrightarrow \mathcal{A} \xrightarrow{d} \mathcal{O}^{1} \mathcal{A} \xrightarrow{d} \cdots \mathcal{O}^{k} \mathcal{A} \xrightarrow{d} \cdots \tag{6.2.14}
\end{equation*}
$$

is called the bigraded de Rham complex of $\mathcal{A}$.
Following the construction of a connection in commutative geometry (see Section 8.2), one comes to the notion of a connection on modules over a real graded commutative $\operatorname{ring} \mathcal{A}$. The following are the straightforward counterparts of Definitions 8.2.3 and 8.2.4.

Definition 6.2.1: A connection on a graded $\mathcal{A}$-module $P$ is a graded $\mathcal{A}$-module morphism

$$
\begin{equation*}
\mathfrak{d} \mathcal{A} \ni u \rightarrow \nabla_{u} \in \operatorname{Diff}_{1}(P, P) \tag{6.2.15}
\end{equation*}
$$

such that the first order differential operators $\nabla_{u}$ obey the Grassmanngraded Leibniz rule

$$
\begin{equation*}
\nabla_{u}(a p)=u(a) p+(-1)^{[a][u]} a \nabla_{u}(p), \quad a \in \mathcal{A}, \quad p \in P . \tag{6.2.16}
\end{equation*}
$$

Definition 6.2.2: Let $P$ in Definition 6.2 .1 be a graded commutative $\mathcal{A}$-ring and $\mathfrak{o} P$ the derivation module of $P$ as a graded commutative $\mathcal{K}$-ring. A connection on a graded commutative $\mathcal{A}$-ring $P$ is a graded $\mathcal{A}$-module morphism

$$
\begin{equation*}
\mathfrak{d} \mathcal{A} \ni u \rightarrow \nabla_{u} \in \mathfrak{o} P, \tag{6.2.17}
\end{equation*}
$$

which is a connection on $P$ as an $\mathcal{A}$-module, i.e., it obeys the graded Leibniz rule (6.2.16).

### 6.3 Graded manifolds

A graded manifold of dimension $(n, m)$ is defined as a local-ringed space $(Z, \mathfrak{A})$ where $Z$ is an $n$-dimensional smooth manifold $Z$ and $\mathfrak{A}=\mathfrak{A}_{0} \oplus \mathfrak{A}_{1}$ is a sheaf of graded commutative algebras of rank $m$ such that:

- there is the exact sequence of sheaves

$$
\begin{equation*}
0 \rightarrow \mathcal{R} \rightarrow \mathfrak{A} \xrightarrow{\sigma} C_{Z}^{\infty} \rightarrow 0, \quad \mathcal{R}=\mathfrak{A}_{1}+\left(\mathfrak{A}_{1}\right)^{2} \tag{6.3.1}
\end{equation*}
$$

where $C_{Z}^{\infty}$ is the sheaf of smooth real functions on $Z$;

- $\mathcal{R} / \mathcal{R}^{2}$ is a locally free sheaf of $C_{Z}^{\infty}$-modules of finite rank (with respect to pointwise operations), and the sheaf $\mathfrak{A}$ is locally isomorphic to the exterior product $\wedge_{C_{Z}^{\infty}}\left(\mathcal{R} / \mathcal{R}^{2}\right)$.

The sheaf $\mathfrak{A}$ is called a structure sheaf of a graded manifold $(Z, \mathfrak{A})$, and a manifold $Z$ is said to be the body of $(Z, \mathfrak{A})$. Sections of the sheaf $\mathfrak{A}$ are called graded functions on a graded manifold $(Z, \mathfrak{A})$. They make up a graded commutative $C^{\infty}(Z)$-ring $\mathfrak{A}(Z)$ called the structure ring of $(Z, \mathfrak{A})$.

A graded manifold $(Z, \mathfrak{A})$ possesses the following local structure. Given a point $z \in Z$, there exists its open neighborhood $U$, called a splitting domain, such that

$$
\begin{equation*}
\mathfrak{A}(U)=C^{\infty}(U) \otimes \wedge \mathbb{R}^{m} \tag{6.3.2}
\end{equation*}
$$

This means that the restriction $\left.\mathfrak{A}\right|_{U}$ of the structure sheaf $\mathfrak{A}$ to $U$ is isomorphic to the sheaf $C_{U}^{\infty} \otimes \wedge \mathbb{R}^{m}$ of sections of some exterior bundle

$$
\wedge E_{U}^{*}=U \times \wedge \mathbb{R}^{m} \rightarrow U
$$

The well-known Batchelor theorem states that such a structure of a graded manifold is global as follows.

Theorem 6.3.1: Let $(Z, \mathfrak{A})$ be a graded manifold. There exists a vector bundle $E \rightarrow Z$ with an $m$-dimensional typical fibre $V$ such that the structure sheaf $\mathfrak{A}$ of $(Z, \mathfrak{A})$ is isomorphic to the structure sheaf $\mathfrak{A}_{E}=S_{\wedge E^{*}}$ of germs of sections of the exterior bundle $\wedge E^{*}$ (1.2.2), whose typical fibre is the Grassmann algebra $\wedge V^{*}$.

Note that Batchelor's isomorphism in Theorem 6.3.1 fails to be canonical. In field models, it however is fixed from the beginning. Therefore, we restrict our consideration to graded manifolds $\left(Z, \mathfrak{A}_{E}\right)$ whose structure sheaf is the sheaf of germs of sections of some exterior bundle $\wedge E^{*}$. We agree to call $\left(Z, \mathfrak{A}_{E}\right)$ a simple graded manifold modelled over a vector bundle $E \rightarrow Z$, called its characteristic vector bundle. Accordingly, the structure ring $\mathcal{A}_{E}$ of a simple graded manifold $\left(Z, \mathfrak{A}_{E}\right)$ is the structure module

$$
\begin{equation*}
\mathcal{A}_{E}=\mathfrak{A}_{E}(Z)=\wedge E^{*}(Z) \tag{6.3.3}
\end{equation*}
$$

of sections of the exterior bundle $\wedge E^{*}$. Automorphisms of a simple graded manifold $\left(Z, \mathfrak{A}_{E}\right)$ are restricted to those induced by automorphisms of its characteristic vector bundles $E \rightarrow Z$ (see Remark 6.3.2).

Combining Batchelor Theorem 6.3.1 and classical Serre-Swan Theorem 8.6.3, we come to the following Serre-Swan theorem for graded manifolds.

Theorem 6.3.2: Let $Z$ be a smooth manifold. A graded commutative $C^{\infty}(Z)$-algebra $\mathcal{A}$ is isomorphic to the structure ring of a graded manifold with a body $Z$ iff it is the exterior algebra of some projective $C^{\infty}(Z)$ module of finite rank.

Given a graded manifold $\left(Z, \mathfrak{A}_{E}\right)$, every trivialization chart $\left(U ; z^{A}, y^{a}\right)$ of the vector bundle $E \rightarrow Z$ yields a splitting domain $\left(U ; z^{A}, c^{a}\right)$ of $\left(Z, \mathfrak{A}_{E}\right)$. Graded functions on such a chart are $\Lambda$-valued functions

$$
\begin{equation*}
f=\sum_{k=0}^{m} \frac{1}{k!} f_{a_{1} \ldots a_{k}}(z) c^{a_{1}} \cdots c^{a_{k}} \tag{6.3.4}
\end{equation*}
$$

where $f_{a_{1} \cdots a_{k}}(z)$ are smooth functions on $U$ and $\left\{c^{a}\right\}$ is the fibre basis for $E^{*}$. In particular, the sheaf epimorphism $\sigma$ in (6.3.1) is induced by the body map of $\Lambda$. One calls $\left\{z^{A}, c^{a}\right\}$ the local basis for the graded manifold $\left(Z, \mathfrak{A}_{E}\right)$. Transition functions $y^{\prime a}=\rho_{b}^{a}\left(z^{A}\right) y^{b}$ of bundle coordinates on $E \rightarrow Z$ induce the corresponding transformation

$$
\begin{equation*}
c^{\prime a}=\rho_{b}^{a}\left(z^{A}\right) c^{b} \tag{6.3.5}
\end{equation*}
$$

of the associated local basis for the graded manifold $\left(Z, \mathfrak{A}_{E}\right)$ and the according coordinate transformation law of graded functions (6.3.4).

Remark 6.3.1: Strictly speaking, elements $c^{a}$ of the local basis for a graded manifold are locally constant sections $c^{a}$ of $E^{*} \rightarrow X$ such that $y_{b} \circ c^{a}=\delta_{b}^{a}$. Therefore, graded functions are locally represented by $\Lambda$-valued functions (6.3.4), but they are not $\Lambda$-valued functions on a manifold $Z$ because of the transformation law (6.3.5).

Remark 6.3.2: In general, automorphisms of a graded manifold read

$$
\begin{equation*}
c^{\prime a}=\rho^{a}\left(z^{A}, c^{b}\right) . \tag{6.3.6}
\end{equation*}
$$

Considering a simple graded manifold $\left(Z, \mathfrak{A}_{E}\right)$, we restrict the class of graded manifold transformations (6.3.6) to the linear ones (6.3.5), compatible with given Batchelor's isomorphism.

Let $E \rightarrow Z$ and $E^{\prime} \rightarrow Z$ be vector bundles and $\Phi: E \rightarrow E^{\prime}$ their bundle morphism over a morphism $\varphi: Z \rightarrow Z^{\prime}$. Then every section $s^{*}$ of the dual bundle $E^{*} \rightarrow Z^{\prime}$ defines the pull-back section $\Phi^{*} s^{*}$ of the dual bundle $E^{*} \rightarrow Z$ by the law

$$
\left.\left.v_{z}\right\rfloor \Phi^{*} s^{*}(z)=\Phi\left(v_{z}\right)\right\rfloor s^{*}(\varphi(z)), \quad v_{z} \in E_{z} .
$$

It follows that the bundle morphism $(\Phi, \varphi)$ yields a morphism of simple graded manifolds

$$
\begin{equation*}
\widehat{\Phi}:\left(Z, \mathfrak{A}_{E}\right) \rightarrow\left(Z^{\prime}, \mathfrak{A}_{E^{\prime}}\right) \tag{6.3.7}
\end{equation*}
$$

as local-ringed spaces. This is a pair $\left(\varphi, \varphi_{*} \circ \Phi^{*}\right)$ of a morphism $\varphi$ of body manifolds and the composition $\varphi_{*} \circ \Phi^{*}$ of the pull-back

$$
\mathcal{A}_{E^{\prime}} \ni f \rightarrow \Phi^{*} f \in \mathcal{A}_{E}
$$

of graded functions and the direct image $\varphi_{*}$ of the sheaf $\mathfrak{A}_{E}$ onto $Z^{\prime}$. Relative to local bases $\left(z^{A}, c^{a}\right)$ and $\left(z^{\prime A}, c^{\prime a}\right)$ for $\left(Z, \mathfrak{A}_{E}\right)$ and $\left(Z^{\prime}, \mathfrak{A}_{E^{\prime}}\right)$, the morphism (6.3.7) of graded manifolds reads

$$
\widehat{\Phi}(z)=\varphi(z), \quad \widehat{\Phi}\left(c^{\prime a}\right)=\Phi_{b}^{a}(z) c^{b}
$$

Given a graded manifold $(Z, \mathfrak{A})$, by the sheaf $\mathfrak{d A}$ of graded derivations of $\mathfrak{A}$ is meant a subsheaf of endomorphisms of the structure sheaf $\mathfrak{A}$ such that any section $u \in \mathfrak{d A}(U)$ of $\mathfrak{d A}$ over an open subset $U \subset Z$ is a graded derivation of the real graded commutative algebra $\mathfrak{A}(U)$, i.e., $u \in \mathfrak{d}(\mathfrak{A}(U))$. Conversely, one can show that, given open sets $U^{\prime} \subset U$, there is a surjection of the graded derivation modules

$$
\mathfrak{d}(\mathfrak{A}(U)) \rightarrow \mathfrak{d}\left(\mathfrak{A}\left(U^{\prime}\right)\right) .
$$

It follows that any graded derivation of the local graded algebra $\mathfrak{A}(U)$ also is a local section over $U$ of the sheaf $\mathfrak{o A}$. Global sections of $\mathfrak{d A}$ are called graded vector fields on the graded manifold $(Z, \mathfrak{A})$. They make up the graded derivation module $\mathfrak{d A}(Z)$ of the real graded commutative ring $\mathfrak{A}(Z)$. This module is a real Lie superalgebra with the superbracket (6.2.4).

A key point is that graded vector fields $u \in \mathfrak{o} \mathcal{A}_{E}$ on a simple graded manifold $\left(Z, \mathfrak{A}_{E}\right)$ can be represented by sections of some vector bundle as follows. Due to the canonical splitting $V E=E \times E$, the vertical tangent bundle $V E$ of $E \rightarrow Z$ can be provided with the fibre bases
$\left\{\partial / \partial c^{a}\right\}$, which are the duals of the bases $\left\{c^{a}\right\}$. Then graded vector fields on a splitting domain $\left(U ; z^{A}, c^{a}\right)$ of $\left(Z, \mathfrak{A}_{E}\right)$ read

$$
\begin{equation*}
u=u^{A} \partial_{A}+u^{a} \frac{\partial}{\partial c^{a}} \tag{6.3.8}
\end{equation*}
$$

where $u^{\lambda}, u^{a}$ are local graded functions on $U$. In particular,

$$
\frac{\partial}{\partial c^{a}} \circ \frac{\partial}{\partial c^{b}}=-\frac{\partial}{\partial c^{b}} \circ \frac{\partial}{\partial c^{a}}, \quad \partial_{A} \circ \frac{\partial}{\partial c^{a}}=\frac{\partial}{\partial c^{a}} \circ \partial_{A}
$$

Graded derivations (6.3.8) act on graded functions $f$ (6.3.4) by the rule

$$
\begin{equation*}
\left.u\left(f_{a \ldots b} c^{a} \cdots c^{b}\right)=u^{A} \partial_{A}\left(f_{a \ldots b}\right) c^{a} \cdots c^{b}+u^{k} f_{a \ldots b} \frac{\partial}{\partial c^{k}}\right\rfloor\left(c^{a} \cdots c^{b}\right) \tag{6.3.9}
\end{equation*}
$$

This rule implies the corresponding coordinate transformation law

$$
u^{\prime A}=u^{A}, \quad u^{\prime a}=\rho_{j}^{a} u^{j}+u^{A} \partial_{A}\left(\rho_{j}^{a}\right) c^{j}
$$

of graded vector fields. It follows that graded vector fields (6.3.8) can be represented by sections of the following vector bundle $\mathcal{V}_{E} \rightarrow Z$. This vector bundle is locally isomorphic to the vector bundle

$$
\begin{equation*}
\left.\left.\mathcal{V}_{E}\right|_{U} \approx \wedge E^{*} \underset{Z}{\otimes}(E \underset{Z}{\oplus} T Z)\right|_{U} \tag{6.3.10}
\end{equation*}
$$

and is characterized by an atlas of bundle coordinates

$$
\left(z^{A}, z_{a_{1} \ldots a_{k}}^{A}, v_{b_{1} \ldots b_{k}}^{i}\right), \quad k=0, \ldots, m
$$

possessing the transition functions

$$
\begin{aligned}
& z_{i_{1} \ldots i_{k}}^{\prime A}=\rho_{i_{1}}^{-1 a_{1}} \cdots \rho^{-1 a_{k}} z_{i_{k} \ldots a_{k}}^{A}, \\
& v_{j_{1} \ldots j_{k}}^{\prime i}=\rho^{-1 b_{1}} \cdots \rho_{j_{1}}^{-1 b_{k}}\left[\rho_{j}^{i} v_{b_{1} \ldots b_{k}}^{j}+\frac{k!}{(k-1)!} z_{b_{1} \ldots b_{k-1}}^{A} \partial_{A} \rho_{b_{k}}^{i}\right],
\end{aligned}
$$

which fulfil the cocycle condition (1.1.4). Thus, the graded derivation module $\mathfrak{d} \mathcal{A}_{E}$ is isomorphic to the structure module $\mathcal{V}_{E}(Z)$ of global sections of the vector bundle $\mathcal{V}_{E} \rightarrow Z$.

There is the exact sequence

$$
\begin{equation*}
0 \rightarrow \wedge E^{*}{\underset{Z}{\otimes}}_{\otimes} E \rightarrow \mathcal{V}_{E} \rightarrow \wedge E^{*}{\underset{Z}{\mid}}_{\otimes} T Z \rightarrow 0 \tag{6.3.11}
\end{equation*}
$$

of vector bundles over $Z$. Its splitting

$$
\begin{equation*}
\tilde{\gamma}: \dot{z}^{A} \partial_{A} \rightarrow \dot{z}^{A}\left(\partial_{A}+\tilde{\gamma}_{A}^{a} \frac{\partial}{\partial c^{a}}\right) \tag{6.3.12}
\end{equation*}
$$

transforms every vector field $\tau$ on $Z$ into the graded vector field

$$
\begin{equation*}
\tau=\tau^{A} \partial_{A} \rightarrow \nabla_{\tau}=\tau^{A}\left(\partial_{A}+\tilde{\gamma}_{A}^{a} \frac{\partial}{\partial c^{a}}\right), \tag{6.3.13}
\end{equation*}
$$

which is a graded derivation of the real graded commutative ring $\mathcal{A}_{E}$ (6.3.3) satisfying the Leibniz rule

$$
\left.\nabla_{\tau}(s f)=(\tau\rfloor d s\right) f+s \nabla_{\tau}(f), \quad f \in \mathcal{A}_{E}, \quad s \in C^{\infty}(Z) .
$$

It follows that the splitting (6.3.12) of the exact sequence (6.3.11) yields a connection on the graded commutative $C^{\infty}(Z)$-ring $\mathcal{A}_{E}$ in accordance with Definition 6.2.2. It is called a graded connection on the simple graded manifold $\left(Z, \mathfrak{A}_{E}\right)$. In particular, this connection provides the corresponding horizontal splitting

$$
u=u^{A} \partial_{A}+u^{a} \frac{\partial}{\partial c^{a}}=u^{A}\left(\partial_{A}+\tilde{\gamma}_{A}^{a} \frac{\partial}{\partial c^{a}}\right)+\left(u^{a}-u^{A} \tilde{\gamma}_{A}^{a}\right) \frac{\partial}{\partial c^{a}}
$$

of graded vector fields. In accordance with Theorem 1.2.2, a graded connection (6.3.12) always exists.

Remark 6.3.3: By virtue of the isomorphism (6.3.2), any connection $\tilde{\gamma}$ on a graded manifold $(Z, \mathfrak{A})$, restricted to a splitting domain $U$, takes the form (6.3.12). Given two splitting domains $U$ and $U^{\prime}$ of $(Z, \mathfrak{A})$ with the transition functions (6.3.6), the connection components $\tilde{\gamma}_{A}^{a}$ obey the transformation law

$$
\begin{equation*}
\tilde{\gamma}_{A}^{\prime a}=\tilde{\gamma}_{A}^{b} \frac{\partial}{\partial c^{b}} \rho^{a}+\partial_{A} \rho^{a} . \tag{6.3.14}
\end{equation*}
$$

If $U$ and $U^{\prime}$ are the trivialization charts of the same vector bundle $E$ in Theorem 6.3.1 together with the transition functions (6.3.5), the transformation law (6.3.14) takes the form

$$
\begin{equation*}
\tilde{\gamma}_{A}^{\prime a}=\rho_{b}^{a}(z) \tilde{\gamma}_{A}^{b}+\partial_{A} \rho_{b}^{a}(z) c^{b} . \tag{6.3.15}
\end{equation*}
$$

Remark 6.3.4: Every linear connection

$$
\gamma=d z^{A} \otimes\left(\partial_{A}+\gamma_{A}{ }^{a}{ }_{b} y^{b} \partial_{a}\right)
$$

on a vector bundle $E \rightarrow Z$ yields the graded connection

$$
\begin{equation*}
\gamma_{S}=d z^{A} \otimes\left(\partial_{A}+\gamma_{A}{ }^{a}{ }_{b} c^{b} \frac{\partial}{\partial c^{a}}\right) \tag{6.3.16}
\end{equation*}
$$

on the simple graded manifold $\left(Z, \mathfrak{A}_{E}\right)$ modelled over $E$. In view of Remark 6.3.3, $\gamma_{S}$ also is a graded connection on the graded manifold $(Z, \mathfrak{A}) \cong\left(Z, \mathfrak{A}_{E}\right)$, but its linear form (6.3.16) is not maintained under the transformation law (6.3.14).

### 6.4 Graded differential forms

Given the structure ring $\mathcal{A}_{E}$ of graded functions on a simple graded manifold $\left(Z, \mathfrak{A}_{E}\right)$ and the real Lie superalgebra $\mathfrak{d} \mathcal{A}_{E}$ of its graded derivations, let us consider the graded Chevalley-Eilenberg differential calculus

$$
\begin{equation*}
\mathcal{S}^{*}[E ; Z]=\mathcal{O}^{*}\left[\mathfrak{d} \mathcal{A}_{E}\right] \tag{6.4.1}
\end{equation*}
$$

over $\mathcal{A}_{E}$. Since the graded derivation module $\mathfrak{o} \mathcal{A}_{E}$ is isomorphic to the structure module of sections of the vector bundle $\mathcal{V}_{E} \rightarrow Z$, elements of $\mathcal{S}^{*}[E ; Z]$ are sections of the exterior bundle $\wedge \overline{\mathcal{V}}_{E}$ of the $\mathcal{A}_{E}$-dual $\overline{\mathcal{V}}_{E} \rightarrow Z$ of $\mathcal{V}_{E}$. The bundle $\overline{\mathcal{V}}_{E}$ is locally isomorphic to the vector bundle

$$
\begin{equation*}
\left.\left.\overline{\mathcal{V}}_{E}\right|_{U} \approx\left(E^{*} \underset{Z}{\oplus} T^{*} Z\right)\right|_{U} \tag{6.4.2}
\end{equation*}
$$

With respect to the dual fibre bases $\left\{d z^{A}\right\}$ for $T^{*} Z$ and $\left\{d c^{b}\right\}$ for $E^{*}$, sections of $\overline{\mathcal{V}}_{E}$ take the coordinate form

$$
\phi=\phi_{A} d z^{A}+\phi_{a} d c^{a}
$$

together with transition functions

$$
\phi_{a}^{\prime}=\rho_{a}^{-1 b} \phi_{b}, \quad \phi_{A}^{\prime}=\phi_{A}+\rho_{a}^{-1 b} \partial_{A}\left(\rho_{j}^{a}\right) \phi_{b} c^{j} .
$$

The duality isomorphism $\mathcal{S}^{1}[E ; Z]=\mathcal{D} \mathcal{A}_{E}^{*}$ (6.2.11) is given by the graded interior product

$$
\begin{equation*}
u\rfloor \phi=u^{A} \phi_{A}+(-1)^{\left[\phi_{a}\right]} u^{a} \phi_{a} . \tag{6.4.3}
\end{equation*}
$$

Elements of $\mathcal{S}^{*}[E ; Z]$ are called graded exterior forms on the graded manifold $\left(Z, \mathfrak{A}_{E}\right)$.

Seen as an $\mathcal{A}_{E}$-algebra, the DBGA $\mathcal{S}^{*}[E ; Z]$ (6.4.1) on a splitting domain $\left(U ; z^{A}, c^{a}\right)$ is locally generated by the graded one-forms $d z^{A}, d c^{i}$ such that

$$
\begin{equation*}
d z^{A} \wedge d c^{i}=-d c^{i} \wedge d z^{A}, \quad d c^{i} \wedge d c^{j}=d c^{j} \wedge d c^{i} \tag{6.4.4}
\end{equation*}
$$

Accordingly a graded Chevalley-Eilenberg coboundary operator $d$ (6.2.7), called the graded exterior differential, reads

$$
d \phi=d z^{A} \wedge \partial_{A} \phi+d c^{a} \wedge \frac{\partial}{\partial c^{a}} \phi,
$$

where the derivatives $\partial_{\lambda}, \partial / \partial c^{a}$ act on coefficients of graded exterior forms by the formula (6.3.9), and they are graded commutative with the graded forms $d z^{A}$ and $d c^{a}$. The formulas (6.2.9) - (6.2.13) hold.

Theorem 6.4.1: The DBGA $\mathcal{S}^{*}[E ; Z]$ (6.4.1) is a minimal differential calculus over $\mathcal{A}_{E}$, i.e., it is generated by elements $d f, f \in \mathcal{A}_{E}$.

The bigraded de Rham complex (6.2.14) of the minimal ChevalleyEilenberg differential calculus $\mathcal{S}^{*}[E ; Z]$ reads

$$
\begin{equation*}
0 \rightarrow \mathbb{R} \rightarrow \mathcal{A}_{E} \xrightarrow{d} \mathcal{S}^{1}[E ; Z] \xrightarrow{d} \cdots \mathcal{S}^{k}[E ; Z] \xrightarrow{d} \cdots . \tag{6.4.5}
\end{equation*}
$$

Its cohomology $H^{*}\left(\mathcal{A}_{E}\right)$ is called the de Rham cohomology of a simple graded manifold $\left(Z, \mathfrak{A}_{E}\right)$.

In particular, given the DGA $\mathcal{O}^{*}(Z)$ of exterior forms on $Z$, there exist the canonical monomorphism

$$
\begin{equation*}
\mathcal{O}^{*}(Z) \rightarrow \mathcal{S}^{*}[E ; Z] \tag{6.4.6}
\end{equation*}
$$

and the body epimorphism $\mathcal{S}^{*}[E ; Z] \rightarrow \mathcal{O}^{*}(Z)$ which are cochain morphisms of the de Rham complexes (6.4.5) and (8.6.5).

Theorem 6.4.2: The de Rham cohomology of a simple graded manifold $\left(Z, \mathfrak{A}_{E}\right)$ equals the de Rham cohomology of its body $Z$.

Corollary 6.4.3: Any closed graded exterior form is decomposed into a sum $\phi=\sigma+d \xi$ where $\sigma$ is a closed exterior form on $Z$.

## Chapter 7

## Lagrangian theory

Lagrangian theory on fibre bundles is algebraically formulated in terms of the variational bicomplex without appealing to the calculus of variations. This formulation is extended to Lagrangian theory on graded manifolds [9, 24].

### 7.1 Variational bicomplex

Let $Y \rightarrow X$ be a fibre bundle. The DGA $\mathcal{O}_{\infty}^{*}(2.4 .6)$, decomposed into the variational bicomplex, describes finite order Lagrangian theories on $Y \rightarrow X$. One also considers the variational bicomplex of the DGA $\mathcal{Q}_{\infty}^{*}$ (2.4.8) and different variants of the variational sequence of finite jet order.

In order to transform the bicomplex $\mathcal{O}_{\infty}^{*, *}$ into the variational one, let us consider the following two operators acting on $\mathcal{O}_{\infty}^{*, n}$.
(i) There exists an $\mathbb{R}$-module endomorphism

$$
\begin{align*}
& \varrho=\sum_{k>0} \frac{1}{k} \bar{\varrho} \circ h_{k} \circ h^{n}: \mathcal{O}_{\infty}^{*>0, n} \rightarrow \mathcal{O}_{\infty}^{*>0, n},  \tag{7.1.1}\\
& \left.\bar{\varrho}(\phi)=\sum_{0 \leq|\Lambda|}(-1)^{|\Lambda|} \theta^{i} \wedge\left[d_{\Lambda}\left(\partial_{i}^{\Lambda}\right\rfloor \phi\right)\right], \quad \phi \in \mathcal{O}_{\infty}^{>0, n},
\end{align*}
$$

possessing the following properties.
Lemma 7.1.1: For any $\phi \in \mathcal{O}_{\infty}^{>0, n}$, the form $\phi-\varrho(\phi)$ is locally $d_{H}$-exact on each coordinate chart (2.4.3).

Lemma 7.1.2: The operator $\varrho$ obeys the relation

$$
\begin{equation*}
\left(\varrho \circ d_{H}\right)(\psi)=0, \quad \psi \in \mathcal{O}_{\infty}^{>0, n-1} \tag{7.1.2}
\end{equation*}
$$

It follows from Lemmas 7.1.1 and 7.1.2 that $\varrho(7.1 .1)$ is a projector.
(ii) One defines the variational operator

$$
\begin{equation*}
\delta=\varrho \circ d: \mathcal{O}_{\infty}^{*, n} \rightarrow \mathcal{O}_{\infty}^{*+1, n} \tag{7.1.3}
\end{equation*}
$$

which is nilpotent, i.e., $\delta \circ \delta=0$, and obeys the relation $\delta \circ \varrho=\delta$.
Let us denote $\mathbf{E}_{k}=\varrho\left(\mathcal{O}_{\infty}^{k, n}\right)$. Provided with the operators $d_{H}, d_{V}, \varrho$ and $\delta$, the $\mathrm{DGA} \mathcal{O}_{\infty}^{*}$ is decomposed into the variational bicomplex

$$
\begin{aligned}
& 0 \rightarrow \mathbb{R} \rightarrow \quad \mathcal{O}_{\infty}^{0} \quad \xrightarrow{d_{V} \uparrow} \quad{ }^{d_{H}} \mathcal{O}_{\infty}^{0,1} \quad \xrightarrow{d_{H}} \cdots \quad \mathcal{O}_{\infty}^{d_{V} \uparrow} \quad \equiv{ }^{-\delta \uparrow} \mathcal{O}_{\infty}^{0, n} \\
& 0 \rightarrow \mathbb{R} \rightarrow \quad{ }^{\uparrow} \mathcal{O}^{0}(X) \xrightarrow{d}{ }^{\uparrow} \mathcal{O}^{1}(X) \xrightarrow{d} \cdots \quad \mathcal{O}^{n}(X) \xrightarrow{d} 0 \\
& 0 \\
& \begin{array}{ll}
\uparrow & \uparrow \\
0 & 0
\end{array}
\end{aligned}
$$

It possesses the following cohomology [20, 24].
Theorem 7.1.3: The second row from the bottom and the last column of the variational bicomplex (7.1.4) make up the variational complex

$$
\begin{equation*}
0 \rightarrow \mathbb{R} \rightarrow \mathcal{O}_{\infty}^{0} \xrightarrow{d_{H}} \mathcal{O}_{\infty}^{0,1} \cdots \xrightarrow{d_{H}} \mathcal{O}_{\infty}^{0, n} \xrightarrow{\delta} \mathbf{E}_{1} \xrightarrow{\delta} \mathbf{E}_{2} \longrightarrow \cdots \tag{7.1.5}
\end{equation*}
$$

Its cohomology is isomorphic to the de Rham cohomology of a fibre bundle $Y$, namely,

$$
\begin{equation*}
H^{k<n}\left(d_{H} ; \mathcal{O}_{\infty}^{*}\right)=H_{\mathrm{DR}}^{k<n}(Y), \quad H^{k \geq n}\left(\delta ; \mathcal{O}_{\infty}^{*}\right)=H_{\mathrm{DR}}^{k \geq n}(Y) \tag{7.1.6}
\end{equation*}
$$

Theorem 7.1.4: The rows of contact forms of the variational bicomplex (7.1.4) are exact sequences.

The cohomology isomorphism (7.1.6) gives something more. Due to the relations $d_{H} \circ h_{0}=h_{0} \circ d$ and $\delta \circ \varrho=\delta$, we have the cochain morphism of the de Rham complex (2.4.7) of the DGA $\mathcal{O}_{\infty}^{*}$ to its variational complex (7.1.5). By virtue of Theorems 2.4.3 and 7.1.3, the corresponding homomorphism of their cohomology groups is an isomorphism. Then the splitting of a closed form $\phi \in \mathcal{O}_{\infty}^{*}$ in Corollary 2.4.4 leads to the following decompositions.

Theorem 7.1.5: Any $d_{H}$-closed form $\phi \in \mathcal{O}^{0, m}, m<n$, is represented by a sum

$$
\begin{equation*}
\phi=h_{0} \sigma+d_{H} \xi, \quad \xi \in \mathcal{O}_{\infty}^{m-1} \tag{7.1.7}
\end{equation*}
$$

where $\sigma$ is a closed $m$-form on $Y$. Any $\delta$-closed form $\phi \in \mathcal{O}^{k, n}$ is split into

$$
\begin{array}{lc}
\phi=h_{0} \sigma+d_{H} \xi, & k=0, \\
\phi=\varrho(\sigma)+\delta(\xi), & k=1, \\
\phi=\varrho(\sigma) \mathcal{O}_{\infty}^{0, n-1}  \tag{7.1.10}\\
\phi=\mathcal{O}_{\infty}^{0, n} \\
\phi(\xi), & k>1,
\end{array}, \xi \in \mathbf{E}_{k-1},
$$

where $\sigma$ is a closed $(n+k)$-form on $Y$.

### 7.2 Lagrangian theory on fibre bundles

In Lagrangian formalism on fibre bundles, a finite order Lagrangian and its Euler-Lagrange operator are defined as elements

$$
\begin{align*}
& L=\mathcal{L} \omega \in \mathcal{O}_{\infty}^{0, n}  \tag{7.2.1}\\
& \delta L=\mathcal{E}_{L}=\mathcal{E}_{i} \theta^{i} \wedge \omega \in \mathbf{E}_{1},  \tag{7.2.2}\\
& \mathcal{E}_{i}=\sum_{0 \leq|\Lambda|}(-1)^{|\Lambda|} d_{\Lambda}\left(\partial_{i}^{\Lambda} \mathcal{L}\right), \tag{7.2.3}
\end{align*}
$$

of the variational complex (7.1.5) (see the notation (1.4.1)). Components $\mathcal{E}_{i}$ (7.2.3) of the Euler-Lagrange operator (7.2.2) are called the variational derivatives. Elements of $\mathbf{E}_{1}$ are called the Euler-Lagrangetype operators.

Hereafter, we call a pair $\left(\mathcal{O}_{\infty}^{*}, L\right)$ the Lagrangian system. The following are corollaries of Theorem 7.1.5.

Corollary 7.2.1: A finite order Lagrangian $L$ (7.2.1) is variationally trivial, i.e., $\delta(L)=0$ iff

$$
\begin{equation*}
L=h_{0} \sigma+d_{H} \xi, \quad \xi \in \mathcal{O}_{\infty}^{0, n-1}, \tag{7.2.4}
\end{equation*}
$$

where $\sigma$ is a closed $n$-form on $Y$.
Corollary 7.2.2: A finite order Euler-Lagrange-type operator $\mathcal{E} \in \mathbf{E}_{1}$ satisfies the Helmholtz condition $\delta(\mathcal{E})=0$ iff

$$
\mathcal{E}=\delta L+\varrho(\sigma), \quad L \in \mathcal{O}_{\infty}^{0, n}
$$

where $\sigma$ is a closed $(n+1)$-form on $Y$.
A glance at the expression (7.2.2) shows that, if a Lagrangian $L$ (7.2.1) is of $r$-order, its Euler-Lagrange operator $\mathcal{E}_{L}$ is of $2 r$-order. Its kernel is called the Euler-Lagrange equation. Euler-Lagrange equations traditionally came from the variational formula

$$
\begin{equation*}
d L=\delta L-d_{H} \Xi_{L} \tag{7.2.5}
\end{equation*}
$$

of the calculus of variations. In formalism of the variational bicomplex, this formula is a corollary of Theorem 7.1.4.

Corollary 7.2.3: The exactness of the row of one-contact forms of the variational bicomplex (7.1.4) at the term $\mathcal{O}_{\infty}^{1, n}$ relative to the projector $\varrho$ provides the $\mathbb{R}$-module decomposition

$$
\mathcal{O}_{\infty}^{1, n}=\mathbf{E}_{1} \oplus d_{H}\left(\mathcal{O}_{\infty}^{1, n-1}\right) .
$$

In particular, any Lagrangian $L$ admits the decomposition (7.2.5).

Defined up to a $d_{H}$-closed term, a form $\Xi_{L} \in \mathcal{O}_{\infty}^{n}$ in the variational formula (7.2.5) reads

$$
\begin{align*}
& \Xi_{L}=L+\left[\left(\partial_{i}^{\lambda} \mathcal{L}-d_{\mu} F_{i}^{\mu \lambda}\right) \theta^{i}+\sum_{s=1} F_{i}^{\lambda \nu_{s} \ldots \nu_{1}} \theta_{\nu_{s} \ldots \nu_{1}}^{i}\right] \wedge \omega_{\lambda},  \tag{7.2.6}\\
& F_{i}^{\nu_{k} \ldots \nu_{1}}=\partial_{i}^{\nu_{k} \ldots \nu_{1}} \mathcal{L}-d_{\mu} F_{i}^{\mu \nu_{k} \ldots \nu_{1}}+\sigma_{i}^{\nu_{k} \ldots \nu_{1}}, \quad k=2,3, \ldots,
\end{align*}
$$

where $\sigma_{i}^{\nu_{k} \ldots \nu_{1}}$ are local functions such that

$$
\sigma_{i}^{\left(\nu_{k} \nu_{k-1}\right) \ldots \nu_{1}}=0
$$

The form $\Xi_{L}(7.2 .6)$ possesses the following properties:

- $h_{0}\left(\Xi_{L}\right)=L$,
- $\left.h_{0}(\vartheta\rfloor d \Xi_{L}\right)=\vartheta^{i} \mathcal{E}_{i} \omega$ for any derivation $\vartheta(2.4 .12)$.

Consequently, $\Xi_{L}$ is a Lepage equivalent of a Lagrangian $L$.
A special interest is concerned with Lagrangian theories on an affine bundle $Y \rightarrow X$. Since $X$ is a strong deformation retract of an affine bundle $Y$, the de Rham cohomology of $Y$ equals that of $X$. In this case, the cohomology (7.1.6) of the variational complex (7.1.5) equals the de Rham cohomology of $X$, namely,

$$
\begin{align*}
& H^{<n}\left(d_{H} ; \mathcal{O}_{\infty}^{*}\right)=H_{\mathrm{DR}}^{<n}(X), \\
& H^{n}\left(\delta ; \mathcal{O}_{\infty}^{*}\right)=H_{\mathrm{DR}}^{n}(X),  \tag{7.2.7}\\
& H^{>n}\left(\delta ; \mathcal{O}_{\infty}^{*}\right)=0
\end{align*}
$$

It follows that every $d_{H}$-closed form $\phi \in \mathcal{O}_{\infty}^{0, m<n}$ is represented by the sum

$$
\begin{equation*}
\phi=\sigma+d_{H} \xi, \quad \xi \in \mathcal{O}_{\infty}^{0, m-1} \tag{7.2.8}
\end{equation*}
$$

where $\sigma$ is a closed $m$-form on $X$. Similarly, any variationally trivial Lagrangian takes the form

$$
\begin{equation*}
L=\sigma+d_{H} \xi, \quad \xi \in \mathcal{O}_{\infty}^{0, n-1} \tag{7.2.9}
\end{equation*}
$$

where $\sigma$ is an $n$-form on $X$.

In view of the cohomology isomorphism (7.2.7), if $Y \rightarrow X$ is an affine bundle, let us restrict our consideration to the short variational complex

$$
\begin{equation*}
0 \rightarrow \mathbb{R} \rightarrow \mathcal{O}_{\infty}^{0} \xrightarrow{d_{H}} \mathcal{O}_{\infty}^{0,1} \cdots \xrightarrow{d_{H}} \mathcal{O}_{\infty}^{0, n} \xrightarrow{\delta} \mathbf{E}_{1}, \tag{7.2.10}
\end{equation*}
$$

whose non-trivial cohomology equals that of the variational complex (7.1.5). Let us consider a DGA $\mathcal{P}_{\infty}^{*} \subset \mathcal{O}_{\infty}^{*}$ of exterior forms whose coefficients are polynomials in jet coordinates $y_{\Lambda}^{i}, 0 \leq|\Lambda|$, of the continuous bundle $J^{\infty} Y \rightarrow X$.

Theorem 7.2.4: The cohomology of the short variational complex

$$
\begin{equation*}
0 \rightarrow \mathbb{R} \rightarrow \mathcal{P}_{\infty}^{0} \xrightarrow{d_{H}} \mathcal{P}_{\infty}^{0,1} \cdots \xrightarrow{d_{H}} \mathcal{P}_{\infty}^{0, n} \xrightarrow{\delta} 0 \tag{7.2.11}
\end{equation*}
$$

of the polynomial algebra $\mathcal{P}_{\infty}^{*}$ equals that of the complex (7.2.10), i.e., the de Rham cohomology of $X$.

Given a Lagrangian system $\left(\mathcal{O}_{\infty}^{*}, L\right)$, its infinitesimal transformations are defined to be contact derivations of the ring $\mathcal{O}_{\infty}^{0}$.

A derivation $\vartheta \in \mathfrak{d} \mathcal{O}_{\infty}^{0}$ (2.4.12) is called contact if the Lie derivative $\mathbf{L}_{v}$ preserves the ideal of contact forms of the DGA $\mathcal{O}_{\infty}^{*}$, i.e., the Lie derivative $\mathbf{L}_{v}$ of a contact form is a contact form.

Lemma 7.2.5: A derivation $\vartheta(2.4 .12)$ is contact iff it takes the form

$$
\begin{equation*}
\vartheta=v^{\lambda} \partial_{\lambda}+v^{i} \partial_{i}+\sum_{0<|\Lambda|}\left[d_{\Lambda}\left(v^{i}-y_{\mu}^{i} v^{\mu}\right)+y_{\mu+\Lambda}^{i} v^{\mu}\right] \partial_{i}^{\Lambda} \tag{7.2.12}
\end{equation*}
$$

The expression (2.2.8) enables one to regard a contact derivation $\vartheta$ (7.2.12) as an infinite order jet prolongation $\vartheta=J^{\infty} v$ of its restriction

$$
\begin{equation*}
v=v^{\lambda} \partial_{\lambda}+v^{i} \partial_{i} \tag{7.2.13}
\end{equation*}
$$

to the ring $C^{\infty}(Y)$. Since coefficients $v^{\lambda}$ and $v^{i}$ of $v$ (7.2.13) depend generally on jet coordinates $y_{\Lambda}^{i}, 0<|\Lambda|$, one calls $v$ (7.2.13) a generalized vector field. It can be represented as a section of some pull-back bundle

$$
J^{r} Y \underset{Y}{\times} T Y \rightarrow J^{r} Y
$$

A contact derivation $\vartheta$ (7.2.12) is called projectable, if the generalized vector field $v(7.2 .13)$ projects onto a vector field $v^{\lambda} \partial_{\lambda}$ on $X$.

Any contact derivation $\vartheta$ (7.2.12) admits the horizontal splitting

$$
\begin{align*}
& \vartheta=\vartheta_{H}+\vartheta_{V}=v^{\lambda} d_{\lambda}+\left[v_{V}^{i} \partial_{i}+\sum_{0<|\Lambda|} d_{\Lambda} v_{V}^{i} \partial_{i}^{\Lambda}\right],  \tag{7.2.14}\\
& v=v_{H}+v_{V}=v^{\lambda} d_{\lambda}+\left(v^{i}-y_{\mu}^{i} v^{\mu}\right) \partial_{i}, \tag{7.2.15}
\end{align*}
$$

relative to the canonical connection $\nabla(2.4 .14)$ on the $C^{\infty}(X)$-ring $\mathcal{O}_{\infty}^{0}$.
Lemma 7.2.6: Any vertical contact derivation

$$
\begin{equation*}
\vartheta=v^{i} \partial_{i}+\sum_{0<|\Lambda|} d_{\Lambda} v^{i} \partial_{i}^{\Lambda} \tag{7.2.16}
\end{equation*}
$$

obeys the relations

$$
\begin{align*}
& \left.\vartheta\rfloor d_{H} \phi=-d_{H}(\vartheta\rfloor \phi\right),  \tag{7.2.17}\\
& \mathbf{L}_{\vartheta}\left(d_{H} \phi\right)=d_{H}\left(\mathbf{L}_{\vartheta} \phi\right), \quad \phi \in \mathcal{O}_{\infty}^{*} . \tag{7.2.18}
\end{align*}
$$

The global decomposition (7.2.5) leads to the following first variational formula (Theorem 7.2.7) and the first Noether theorem (Theorem 7.2.9).

Theorem 7.2.7: Given a Lagrangian $L \in \mathcal{O}_{\infty}^{0, n}$, its Lie derivative $\mathbf{L}_{v} L$ along a contact derivation $v$ (7.2.14) fulfils the first variational formula

$$
\begin{equation*}
\left.\left.\left.\mathbf{L}_{\vartheta} L=v_{V}\right\rfloor \delta L+d_{H}\left(h_{0}(\vartheta\rfloor \Xi_{L}\right)\right)+\mathcal{L} d_{V}\left(v_{H}\right\rfloor \omega\right), \tag{7.2.19}
\end{equation*}
$$

where $\Xi_{L}$ is the Lepage equivalent $(7.2 .6)$ of $L$.
A contact derivation $\vartheta(7.2 .12)$ is called a variational symmetry of a Lagrangian $L$ if the Lie derivative $\mathbf{L}_{\vartheta} L$ is $d_{H}$-exact, i.e.,

$$
\begin{equation*}
\mathbf{L}_{\vartheta} L=d_{H} \sigma . \tag{7.2.20}
\end{equation*}
$$

Lemma 7.2.8: A glance at the expression (7.2.19) shows the following.
(i) A contact derivation $\vartheta$ is a variational symmetry only if it is projectable.
(ii) Any projectable contact derivation is a variational symmetry of a variationally trivial Lagrangian. It follows that, if $\vartheta$ is a variational symmetry of a Lagrangian $L$, it also is a variational symmetry of a Lagrangian $L+L_{0}$, where $L_{0}$ is a variationally trivial Lagrangian.
(iii) A projectable contact derivations $\vartheta$ is a variational symmetry iff its vertical part $v_{V}(7.2 .14)$ is well.
(iv) A projectable contact derivations $\vartheta$ is a variational symmetry iff the density $\left.v_{V}\right\rfloor \delta L$ is $d_{H^{-}}$-exact.

It is readily observed that variational symmetries of a Lagrangian $L$ constitute a real vector subspace $\mathcal{G}_{L}$ of the derivation module $\mathfrak{d} \mathcal{O}_{\infty}^{0}$. By virtue of item (ii) of Lemma 7.2.8, the Lie bracket

$$
\mathbf{L}_{\left[\vartheta, v^{\prime}\right]}=\left[\mathbf{L}_{\vartheta}, \mathbf{L}_{\vartheta^{\prime}}\right]
$$

of variational symmetries is a variational symmetry and, therefore, their vector space $\mathcal{G}_{L}$ is a real Lie algebra. The following is the first Noether theorem.

Theorem 7.2.9: If a contact derivation $\vartheta(7.2 .12)$ is a variational symmetry (7.2.20) of a Lagrangian $L$, the first variational formula (7.2.19) restricted to the kernel of the Euler-Lagrange operator $\operatorname{Ker} \mathcal{E}_{L}$ leads to the weak conservation law

$$
\begin{equation*}
\left.0 \approx d_{H}\left(h_{0}(\vartheta\rfloor \Xi_{L}\right)-\sigma\right) \tag{7.2.21}
\end{equation*}
$$

on the shell $\delta L=0$.
A variational symmetry $\vartheta$ of a Lagrangian $L$ is called its exact symmetry or, simply, a symmetry if

$$
\begin{equation*}
\mathbf{L}_{\vartheta} L=0 . \tag{7.2.22}
\end{equation*}
$$

Symmetries of a Lagrangian $L$ constitute a real vector space, which is a real Lie algebra. Its vertical symmetries $v(7.2 .16)$ obey the relation

$$
\left.\mathbf{L}_{v} L=v\right\rfloor d L
$$

and, therefore, make up a $\mathcal{O}_{\infty}^{0}$-module which is a Lie $C^{\infty}(X)$-algebra.
If $\vartheta$ is an exact symmetry of a Lagrangian $L$, the weak conservation law (7.2.21) takes the form

$$
\begin{equation*}
\left.0 \approx d_{H}\left(h_{0}(\vartheta\rfloor \Xi_{L}\right)\right)=-d_{H} \mathcal{J}_{v} \tag{7.2.23}
\end{equation*}
$$

where

$$
\begin{equation*}
\left.\mathcal{J}_{v}=\mathcal{J}_{v}^{\mu} \omega_{\mu}=-h_{0}(\vartheta\rfloor \Xi_{L}\right) \tag{7.2.24}
\end{equation*}
$$

is called the symmetry current. Of course, the symmetry current (7.2.24) is defined with the accuracy of a $d_{H}$-closed term.

Let $\vartheta$ be an exact symmetry of a Lagrangian $L$. Whenever $L_{0}$ is a variationally trivial Lagrangian, $\vartheta$ is a variational symmetry of the Lagrangian $L+L_{0}$ such that the weak conservation law (7.2.21) for this Lagrangian is reduced to the weak conservation law (7.2.23) for a Lagrangian $L$ as follows:

$$
\mathbf{L}_{\vartheta}\left(L+L_{0}\right)=d_{H} \sigma \approx d_{H} \sigma-d_{H} \mathcal{J}_{v}
$$

Remark 7.2.1: In accordance with standard terminology, variational and exact symmetries generated by generalized vector fields (7.2.13) are called generalized symmetries because they depend on derivatives of variables. Accordingly, by variational symmetries and symmetries one means only those generated by vector fields $u$ on $Y$. We agree to call them classical symmetries.

Let $\vartheta$ be a classical variational symmetry of a Lagrangian $L$, i.e., $\vartheta$ (7.2.12) is the jet prolongation of a vector field $u$ on $Y$. Then the relation

$$
\begin{equation*}
\mathbf{L}_{\vartheta} \mathcal{E}_{L}=\delta\left(\mathbf{L}_{\vartheta} L\right) \tag{7.2.25}
\end{equation*}
$$

holds. It follows that $\vartheta$ also is a symmetry of the Euler-Lagrange operator $\mathcal{E}_{L}$ of $L$, i.e., $\mathbf{L}_{\vartheta} \mathcal{E}_{L}=0$. However, the equality (7.2.25) fails to be true in the case of generalized symmetries.

Definition 7.2.10: Let $E \rightarrow X$ be a vector bundle and $E(X)$ the $C^{\infty}(X)$ module $E(X)$ of sections of $E \rightarrow X$. Let $\zeta$ be a linear differential operator on $E(X)$ taking values into the vector space $\mathcal{G}_{L}$ of variational symmetries of a Lagrangian $L$ (see Definition 8.2.1). Elements

$$
\begin{equation*}
u_{\xi}=\zeta(\xi) \tag{7.2.26}
\end{equation*}
$$

of $\operatorname{Im} \zeta$ are called the gauge symmetry of a Lagrangian $L$ parameterized by sections $\xi$ of $E \rightarrow X$. They are called the gauge parameters.

Remark 7.2.2: The differential operator $\zeta$ in Definition 7.2.10 takes its values into the vector space $\mathcal{G}_{L}$ as a subspace of the $C^{\infty}(X)$-module $\mathfrak{d} \mathcal{O}_{\infty}^{0}$, but it sends the $C^{\infty}(X)$-module $E(X)$ into the real vector space $\mathcal{G}_{L} \subset \mathfrak{d} \mathcal{O}_{\infty}^{0}$. The differential operator $\zeta$ is assumed to be at least of first order.

Equivalently, the gauge symmetry (7.2.26) is given by a section $\tilde{\zeta}$ of the fibre bundle

$$
\left(J^{r} Y \underset{Y}{\times} J^{m} E\right) \underset{Y}{\times} T Y \rightarrow J^{r} Y \underset{Y}{\times} J^{m} E
$$

(see Definition 2.3.2) such that

$$
u_{\xi}=\zeta(\xi)=\tilde{\zeta} \circ \xi
$$

for any section $\xi$ of $E \rightarrow X$. Hence, it is a generalized vector field $u_{\zeta}$ on the product $Y \times E$ represented by a section of the pull-back bundle

$$
J^{k}(Y \underset{X}{\times} E) \underset{Y}{\times} T(Y \underset{X}{\times} E) \rightarrow J^{k}(Y \underset{X}{\times} E), \quad k=\max (r, m),
$$

which lives in $T Y \subset T(Y \times E)$. This generalized vector field yields a contact derivation $J^{\infty} u_{\zeta}(7.2 .12)$ of the real ring $\mathcal{O}_{\infty}^{0}[Y \times E]$ which obeys the following condition.

Condition 7.2.11: Given a Lagrangian

$$
L \in \mathcal{O}_{\infty}^{0, n} E \subset \mathcal{O}_{\infty}^{0, n}[Y \times E],
$$

let us consider its Lie derivative

$$
\begin{equation*}
\left.\left.\mathbf{L}_{J^{\infty} u_{\zeta}} L=J^{\infty} u_{\zeta}\right\rfloor d L+d\left(J^{\infty} u_{\zeta}\right\rfloor L\right) \tag{7.2.27}
\end{equation*}
$$

where $d$ is the exterior differential of $\mathcal{O}_{\infty}^{0}[Y \times E]$. Then, for any section $\xi$ of $E \rightarrow X$, the pull-back $\xi^{*} \mathbf{L}_{\vartheta}$ is $d_{H^{-}}$-exact.

It follows from the first variational formula (7.2.19) for the Lie derivative (7.2.27) that Condition 7.2 .11 holds only if $u_{\zeta}$ projects onto a generalized vector field on $E$ and, in this case, iff the density $\left.\left(u_{\zeta}\right)_{V}\right\rfloor \mathcal{E}$ is $d_{H^{-}}$exact. Thus, we come to the following equivalent definition of gauge symmetries.

Definition 7.2.12: Let $E \rightarrow X$ be a vector bundle. A gauge symmetry of a Lagrangian $L$ parameterized by sections $\xi$ of $E \rightarrow X$ is defined as a contact derivation $\vartheta=J^{\infty} u$ of the real ring $\mathcal{O}_{\infty}^{0}[Y \times E]$ such that:
(i) it vanishes on the subring $\mathcal{O}_{\infty}^{0} E$,
(ii) the generalized vector field $u$ is linear in coordinates $\chi_{\Lambda}^{a}$ on $J^{\infty} E$, and it projects onto a generalized vector field on $E$, i.e., it takes the form

$$
\begin{equation*}
u=\left(\sum_{0 \leq|\Lambda| \leq m} u_{a}^{\lambda \Lambda}\left(x^{\mu}\right) \chi_{\Lambda}^{a}\right) \partial_{\lambda}+\left(\sum_{0 \leq|\Lambda| \leq m} u_{a}^{i \Lambda}\left(x^{\mu}, y_{\Sigma}^{j}\right) \chi_{\Lambda}^{a}\right) \partial_{i} \tag{7.2.28}
\end{equation*}
$$

(iii) the vertical part of $u$ (7.2.28) obeys the equality

$$
\begin{equation*}
\left.u_{V}\right\rfloor \mathcal{E}=d_{H} \sigma \tag{7.2.29}
\end{equation*}
$$

For the sake of convenience, we also call a generalized vector field (7.2.28) the gauge symmetry. By virtue of item (iii) of Definition 7.2.12, $u$ (7.2.28) is a gauge symmetry iff its vertical part is so.

Gauge symmetries possess the following particular properties.
(i) Let $E^{\prime} \rightarrow X$ be a vector bundle and $\zeta^{\prime}$ a linear $E(X)$-valued differential operator on the $C^{\infty}(X)$-module $E^{\prime}(X)$ of sections of $E^{\prime} \rightarrow X$. Then

$$
u_{\zeta^{\prime}\left(\xi^{\prime}\right)}=\left(\zeta \circ \zeta^{\prime}\right)\left(\xi^{\prime}\right)
$$

also is a gauge symmetry of $L$ parameterized by sections $\xi^{\prime}$ of $E^{\prime} \rightarrow X$. It factorizes through the gauge symmetries $u_{\phi}$ (7.2.26).
(ii) If a gauge symmetry is an exact Lagrangian symmetry, the corresponding conserved symmetry current $\mathcal{J}_{u}(7.2 .24)$ is reduced to a superpotential (see Theorem 7.5.4).
(iii) The direct second Noether theorem associates to a gauge symmetry of a Lagrangian $L$ the Noether identities of its Euler-Lagrange operator $\delta L$.

Theorem 7.2.13: Let $u(7.2 .28)$ be a gauge symmetry of a Lagrangian $L$, then its Euler-Lagrange operator $\delta L$ obeys the Noether identities

$$
\begin{align*}
\mathcal{E}_{a}= & \sum_{0 \leq|\Lambda|}(-1)^{|\Lambda|} d_{\Lambda}\left[\left(u_{a}^{i \Lambda}-y_{\lambda}^{i} u_{a}^{\lambda \Lambda}\right) \mathcal{E}_{i}\right]=  \tag{7.2.30}\\
& \sum_{0 \leq|\Lambda|} \eta\left(u_{a}^{i}-y_{\lambda}^{i} u_{a}^{\lambda}\right)^{\Lambda} d_{\Lambda} \mathcal{E}_{i}=0
\end{align*}
$$

(see Notation 7.5.2).
It follows from direct second Noether Theorem 7.2.13 that gauge symmetries of Lagrangian field theory characterize its degeneracy. A problem is that any Lagrangian possesses gauge symmetries and, therefore, one must separate them into the trivial and non-trivial ones. Moreover, gauge symmetries can be reducible, i.e., $\operatorname{Ker} \zeta \neq 0$. To solve these problems, we follow a different definition of gauge symmetries as those associated to non-trivial Noether identities by means of inverse second Noether Theorem 7.5.3.

### 7.3 Grassmann-graded Lagrangian theory

We start with the following definition of jets of odd variables. Let us consider a vector bundle $F \rightarrow X$ and the simple graded manifolds $\left(X, \mathcal{A}_{J^{r} F}\right)$ modelled over the vector bundles $J^{r} F \rightarrow X$. There is the direct system of the corresponding DBGA

$$
\mathcal{S}^{*}[F ; X] \longrightarrow \mathcal{S}^{*}\left[J^{1} F ; X\right] \longrightarrow \cdots \mathcal{S}^{*}\left[J^{r} F ; X\right] \longrightarrow \cdots
$$

of graded exterior forms on graded manifolds $\left(X, \mathcal{A}_{J^{r} F}\right)$. Its direct limit $\mathcal{S}_{\infty}^{*}[F ; X]$ is the Grassmann-graded counterpart of the DGA $\mathcal{P}_{\infty}^{*}$.

In order to describe Lagrangian theories both of even and odd variables, let us consider a composite bundle

$$
\begin{equation*}
F \rightarrow Y \rightarrow X \tag{7.3.1}
\end{equation*}
$$

where $F \rightarrow Y$ is a vector bundle provided with bundle coordinates $\left(x^{\lambda}, y^{i}, q^{a}\right)$. We call the simple graded manifold $\left(Y, \mathfrak{A}_{F}\right)$ modelled over $F \rightarrow Y$ the composite graded manifold. Let us associate to this graded manifold the following DBGA $\mathcal{S}_{\infty}^{*}[F ; Y]$.

It is readily observed that the jet manifold $J^{r} F$ of $F \rightarrow X$ is a vector bundle $J^{r} F \rightarrow J^{r} Y$ coordinated by $\left(x^{\lambda}, y_{\Lambda}^{i}, q_{\Lambda}^{a}\right), 0 \leq|\Lambda| \leq r$. Let $\left(J^{r} Y, \mathfrak{A}_{r}\right)$ be a simple graded manifold modelled over this vector bundle. Its local basis is $\left(x^{\lambda}, y_{\Lambda}^{i}, c_{\Lambda}^{a}\right), 0 \leq|\Lambda| \leq r$. Let

$$
\begin{equation*}
\mathcal{S}_{r}^{*}[F ; Y]=\mathcal{S}_{r}^{*}\left[J^{r} F ; J^{r} Y\right] \tag{7.3.2}
\end{equation*}
$$

denote the DBGA of graded exterior forms on the simple graded manifold $\left(J^{r} Y, \mathfrak{A}_{r}\right)$. In particular, there is a cochain monomorphism

$$
\begin{equation*}
\mathcal{O}_{r}^{*}=\mathcal{O}^{*}\left(J^{r} Y\right) \rightarrow \mathcal{S}_{r}^{*}[F ; Y] \tag{7.3.3}
\end{equation*}
$$

The surjection

$$
\pi_{r}^{r+1}: J^{r+1} Y \rightarrow J^{r} Y
$$

yields an epimorphism of graded manifolds

$$
\left(\pi_{r}^{r+1}, \hat{\pi}_{r}^{r+1}\right):\left(J^{r+1} Y, \mathfrak{A}_{r+1}\right) \rightarrow\left(J^{r} Y, \mathfrak{A}_{r}\right)
$$

including the sheaf monomorphism

$$
\widehat{\pi}_{r}^{r+1}: \pi_{r}^{r+1 *} \mathfrak{A}_{r} \rightarrow \mathfrak{A}_{r+1}
$$

where $\pi_{r}^{r+1 *} \mathfrak{A}_{r}$ is the pull-back onto $J^{r+1} Y$ of the continuous fibre bundle $\mathfrak{A}_{r} \rightarrow J^{r} Y$. This sheaf monomorphism induces the monomorphism of the canonical presheaves $\overline{\mathfrak{A}}_{r} \rightarrow \overline{\mathfrak{A}}_{r+1}$, which associates to each open subset
$U \subset J^{r+1} Y$ the ring of sections of $\mathfrak{A}_{r}$ over $\pi_{r}^{r+1}(U)$. Accordingly, there is a monomorphism of the structure rings

$$
\begin{equation*}
\pi_{r}^{r+1 *}: \mathcal{S}_{r}^{0}[F ; Y] \rightarrow \mathcal{S}_{r+1}^{0}[F ; Y] \tag{7.3.4}
\end{equation*}
$$

of graded functions on graded manifolds $\left(J^{r} Y, \mathfrak{A}_{r}\right)$ and $\left(J^{r+1} Y, \mathfrak{A}_{r+1}\right)$. By virtue of Lemma 6.4.1, the differential calculus $\mathcal{S}_{r}^{*}[F ; Y]$ and $\mathcal{S}_{r+1}^{*}[F ; Y]$ are minimal. Therefore, the monomorphism (7.3.4) yields that of the DBGA

$$
\begin{equation*}
\pi_{r}^{r+1 *}: \mathcal{S}_{r}^{*}[F ; Y] \rightarrow \mathcal{S}_{r+1}^{*}[F ; Y] . \tag{7.3.5}
\end{equation*}
$$

As a consequence, we have the direct system of DBGAs

$$
\begin{align*}
& \mathcal{S}^{*}[F ; Y] \xrightarrow{\pi^{*}} \mathcal{S}_{1}^{*}[F ; Y] \longrightarrow \cdots \mathcal{S}_{r-1}^{*}[F ; Y] \xrightarrow{\pi_{r-1}^{r *}}  \tag{7.3.6}\\
& \mathcal{S}_{r}^{*}[F ; Y] \longrightarrow \cdots
\end{align*}
$$

The DBGA $\mathcal{S}_{\infty}^{*}[F ; Y]$ that we associate to the composite graded manifold $\left(Y, \mathfrak{A}_{F}\right)$ is defined as the direct limit

$$
\begin{equation*}
\mathcal{S}_{\infty}^{*}[F ; Y]=\overrightarrow{\lim } \mathcal{S}_{r}^{*}[F ; Y] \tag{7.3.7}
\end{equation*}
$$

of the direct system (7.3.6). It consists of all graded exterior forms $\phi \in$ $\mathcal{S}^{*}\left[F_{r} ; J^{r} Y\right]$ on graded manifolds $\left(J^{r} Y, \mathfrak{A}_{r}\right)$ modulo the monomorphisms (7.3.5). Its elements obey the relations (6.2.9) - (6.2.10).

Cochain monomorphisms $\mathcal{O}_{r}^{*} \rightarrow \mathcal{S}_{r}^{*}[F ; Y]$ (7.3.3) provide a monomorphism of the direct system (2.4.5) to the direct system (7.3.6) and, consequently, the monomorphism

$$
\begin{equation*}
\mathcal{O}_{\infty}^{*} \rightarrow \mathcal{S}_{\infty}^{*}[F ; Y] \tag{7.3.8}
\end{equation*}
$$

of their direct limits. In particular, $\mathcal{S}_{\infty}^{*}[F ; Y]$ is an $\mathcal{O}_{\infty}^{0}$-algebra. Accordingly, the body epimorphisms $\mathcal{S}_{r}^{*}[F ; Y] \rightarrow \mathcal{O}_{r}^{*}$ yield the epimorphism of $\mathcal{O}_{\infty}^{0}$-algebras

$$
\begin{equation*}
\mathcal{S}_{\infty}^{*}[F ; Y] \rightarrow \mathcal{O}_{\infty}^{*} \tag{7.3.9}
\end{equation*}
$$

It is readily observed that the morphisms (7.3.8) and (7.3.9) are cochain morphisms between the de Rham complex (2.4.7) of the DGA $\mathcal{O}_{\infty}^{*}$ (2.4.6) and the de Rham complex

$$
\begin{equation*}
0 \rightarrow \mathbb{R} \rightarrow \mathcal{S}_{\infty}^{0}[F ; Y] \xrightarrow{d} \mathcal{S}_{\infty}^{1}[F ; Y] \cdots \xrightarrow{d} \mathcal{S}_{\infty}^{k}[F ; Y] \longrightarrow \cdots \tag{7.3.10}
\end{equation*}
$$

of the DBGA $\mathcal{S}_{\infty}^{0}[F ; Y]$. Moreover, the corresponding homomorphisms of cohomology groups of these complexes are isomorphisms as follows.

Theorem 7.3.1: There is an isomorphism

$$
\begin{equation*}
H^{*}\left(\mathcal{S}_{\infty}^{*}[F ; Y]\right)=H_{D R}^{*}(Y) \tag{7.3.11}
\end{equation*}
$$

of the cohomology $H^{*}\left(\mathcal{S}_{\infty}^{*}[F ; Y]\right)$ of the de Rham complex (7.3.10) to the de Rham cohomology $H_{D R}^{*}(Y)$ of $Y$.

Corollary 7.3.2: Any closed graded form $\phi \in \mathcal{S}_{\infty}^{*}[F ; Y]$ is decomposed into the sum $\phi=\sigma+d \xi$ where $\sigma$ is a closed exterior form on $Y$.

Similarly to the DGA $\mathcal{O}_{\infty}^{*}(2.4 .6)$, one thinks of elements of $\mathcal{S}_{\infty}^{*}[F ; Y]$ as being graded differential forms on the infinite order jet manifold $J^{\infty} Y$. We can restrict $\mathcal{S}_{\infty}^{*}[F ; Y]$ to the coordinate chart (2.4.3) of $J^{\infty} Y$ and say that $\mathcal{S}_{\infty}^{*}[F ; Y]$ as an $\mathcal{O}_{\infty}^{0}$-algebra is locally generated by the elements

$$
\left(c_{\Lambda}^{a}, d x^{\lambda}, \theta_{\Lambda}^{a}=d c_{\Lambda}^{a}-c_{\lambda+\Lambda}^{a} d x^{\lambda}, \theta_{\Lambda}^{i}=d y_{\Lambda}^{i}-y_{\lambda+\Lambda}^{i} d x^{\lambda}\right), \quad 0 \leq|\Lambda|
$$

where $c_{\Lambda}^{a}, \theta_{\Lambda}^{a}$ are odd and $d x^{\lambda}, \theta_{\Lambda}^{i}$ are even. We agree to call $\left(y^{i}, c^{a}\right)$ the local generating basis for $\mathcal{S}_{\infty}^{*}[F ; Y]$. Let the collective symbol $s^{A}$ stand for its elements. Accordingly, the notation $s_{\Lambda}^{A}$ and

$$
\theta_{\Lambda}^{A}=d s_{\Lambda}^{A}-s_{\lambda+\Lambda}^{A} d x^{\lambda}
$$

is introduced. For the sake of simplicity, we further denote $[A]=\left[s^{A}\right]$.
The DBGA $\mathcal{S}_{\infty}^{*}[F ; Y]$ is split into $\mathcal{S}_{\infty}^{0}[F ; Y]$-modules $\mathcal{S}_{\infty}^{k, r}[F ; Y]$ of $k$ contact and r-horizontal graded forms together with the corresponding projections

$$
h_{k}: \mathcal{S}_{\infty}^{*}[F ; Y] \rightarrow \mathcal{S}_{\infty}^{k, *}[F ; Y], \quad h^{m}: \mathcal{S}_{\infty}^{*}[F ; Y] \rightarrow \mathcal{S}_{\infty}^{*, m}[F ; Y]
$$

Accordingly, the graded exterior differential $d$ on $\mathcal{S}_{\infty}^{*}[F ; Y]$ falls into the sum $d=d_{V}+d_{H}$ of the vertical graded differential

$$
d_{V} \circ h^{m}=h^{m} \circ d \circ h^{m}, \quad d_{V}(\phi)=\theta_{\Lambda}^{A} \wedge \partial_{A}^{\Lambda} \phi, \quad \phi \in \mathcal{S}_{\infty}^{*}[F ; Y],
$$

and the total graded differential

$$
d_{H} \circ h_{k}=h_{k} \circ d \circ h_{k}, \quad d_{H} \circ h_{0}=h_{0} \circ d, \quad d_{H}(\phi)=d x^{\lambda} \wedge d_{\lambda}(\phi),
$$

where

$$
d_{\lambda}=\partial_{\lambda}+\sum_{0 \leq|\Lambda|} s_{\lambda+\Lambda}^{A} \partial_{A}^{\Lambda}
$$

are the graded total derivatives. These differentials obey the nilpotent relations (2.4.11).

Similarly to the DGA $\mathcal{O}_{\infty}^{*}$, the DBGA $\mathcal{S}_{\infty}^{*}[F ; Y]$ is provided with the graded projection endomorphism

$$
\begin{aligned}
& \varrho=\sum_{k>0} \frac{1}{k} \bar{\varrho} \circ h_{k} \circ h^{n}: \mathcal{S}_{\infty}^{*>0, n}[F ; Y] \rightarrow \mathcal{S}_{\infty}^{*>0, n}[F ; Y], \\
& \left.\bar{\varrho}(\phi)=\sum_{0 \leq \backslash \Lambda \mid}(-1)^{|\Lambda|} \theta^{A} \wedge\left[d_{\Lambda}\left(\partial_{A}^{\Lambda}\right\rfloor \phi\right)\right], \quad \phi \in \mathcal{S}_{\infty}^{>0, n}[F ; Y],
\end{aligned}
$$

such that $\varrho \circ d_{H}=0$, and with the nilpotent graded variational operator

$$
\begin{equation*}
\delta=\varrho \circ d \mathcal{S}_{\infty}^{*, n}[F ; Y] \rightarrow \mathcal{S}_{\infty}^{*+1, n}[F ; Y] . \tag{7.3.12}
\end{equation*}
$$

With these operators the DBGA $\mathcal{S}_{\infty}^{*}[F ; Y]$ is split into the Grassmanngraded variational bicomplex. We restrict our consideration to its short variational subcomplex

$$
\begin{align*}
0 \rightarrow & \mathbb{R} \rightarrow \mathcal{S}_{\infty}^{0}[F ; Y] \xrightarrow{d_{H}} \mathcal{S}_{\infty}^{0,1}[F ; Y] \cdots \xrightarrow{d_{H}}  \tag{7.3.13}\\
& \mathcal{S}_{\infty}^{0, n}[F ; Y] \xrightarrow{\delta} \mathbf{E}_{1}, \quad \mathbf{E}_{1}=\varrho\left(\mathcal{S}_{\infty}^{1, n}[F ; Y]\right),
\end{align*}
$$

and the subcomplex of one-contact graded forms

$$
\begin{align*}
0 \rightarrow & \mathcal{S}_{\infty}^{1,0}[F ; Y] \xrightarrow{d_{H}} \mathcal{S}_{\infty}^{1,1}[F ; Y] \cdots \xrightarrow{d_{H}} \mathcal{S}_{\infty}^{1, n}[F ; Y]  \tag{7.3.14}\\
& \xrightarrow{\varrho} \mathbf{E}_{1} \rightarrow 0 .
\end{align*}
$$

Theorem 7.3.3: Cohomology of the complex (7.3.13) equals the de Rham cohomology $H_{D R}^{*}(Y)$ of $Y$.

Theorem 7.3.4: The complex (7.3.14) is exact.
Decomposed into the variational bicomplex, the DBGA $\mathcal{S}_{\infty}^{*}[F ; Y]$ describes Grassmann-graded Lagrangian theory on the composite graded manifold $\left(Y, \mathfrak{A}_{F}\right)$. Its graded Lagrangian is defined as an element

$$
\begin{equation*}
L=\mathcal{L} \omega \in \mathcal{S}_{\infty}^{0, n}[F ; Y] \tag{7.3.15}
\end{equation*}
$$

of the graded variational complex (7.3.13), while the graded exterior form

$$
\begin{equation*}
\delta L=\theta^{A} \wedge \mathcal{E}_{A} \omega=\sum_{0 \leq|\Lambda|}(-1)^{|\Lambda|} \theta^{A} \wedge d_{\Lambda}\left(\partial_{A}^{\Lambda} L\right) \omega \in \mathbf{E}_{1} \tag{7.3.16}
\end{equation*}
$$

is said to be its graded Euler-Lagrange operator. We agree to call a pair $\left(\mathcal{S}_{\infty}^{0, n}[F ; Y], L\right)$ the Grassmann-graded Lagrangian system.

The following is a corollary of Theorem 7.3.3.
Theorem 7.3.5: Every $d_{H}$-closed graded form $\phi \in \mathcal{S}_{\infty}^{0, m<n}[F ; Y]$ falls into the sum

$$
\begin{equation*}
\phi=h_{0} \sigma+d_{H} \xi, \quad \xi \in \mathcal{S}_{\infty}^{0, m-1}[F ; Y] \tag{7.3.17}
\end{equation*}
$$

where $\sigma$ is a closed $m$-form on $Y$. Any $\delta$-closed (i.e., variationally trivial) Grassmann-graded Lagrangian $L \in \mathcal{S}_{\infty}^{0, n}[F ; Y]$ is the sum

$$
\begin{equation*}
L=h_{0} \sigma+d_{H} \xi, \quad \xi \in \mathcal{S}_{\infty}^{0, n-1}[F ; Y] \tag{7.3.18}
\end{equation*}
$$

where $\sigma$ is a closed $n$-form on $Y$.

Corollary 7.3.6: Any variationally trivial odd Lagrangian is $d_{H}$-exact.

The exactness of the complex (7.3.14) at the term $\mathcal{S}_{\infty}^{1, n}[F ; Y]$ results in the following.

Theorem 7.3.7: Given a graded Lagrangian $L$, there is the decomposition

$$
\begin{align*}
& d L=\delta L-d_{H} \Xi_{L}, \quad \Xi \in \mathcal{S}_{\infty}^{n-1}[F ; Y]  \tag{7.3.19}\\
& \Xi_{L}=L+\sum_{s=0} \theta_{\nu_{s} \ldots \nu_{1}}^{A} \wedge F_{A}^{\lambda \nu_{s} \ldots \nu_{1}} \omega_{\lambda},  \tag{7.3.20}\\
& F_{A}^{\nu_{k} \ldots \nu_{1}}=\partial_{A}^{\nu_{k} \ldots \nu_{1}} \mathcal{L}-d_{\lambda} F_{A}^{\lambda \nu_{k} \ldots \nu_{1}}+\sigma_{A}^{\nu_{k} \ldots \nu_{1}}, \quad k=1,2, \ldots,
\end{align*}
$$

where local graded functions $\sigma$ obey the relations

$$
\sigma_{A}^{\nu}=0, \quad \sigma_{A}^{\left(\nu_{k} \nu_{k-1}\right) \ldots \nu_{1}}=0
$$

The form $\Xi_{L}(7.3 .20)$ provides a global Lepage equivalent of a graded Lagrangian $L$.

Given a Grassmann-graded Lagrangian system $\left(\mathcal{S}_{\infty}^{*}[F ; Y], L\right)$, by its infinitesimal transformations are meant contact graded derivations of the real graded commutative ring $\mathcal{S}_{\infty}^{0}[F ; Y]$. They constitute a $\mathcal{S}_{\infty}^{0}[F ; Y]$ module $\mathfrak{o} \mathcal{S}_{\infty}^{0}[F ; Y]$ which is a real Lie superalgebra with the Lie superbracket (6.2.4).

Theorem 7.3.8: The derivation module $\mathcal{J}_{\infty}^{0}[F ; Y]$ is isomorphic to the $\mathcal{S}_{\infty}^{0}[F ; Y]$-dual $\left(\mathcal{S}_{\infty}^{1}[F ; Y]\right)^{*}$ of the module of graded one-forms $\mathcal{S}_{\infty}^{1}[F ; Y]$. It follows that the DBGA $\mathcal{S}_{\infty}^{*}[F ; Y]$ is minimal differential calculus over the real graded commutative ring $\mathcal{S}_{\infty}^{0}[F ; Y]$.

Let $\vartheta\rfloor \phi, \vartheta \in \mathfrak{o} \mathcal{S}_{\infty}^{0}[F ; Y], \phi \in \mathcal{S}_{\infty}^{1}[F ; Y]$, denote the corresponding interior product. Extended to the DBGA $\mathcal{S}_{\infty}^{*}[F ; Y]$, it obeys the rule

$$
\left.\vartheta\rfloor(\phi \wedge \sigma)=(\vartheta\rfloor \phi) \wedge \sigma+(-1)^{|\phi|+[\phi][\vartheta]} \phi \wedge(\vartheta\rfloor \sigma\right), \quad \phi, \sigma \in \mathcal{S}_{\infty}^{*}[F ; Y] .
$$

Restricted to a coordinate chart (2.4.3) of $J^{\infty} Y$, the algebra $\mathcal{S}_{\infty}^{*}[F ; Y]$ is a free $\mathcal{S}_{\infty}^{0}[F ; Y]$-module generated by one-forms $d x^{\lambda}, \theta_{\Lambda}^{A}$. Due to the isomorphism stated in Theorem 7.3.8, any graded derivation $\vartheta \in$ $\mathfrak{d} \mathcal{S}_{\infty}^{0}[F ; Y]$ takes the local form

$$
\begin{align*}
& \vartheta=\vartheta^{\lambda} \partial_{\lambda}+\vartheta^{A} \partial_{A}+\sum_{0<|\Lambda|} \vartheta_{\Lambda}^{A} \partial_{A}^{\Lambda},  \tag{7.3.21}\\
& \left.\partial_{A}^{\Lambda}\right\rfloor d y_{\Sigma}^{B}=\delta_{A}^{B} \delta_{\Sigma}^{\Lambda} . \tag{7.3.22}
\end{align*}
$$

Every graded derivation $\vartheta(7.3 .21)$ yields the graded Lie derivative

$$
\begin{aligned}
& \left.\left.\mathbf{L}_{\vartheta} \phi=\vartheta\right\rfloor d \phi+d(\vartheta\rfloor \phi\right), \quad \phi \in \mathcal{S}_{\infty}^{*}[F ; Y] \\
& \mathbf{L}_{\vartheta}(\phi \wedge \sigma)=\mathbf{L}_{\vartheta}(\phi) \wedge \sigma+(-1)^{[\vartheta][\phi]} \phi \wedge \mathbf{L}_{\vartheta}(\sigma)
\end{aligned}
$$

of the DBGA $\mathcal{S}_{\infty}^{*}[F ; Y]$. A graded derivation $\vartheta$ (7.3.21) is called contact if the Lie derivative $\mathbf{L}_{\vartheta}$ preserves the ideal of contact graded forms.

Lemma 7.3.9: With respect to the local generating basis $\left(s^{A}\right)$ for the DBGA $\mathcal{S}_{\infty}^{*}[F ; Y]$, any its contact graded derivation takes the form

$$
\begin{equation*}
\vartheta=v_{H}+v_{V}=v^{\lambda} d_{\lambda}+\left[v^{A} \partial_{A}+\sum_{|\Lambda|>0} d_{\Lambda}\left(v^{A}-s_{\mu}^{A} v^{\mu}\right) \partial_{A}^{\Lambda}\right] \tag{7.3.23}
\end{equation*}
$$

where $v_{H}$ and $v_{V}$ denotes the horizontal and vertical parts of $\vartheta$.
A glance at the expression (7.3.23) shows that a contact graded derivation $\vartheta$ is an infinite order jet prolongation of its restriction

$$
\begin{equation*}
v=v^{\lambda} \partial_{\lambda}+v^{A} \partial_{A} \tag{7.3.24}
\end{equation*}
$$

to the graded commutative ring $S^{0}[F ; Y]$. We call $v(7.3 .24)$ the generalized graded vector field. It is readily justified the following (see Lemma 7.2.16).

Lemma 7.3.10: Any vertical contact graded derivation

$$
\begin{equation*}
\vartheta=v^{A} \partial_{A}+\sum_{|\Lambda|>0} d_{\Lambda} v^{A} \partial_{A}^{\Lambda} \tag{7.3.25}
\end{equation*}
$$

satisfies the relations

$$
\begin{align*}
& \left.\vartheta\rfloor d_{H} \phi=-d_{H}(\vartheta\rfloor \phi\right),  \tag{7.3.26}\\
& \mathbf{L}_{\vartheta}\left(d_{H} \phi\right)=d_{H}\left(\mathbf{L}_{\vartheta} \phi\right) \tag{7.3.27}
\end{align*}
$$

for all $\phi \in \mathcal{S}_{\infty}^{*}[F ; Y]$.
Then the forthcoming assertions are the straightforward generalizations of Theorem 7.2.7, Lemma 7.2.8 and Theorem 7.2.9.

A corollary of the decomposition (7.3.19) is the first variational formula for a graded Lagrangian.

Theorem 7.3.11: The Lie derivative of a graded Lagrangian along any contact graded derivation (7.3.23) obeys the first variational formula

$$
\begin{equation*}
\left.\left.\left.\mathbf{L}_{\vartheta} L=v_{V}\right\rfloor \delta L+d_{H}\left(h_{0}(\vartheta\rfloor \Xi_{L}\right)\right)+d_{V}\left(v_{H}\right\rfloor \omega\right) \mathcal{L}, \tag{7.3.28}
\end{equation*}
$$

where $\Xi_{L}$ is the Lepage equivalent (7.3.20) of $L$.
A contact graded derivation $\vartheta$ (7.3.23) is called a variational symmetry (strictly speaking, a variational supersymmetry) of a graded Lagrangian $L$ if the Lie derivative $\mathbf{L}_{\vartheta} L$ is $d_{H}$-exact, i.e.,

$$
\begin{equation*}
\mathbf{L}_{\vartheta} L=d_{H} \sigma . \tag{7.3.29}
\end{equation*}
$$

Lemma 7.3.12: A glance at the expression (7.3.28) shows the following.
(i) A contact graded derivation $\vartheta$ is a variational symmetry only if it is projected onto $X$.
(ii) Any projectable contact graded derivation is a variational symmetry of a variationally trivial graded Lagrangian. It follows that, if $\vartheta$ is a variational symmetry of a graded Lagrangian $L$, it also is a variational symmetry of a Lagrangian $L+L_{0}$, where $L_{0}$ is a variationally trivial graded Lagrangian.
(iii) A contact graded derivations $\vartheta$ is a variational symmetry iff its vertical part $v_{V}$ (7.3.23) is well.
(iv) It is a variational symmetry iff the graded density $\left.v_{V}\right\rfloor \delta L$ is $d_{H^{-}}$ exact.

Variational symmetries of a graded Lagrangian $L$ constitute a real vector subspace $\mathcal{G}_{L}$ of the graded derivation module $\boldsymbol{\mathcal { S }} \mathcal{S}_{\infty}^{0}[F ; Y]$. By virtue of item (ii) of Lemma 7.3.12, the Lie superbracket

$$
\mathbf{L}_{\left[\vartheta, \vartheta^{\prime}\right]}=\left[\mathbf{L}_{\vartheta}, \mathbf{L}_{\vartheta^{\prime}}\right]
$$

of variational symmetries is a variational symmetry and, therefore, their vector space $\mathcal{G}_{L}$ is a real Lie superalgebra.

A corollary of the first variational formula (7.3.28) is the first Noether theorem for graded Lagrangians.

Theorem 7.3.13: If a contact graded derivation $\vartheta(7.3 .23)$ is a variational symmetry (7.3.29) of a graded Lagrangian $L$, the first variational formula (7.3.28) restricted to Ker $\delta L$ leads to the weak conservation law

$$
\begin{equation*}
\left.0 \approx d_{H}\left(h_{0}(\vartheta\rfloor \Xi_{L}\right)-\sigma\right) \tag{7.3.30}
\end{equation*}
$$

A vertical contact graded derivation $\vartheta(7.3 .25)$ is said to be nilpotent if

$$
\begin{gather*}
\mathbf{L}_{\vartheta}\left(\mathbf{L}_{\vartheta} \phi\right)=\sum_{0 \leq|\Sigma|, 0 \leq|\Lambda|}\left(v_{\Sigma}^{B} \partial_{B}^{\Sigma}\left(v_{\Lambda}^{A}\right) \partial_{A}^{\Lambda}+\right.  \tag{7.3.31}\\
\left.(-1)^{\left[s^{B}\right]\left[v^{A}\right]} v_{\Sigma}^{B} v_{\Lambda}^{A} \partial_{B}^{\Sigma} \partial_{A}^{\Lambda}\right) \phi=0
\end{gather*}
$$

for any horizontal graded form $\phi \in S_{\infty}^{0, *}$.
Lemma 7.3.14: A vertical contact graded derivation (7.3.25) is nilpotent only if it is odd and iff the equality

$$
\mathbf{L}_{\vartheta}\left(v^{A}\right)=\sum_{0 \leq|\Sigma|} v_{\Sigma}^{B} \partial_{B}^{\Sigma}\left(v^{A}\right)=0
$$

holds for all $v^{A}$.
For the sake of brevity, the common symbol $v$ further stands for a generalized graded vector field $v$, the contact graded derivation $\vartheta$ determined by $v$, and the Lie derivative $\mathbf{L}_{\vartheta}$. We agree to call all these operators, simply, a graded derivation of a field system algebra.

Remark 7.3.1: For the sake of convenience, right derivations

$$
\begin{equation*}
\overleftarrow{v}=\overleftarrow{\partial}_{A} v^{A} \tag{7.3.32}
\end{equation*}
$$

also are considered. They act on graded functions and differential forms $\phi$ on the right by the rules

$$
\begin{aligned}
& \overleftarrow{v}(\phi)=d \phi\lfloor\overleftarrow{v}+d(\phi\lfloor\overleftarrow{v}) \\
& \overleftarrow{v}\left(\phi \wedge \phi^{\prime}\right)=(-1)^{\left[\phi^{\prime}\right]} \overleftarrow{v}(\phi) \wedge \phi^{\prime}+\phi \wedge \overleftarrow{v}\left(\phi^{\prime}\right) \\
& \theta_{\Lambda A}\left\lfloor\overleftarrow{\partial}{ }^{\Sigma B}=\delta_{B}^{A} \delta_{\Lambda}^{\Sigma}\right.
\end{aligned}
$$

One associates to any graded right derivation $\overleftarrow{v}(7.3 .32)$ the left one

$$
\begin{align*}
& v^{l}=(-1)^{[v][A]} v^{A} \partial_{A}  \tag{7.3.33}\\
& v^{l}(f)=(-1)^{[v][f]} v(f), \quad f \in \mathcal{S}_{\infty}^{0}[F ; Y]
\end{align*}
$$

### 7.4 Noether identities

The degeneracy of Lagrangian theory is characterized by a set of nontrivial reducible Noether identities. Any Euler-Lagrange operator satisfies Noether identities (henceforth NI) which therefore must be separated into the trivial and non-trivial ones. These NI can obey first-stage NI, which in turn are subject to the second-stage ones, and so on. Thus, there is a hierarchy of higher-stage NI which also are separated into the trivial and non-trivial ones. If certain conditions hold, one can associate to a Grassmann-graded Lagrangian system the exact Koszul-Tate complex possessing the boundary operator whose nilpotentness is equivalent to all non-trivial NI and higher-stage NI. The inverse second Noether theorem formulated in homology terms associates to this Koszul-Tate complex the cochain sequence of ghosts with the ascent operator, called the gauge operator, whose components are non-trivial gauge and higherstage gauge symmetries of Lagrangian theory.

Let $\left(\mathcal{S}_{\infty}^{*}[F ; Y], L\right)$ be a Grassmann-graded Lagrangian system. Without a lose of generality, let a Lagrangian $L$ be even. Its Euler-Lagrange operator $\delta L$ (7.3.16) is assumed to be at least of order 1 in order to guarantee that transition functions of $Y$ do not vanish on-shell. This Euler-Lagrange operator $\delta L \in \mathbf{E}_{1}$ takes its values into the graded vector bundle

$$
\begin{equation*}
\overline{V F}=V^{*} F \underset{F}{\otimes} \wedge T^{*} X \rightarrow F \tag{7.4.1}
\end{equation*}
$$

where $V^{*} F$ is the vertical cotangent bundle of $F \rightarrow X$. It however is not a vector bundle over $Y$. Therefore, we restrict our consideration to the case of a pull-back composite bundle $F$ (7.3.1), that is,

$$
\begin{equation*}
F=\underset{X}{\times} \underset{X}{\times} F^{1} Y \rightarrow X \tag{7.4.2}
\end{equation*}
$$

where $F^{1} \rightarrow X$ is a vector bundle. Let us introduce the following notation.

Notation 7.4.1: Given the vertical tangent bundle $V E$ of a fibre bundle $E \rightarrow X$, by its density-dual bundle is meant the fibre bundle

$$
\begin{equation*}
\overline{V E}=V^{*} E \otimes_{E}^{\otimes} \wedge T^{*} X \tag{7.4.3}
\end{equation*}
$$

If $E \rightarrow X$ is a vector bundle, we have

$$
\begin{equation*}
\overline{V E}=\bar{E} \times \underset{X}{\times} E, \quad \bar{E}=E^{*} \otimes_{X}^{\otimes} \wedge T^{*} X \tag{7.4.4}
\end{equation*}
$$

where $\bar{E}$ is called the density-dual of $E$. Let

$$
E=E^{0} \underset{X}{\oplus} E^{1}
$$

be a graded vector bundle over $X$. Its graded density-dual is defined to be

$$
\bar{E}=\bar{E}^{1} \oplus_{X} \bar{E}^{0}
$$

In these terms, we treat the composite bundle $F$ (7.3.1) as a graded vector bundle over $Y$ possessing only odd part. The density-dual $\overline{V F}$ (7.4.3) of the vertical tangent bundle $V F$ of $F \rightarrow X$ is $\overline{V F}$ (7.4.1). If $F$ (7.3.1) is the pull-back bundle (7.4.2), then

$$
\begin{equation*}
\overline{V F}=\left(\left(\bar{F}^{1} \underset{Y}{\oplus} V^{*} Y\right) \underset{Y}{\otimes} \wedge{ }^{n} T^{*} X\right) \underset{Y}{\oplus} F^{1} \tag{7.4.5}
\end{equation*}
$$

is a graded vector bundle over $Y$. Given a graded vector bundle

$$
E=E^{0} \underset{Y}{\oplus} E^{1} \rightarrow Y
$$

we consider the composite bundle $E \rightarrow E^{0} \rightarrow X$ and the DBGA (7.3.7):

$$
\begin{equation*}
\mathcal{P}_{\infty}^{*}[E ; Y]=\mathcal{S}_{\infty}^{*}\left[E ; E^{0}\right] \tag{7.4.6}
\end{equation*}
$$

Let us consider the density-dual $\overline{V F}$ (7.4.5) of the vertical tangent bundle $V F \rightarrow F$, and let us enlarge the DBGA $\mathcal{S}_{\infty}^{*}[F ; Y]$ to the DBGA $\mathcal{P}_{\infty}^{*}[\overline{V F} ; Y]$ (7.4.6) with the local generating basis $\left(s^{A}, \bar{s}_{A}\right), \quad\left[\bar{s}_{A}\right]=$ $([A]+1) \bmod 2$. Following the physical terminology, we agree to call its elements $\bar{s}_{A}$ the antifields of antifield number $\operatorname{Ant}\left[\bar{s}_{A}\right]=1$. The DBGA $\mathcal{P}_{\infty}^{*}[\overline{V F} ; Y]$ is endowed with the nilpotent right graded derivation $\bar{\delta}=\overleftarrow{\partial}^{A} \mathcal{E}_{A}$, where $\mathcal{E}_{A}$ are the variational derivatives (7.3.16). Then we have the chain complex

$$
\begin{equation*}
0 \leftarrow \operatorname{Im} \bar{\delta} \stackrel{\bar{\delta}}{\leftrightarrows} \mathcal{P}_{\infty}^{0, n}[\overline{V F} ; Y]_{1} \stackrel{\bar{\delta}}{\leftrightarrows} \mathcal{P}_{\infty}^{0, n}[\overline{V F} ; Y]_{2} \tag{7.4.7}
\end{equation*}
$$

of graded densities of antifield number $\leq 2$. Its one-boundaries $\bar{\delta} \Phi$, $\Phi \in \mathcal{P}_{\infty}^{0, n}[\overline{V F} ; Y]_{2}$, by very definition, vanish on-shell.

Lemma 7.4.2: One can associate to any Grassmann-graded Lagrangian system $\left(\mathcal{S}_{\infty}^{*}[F ; Y], L\right)$ the chain complex (7.4.7).

Any one-cycle

$$
\begin{equation*}
\Phi=\sum_{0 \leq|\Lambda|} \Phi^{A, \Lambda_{\bar{s}_{\Lambda A}} \omega \in \mathcal{P}_{\infty}^{0, n}[\overline{V F} ; Y]_{1} .} \tag{7.4.8}
\end{equation*}
$$

of the complex (7.4.7) is a differential operator on the bundle $\overline{V F}$ such that it is linear on fibres of $\overline{V F} \rightarrow F$ and its kernel contains the graded Euler-Lagrange operator $\delta L$ (7.3.16), i.e.,

$$
\begin{equation*}
\bar{\delta} \Phi=0, \quad \sum_{0 \leq|\Lambda|} \Phi^{A, \Lambda} d_{\Lambda} \mathcal{E}_{A} \omega=0 \tag{7.4.9}
\end{equation*}
$$

Thus, the one-cycles (7.4.8) define the NI (7.4.9) of the Euler-Lagrange operator $\delta L$, which we call Noether identities (NI) of the Grassmanngraded Lagrangian system $\left(\mathcal{S}_{\infty}^{*}[F ; Y], L\right)$.

In particular, one-chains $\Phi(7.4 .8)$ are necessarily NI if they are boundaries. Accordingly, non-trivial NI modulo the trivial ones are associated to elements of the first homology $H_{1}(\bar{\delta})$ of the complex (7.4.7). A Lagrangian $L$ is called degenerate if there are non-trivial NI.

Non-trivial NI obey first-stage NI. To describe them, let us assume that the module $H_{1}(\bar{\delta})$ is finitely generated. Namely, there exists a graded projective $C^{\infty}(X)$-module $\mathcal{C}_{(0)} \subset H_{1}(\bar{\delta})$ of finite rank with a local basis $\left\{\Delta_{r} \omega\right\}$ :

$$
\begin{equation*}
\Delta_{r} \omega=\sum_{0 \leq|\Lambda|} \Delta_{r}^{A, \Lambda} \bar{s}_{\Lambda A} \omega, \quad \Delta_{r}^{A, \Lambda} \in \mathcal{S}_{\infty}^{0}[F ; Y] \tag{7.4.10}
\end{equation*}
$$

such that any element $\Phi \in H_{1}(\bar{\delta})$ factorizes as

$$
\begin{equation*}
\Phi=\sum_{0 \leq|\Xi|} \Phi^{r, \Xi} d_{\Xi} \Delta_{r} \omega, \quad \Phi^{r, \Xi} \in \mathcal{S}_{\infty}^{0}[F ; Y] \tag{7.4.11}
\end{equation*}
$$

through elements (7.4.10) of $\mathcal{C}_{(0)}$. Thus, all non-trivial NI (7.4.9) result from the NI

$$
\begin{equation*}
\bar{\delta} \Delta_{r}=\sum_{0 \leq|\Lambda|} \Delta_{r}^{A, \Lambda} d_{\Lambda} \mathcal{E}_{A}=0 \tag{7.4.12}
\end{equation*}
$$

called the complete NI. Clearly, the factorization (7.4.11) is independent of specification of a local basis $\left\{\Delta_{r} \omega\right\}$.

A Lagrangian system whose non-trivial NI are finitely generated is called finitely degenerate. Hereafter, degenerate Lagrangian systems only of this type are considered.

By virtue of Serre-Swan Theorem 6.3.2, the graded module $\mathcal{C}_{(0)}$ is isomorphic to a module of sections of the density-dual $\bar{E}_{0}$ of some graded vector bundle $E_{0} \rightarrow X$. Let us enlarge $\mathcal{P}_{\infty}^{*}[\overline{V F} ; Y]$ to the DBGA

$$
\begin{equation*}
\overline{\mathcal{P}}_{\infty}^{*}\{0\}=\mathcal{P}_{\infty}^{*}\left[\overline{V F} \underset{Y}{\oplus} \bar{E}_{0} ; Y\right] \tag{7.4.13}
\end{equation*}
$$

possessing the local generating basis $\left(s^{A}, \bar{s}_{A}, \bar{c}_{r}\right)$ where $\bar{c}_{r}$ are Noether antifields of Grassmann parity $\left[\bar{c}_{r}\right]=\left(\left[\Delta_{r}\right]+1\right) \bmod 2$ and antifield number $\operatorname{Ant}\left[\bar{c}_{r}\right]=2$. The DBGA (7.4.13) is provided with the odd right graded derivation $\delta_{0}=\bar{\delta}+\overleftarrow{\partial}^{r} \Delta_{r}$ which is nilpotent iff the complete NI (7.4.12) hold. Then $\delta_{0}$ is a boundary operator of the chain complex

$$
\begin{equation*}
0 \leftarrow \operatorname{Im} \bar{\delta} \stackrel{\bar{\delta}}{\leftarrow} \mathcal{P}_{\infty}^{0, n}[\overline{V F} ; Y]_{1} \stackrel{\delta_{0}}{\leftarrow} \overline{\mathcal{P}}_{\infty}^{0, n}\{0\}_{2} \stackrel{\delta_{0}}{\leftarrow} \overline{\mathcal{P}}_{\infty}^{0, n}\{0\}_{3} \tag{7.4.14}
\end{equation*}
$$

of graded densities of antifield number $\leq 3$. Let $H_{*}\left(\delta_{0}\right)$ denote its homology. We have

$$
H_{0}\left(\delta_{0}\right)=H_{0}(\bar{\delta})=0 .
$$

Furthermore, any one-cycle $\Phi$ up to a boundary takes the form (7.4.11) and, therefore, it is a $\delta_{0}$-boundary. Hence, $H_{1}\left(\delta_{0}\right)=0$, i.e., the complex (7.4.14) is one-exact.

Lemma 7.4.3: If the homology $H_{1}(\bar{\delta})$ of the complex (7.4.7) is finitely generated in the above mentioned sense, this complex can be extended to the one-exact chain complex (7.4.14) with a boundary operator whose nilpotency conditions are equivalent to the complete NI (7.4.12).

Let us consider the second homology $H_{2}\left(\delta_{0}\right)$ of the complex (7.4.14). Its two-chains read

$$
\begin{equation*}
\Phi=G+H=\sum_{0 \leq|\Lambda|} G^{r, \Lambda_{\bar{c}_{\Lambda r}} \omega+\sum_{0 \leq|\Lambda|,|\Sigma|} H^{(A, \Lambda)(B, \Sigma)} \bar{s}_{\Lambda A} \bar{s}_{\Sigma B} \omega . . . . ~ . ~} \tag{7.4.15}
\end{equation*}
$$

Its two-cycles define the first-stage NI

$$
\begin{equation*}
\delta_{0} \Phi=0, \quad \sum_{0 \leq|\Lambda|} G^{r, \Lambda} d_{\Lambda} \Delta_{r} \omega=-\bar{\delta} H \tag{7.4.16}
\end{equation*}
$$

The first-stage NI (7.4.16) are trivial either if a two-cycle $\Phi$ (7.4.15) is a $\delta_{0}$-boundary or its summand $G$ vanishes on-shell. Therefore, nontrivial first-stage NI fails to exhaust the second homology $H_{2}\left(\delta_{0}\right)$ the complex (7.4.14) in general.

Lemma 7.4.4: Non-trivial first-stage NI modulo the trivial ones are identified with elements of the homology $H_{2}\left(\delta_{0}\right)$ iff any $\bar{\delta}$-cycle $\phi \in$ $\overline{\mathcal{P}}_{\infty}^{0, n}\{0\}_{2}$ is a $\delta_{0}$-boundary.

A degenerate Lagrangian system is called reducible (resp. irreducible) if it admits (resp. does not admit) non-trivial first stage NI.

If the condition of Lemma 7.4.4 is satisfied, let us assume that nontrivial first-stage NI are finitely generated as follows. There exists a
graded projective $C^{\infty}(X)$-module $\mathcal{C}_{(1)} \subset H_{2}\left(\delta_{0}\right)$ of finite rank with a local basis $\left\{\Delta_{r_{1}} \omega\right\}$ :

$$
\begin{equation*}
\Delta_{r_{1}} \omega=\sum_{0 \leq|\Lambda|} \Delta_{r_{1}}^{r, \Lambda_{\Lambda r}} \bar{c}_{\Lambda r} \omega+h_{r_{1}} \omega \tag{7.4.17}
\end{equation*}
$$

such that any element $\Phi \in H_{2}\left(\delta_{0}\right)$ factorizes as

$$
\begin{equation*}
\Phi=\sum_{0 \leq|\Xi|} \Phi^{r_{1}, \Xi} d_{\Xi} \Delta_{r_{1}} \omega, \quad \Phi^{r_{1}, \Xi} \in \mathcal{S}_{\infty}^{0}[F ; Y] \tag{7.4.18}
\end{equation*}
$$

through elements (7.4.17) of $\mathcal{C}_{(1)}$. Thus, all non-trivial first-stage NI (7.4.16) result from the equalities

$$
\begin{equation*}
\sum_{0 \leq|\Lambda|} \Delta_{r_{1}}^{r, \Lambda} d_{\Lambda} \Delta_{r}+\bar{\delta} h_{r_{1}}=0 \tag{7.4.19}
\end{equation*}
$$

called the complete first-stage NI.
The complete first-stage NI obey second-stage NI, and so on. Iterating the arguments, one comes to the following.

A degenerate Grassmann-graded Lagrangian system $\left(\mathcal{S}_{\infty}^{*}[F ; Y], L\right)$ is called $N$-stage reducible if it admits finitely generated non-trivial $N$ stage NI, but no non-trivial $(N+1)$-stage ones. It is characterized as follows.

- There are graded vector bundles $E_{0}, \ldots, E_{N}$ over $X$ and a DBGA $\mathcal{P}_{\infty}^{*}[\overline{V F} ; Y]$ is enlarged to the DBGA

$$
\begin{equation*}
\overline{\mathcal{P}}_{\infty}^{*}\{N\}=\mathcal{P}_{\infty}^{*}\left[\overline{V F} \underset{Y}{\oplus} \bar{E}_{0} \underset{Y}{\oplus} \cdots \oplus_{Y} \bar{E}_{N} ; Y\right] \tag{7.4.20}
\end{equation*}
$$

with the local generating basis

$$
\left(s^{A}, \bar{s}_{A}, \bar{c}_{r}, \bar{c}_{r_{1}}, \ldots, \bar{c}_{r_{N}}\right)
$$

where $\bar{c}_{r_{k}}$ are Noether $k$-stage antifields of antifield number Ant $\left[\bar{c}_{r_{k}}\right]=$ $k+2$.

- The DBGA (7.4.20) admits with the nilpotent right graded derivation

$$
\begin{equation*}
\delta_{\mathrm{KT}}=\delta_{N}=\bar{\delta}+\sum_{0 \leq|\Lambda|} \overleftarrow{\partial}^{r} \Delta_{r}^{A, \Lambda} \bar{s}_{\Lambda A}+\sum_{1 \leq k \leq N} \overleftarrow{\partial}^{r_{k}} \Delta_{r_{k}} \tag{7.4.21}
\end{equation*}
$$

$$
\begin{align*}
\Delta_{r_{k}} \omega= & \sum_{0 \leq|\Lambda|} \Delta_{r_{k}}^{r_{k-1}, \Lambda} \bar{c}_{\Lambda r_{k-1}} \omega+  \tag{7.4.22}\\
& \sum_{0 \leq|\Sigma|,|\Xi|}\left(h_{r_{k}}^{\left.\left(r_{k-2}, \Sigma\right)(A, \Xi)_{\bar{c}_{\Sigma r_{k-2}}} \bar{S}_{\Xi A}+\ldots\right) \omega \in \overline{\mathcal{P}}_{\infty}^{0, n}\{k-1\}_{k+1},}\right.
\end{align*}
$$

of antifield number -1 . The index $k=-1$ here stands for $\bar{s}_{A}$. The nilpotent derivation $\delta_{\mathrm{KT}}$ (7.4.21) is called the Koszul-Tate operator.

- With this graded derivation, the module $\overline{\mathcal{P}}_{\infty}^{0, n}\{N\}_{\leq N+3}$ of densities of antifield number $\leq(N+3)$ is decomposed into the exact Koszul-Tate chain complex

$$
\begin{align*}
& 0 \leftarrow \operatorname{Im} \bar{\delta} \stackrel{\bar{\delta}}{\leftrightarrows} \mathcal{P}_{\infty}^{0, n}[\overline{V F} ; Y]_{1} \stackrel{\delta_{0}}{\leftrightarrows} \overline{\mathcal{P}}_{\infty}^{0, n}\{0\}_{2} \stackrel{\delta_{1}}{\longleftarrow} \overline{\mathcal{P}}_{\infty}^{0, n}\{1\}_{3} \cdots  \tag{7.4.23}\\
& \delta_{N-1} \overline{\mathcal{P}}_{\infty}^{0, n}\{N-1\}_{N+1} \\
& \stackrel{\delta_{\mathrm{KT}}}{\leftrightarrows} \overline{\mathcal{P}}_{\infty}^{0, n}\{N\}_{N+2} \stackrel{\delta_{\mathrm{KT}}}{\leftrightarrows} \overline{\mathcal{P}}_{\infty}^{0, n}\{N\}_{N+3}
\end{align*}
$$

which satisfies the following homology regularity condition.
Condition 7.4.5: Any $\delta_{k<N}$-cycle

$$
\phi \in \overline{\mathcal{P}}_{\infty}^{0, n}\{k\}_{k+3} \subset \overline{\mathcal{P}}_{\infty}^{0, n}\{k+1\}_{k+3}
$$

is a $\delta_{k+1}$-boundary.

- The nilpotentness $\delta_{\mathrm{KT}}^{2}=0$ of the Koszul-Tate operator (7.4.21) is equivalent to the complete non-trivial NI (7.4.12) and the complete non-trivial $(k \leq N)$-stage $N I$

$$
\begin{align*}
\sum_{0 \leq|\Lambda|} & \Delta_{r_{k}}^{r_{k-1}, \Lambda} d_{\Lambda}\left(\sum_{0 \leq|\Sigma|} \Delta_{r_{k-1}}^{r_{k-2}, \Sigma \bar{c}_{\Sigma r_{k-2}}}\right)=  \tag{7.4.24}\\
& -\bar{\delta}\left(\sum_{0 \leq|\Sigma|,|\Xi|} h_{r_{k}}^{\left.\left(r_{k-2}, \Sigma\right)(A, \Xi)_{\bar{c}_{\Sigma r_{k-2}}} \bar{s}_{\Xi A}\right)} .\right.
\end{align*}
$$

This item means the following.
Theorem 7.4.6: Any $\delta_{k}$-cocycle $\Phi \in \mathcal{P}_{\infty}^{0, n}\{k\}_{k+2}$ is a $k$-stage NI, and vice versa.

Theorem 7.4.7: Any trivial $k$-stage NI is a $\delta_{k}$-boundary $\Phi \in \mathcal{P}_{\infty}^{0, n}\{k\}_{k+2}$.

Theorem 7.4.8: All non-trivial $k$-stage NI, by assumption, factorize as

$$
\Phi=\sum_{0 \leq|\Xi|} \Phi^{r_{k}, \Xi} d_{\Xi} \Delta_{r_{k}} \omega, \quad \Phi^{r_{1}, \Xi} \in \mathcal{S}_{\infty}^{0}[F ; Y]
$$

through the complete ones (7.4.24).
It may happen that a Grassmann-graded Lagrangian field system possesses non-trivial NI of any stage. However, we restrict our consideration to $N$-reducible Lagrangian systems for a finite integer $N$.

### 7.5 Gauge symmetries

Different variants of the second Noether theorem have been suggested in order to relate reducible NI and gauge symmetries. The inverse second Noether theorem (Theorem 7.5.3), that we formulate in homology terms, associates to the Koszul-Tate complex (7.4.23) of non-trivial NI the cochain sequence (7.5.7) with the ascent operator $\mathbf{u}$ (7.5.8) whose components are non-trivial gauge and higher-stage gauge symmetries of Lagrangian system. Let us start with the following notation.

Notation 7.5.1: Given the DBGA $\overline{\mathcal{P}}_{\infty}^{*}\{N\}$ (7.4.20), we consider the DBGA

$$
\begin{equation*}
\mathcal{P}_{\infty}^{*}\{N\}=\mathcal{P}_{\infty}^{*}\left[F \underset{Y}{\oplus} E_{0} \oplus_{Y} \cdots \oplus_{Y} E_{N} ; Y\right] \tag{7.5.1}
\end{equation*}
$$

possessing the local generating basis

$$
\left(s^{A}, c^{r}, c^{r_{1}}, \ldots, c^{r_{N}}\right), \quad\left[c^{r_{k}}\right]=\left(\left[\bar{c}_{r_{k}}\right]+1\right) \bmod 2
$$

and the DBGA

$$
\begin{equation*}
P_{\infty}^{*}\{N\}=\mathcal{P}_{\infty}^{*}\left[\overline{V F} \underset{Y}{\oplus} E_{0} \oplus \cdots \underset{Y}{\oplus} E_{N} \oplus_{Y} \bar{E}_{0} \underset{Y}{\oplus} \cdots \oplus_{Y} \bar{E}_{N} ; Y\right] \tag{7.5.2}
\end{equation*}
$$

with the local generating basis

$$
\left(s^{A}, \bar{s}_{A}, c^{r}, c^{r_{1}}, \ldots, c^{r_{N}}, \bar{c}_{r}, \bar{c}_{r_{1}}, \ldots, \bar{c}_{r_{N}}\right)
$$

Their elements $c^{r_{k}}$ are called $k$-stage ghosts of ghost number $\operatorname{gh}\left[c^{r_{k}}\right]=$ $k+1$ and antifield number

$$
\operatorname{Ant}\left[c^{r_{k}}\right]=-(k+1)
$$

The $C^{\infty}(X)$-module $\mathcal{C}^{(k)}$ of $k$-stage ghosts is the density-dual of the module $\left.\mathcal{C}_{(k+1}\right)$ of $(k+1)$-stage antifields. The DBGAs $\overline{\mathcal{P}}_{\infty}^{*}\{N\}$ (7.4.20) and $\mathcal{P}_{\infty}^{*}\{N\}$ (7.5.1) are subalgebras of $P_{\infty}^{*}\{N\}$ (7.5.2). The Koszul-Tate operator $\delta_{\mathrm{KT}}(7.4 .21)$ is naturally extended to a graded derivation of the DBGA $P_{\infty}^{*}\{N\}$.

Notation 7.5.2: Any graded differential form $\phi \in \mathcal{S}_{\infty}^{*}[F ; Y]$ and any finite tuple $\left(f^{\Lambda}\right), 0 \leq|\Lambda| \leq k$, of local graded functions $f^{\Lambda} \in \mathcal{S}_{\infty}^{0}[F ; Y]$ obey the following relations:

$$
\begin{align*}
& \sum_{0 \leq|\Lambda| \leq k} f^{\Lambda} d_{\Lambda} \phi \wedge \omega=\sum_{0 \leq|\Lambda|}(-1)^{|\Lambda|} d_{\Lambda}\left(f^{\Lambda}\right) \phi \wedge \omega+d_{H} \sigma,  \tag{7.5.3}\\
& \sum_{0 \leq|\Lambda| \leq k}(-1)^{|\Lambda|} d_{\Lambda}\left(f^{\Lambda} \phi\right)=\sum_{0 \leq|\Lambda| \leq k} \eta(f)^{\Lambda} d_{\Lambda} \phi,  \tag{7.5.4}\\
& \eta(f)^{\Lambda}=\sum_{0 \leq|\Sigma| \leq k-|\Lambda|}(-1)^{|\Sigma+\Lambda|} \frac{(|\Sigma+\Lambda|)!}{|\Sigma|!|\Lambda|!} d_{\Sigma} f^{\Sigma+\Lambda},  \tag{7.5.5}\\
& \eta(\eta(f))^{\Lambda}=f^{\Lambda} . \tag{7.5.6}
\end{align*}
$$

Theorem 7.5.3: Given the Koszul-Tate complex (7.4.23), the module of graded densities $\mathcal{P}_{\infty}^{0, n}\{N\}$ is decomposed into the cochain sequence

$$
\begin{align*}
0 \rightarrow & \mathcal{S}_{\infty}^{0, n}[F ; Y] \xrightarrow{\mathbf{u}} \mathcal{P}_{\infty}^{0, n}\{N\}^{1} \xrightarrow{\mathbf{u}} \mathcal{P}_{\infty}^{0, n}\{N\}^{2} \xrightarrow{\mathbf{u}} \cdots,  \tag{7.5.7}\\
\mathbf{u}= & u+u^{(1)}+\cdots+u^{(N)}=  \tag{7.5.8}\\
& u^{A} \frac{\partial}{\partial s^{A}}+u^{r} \frac{\partial}{\partial c^{r}}+\cdots+u^{r_{N-1}} \frac{\partial}{\partial c^{r_{N-1}}},
\end{align*}
$$

graded in ghost number. Its ascent operator $\mathbf{u}$ (7.5.8) is an odd graded derivation of ghost number 1 where

$$
\begin{equation*}
u=u^{A} \frac{\partial}{\partial s^{A}}, \quad u^{A}=\sum_{0 \leq|\Lambda|} c_{\Lambda}^{r} \eta\left(\Delta_{r}^{A}\right)^{\Lambda} \tag{7.5.9}
\end{equation*}
$$

is a variational symmetry of a graded Lagrangian $L$ and the graded derivations

$$
u^{(k)}=u^{r_{k-1}} \frac{\partial}{\partial c^{r_{k-1}}}=\sum_{0 \leq|\Lambda|} c_{\Lambda}^{r_{k}} \eta\left(\Delta_{r_{k}}^{r_{k-1}}\right)^{\Lambda} \frac{\partial}{\partial c^{r_{k-1}}}, \quad k=1, \ldots, N,(7.5 .10)
$$

obey the relations

$$
\begin{align*}
& \sum_{0 \leq|\Sigma|} d_{\Sigma} u^{r_{k-1}} \frac{\partial}{\partial c_{\Sigma}^{r_{k-1}}} u^{r_{k-2}}=\bar{\delta}\left(\alpha^{r_{k-2}}\right),  \tag{7.5.11}\\
& \alpha^{r_{k-2}}=-\sum_{0 \leq|\Sigma|} \eta\left(h_{r_{k}}^{\left(r_{k-2}\right)(A, \Xi)}\right)^{\Sigma} d_{\Sigma}\left(c^{r_{k}} \bar{s}_{\Xi A}\right) .
\end{align*}
$$

A glance at the expression (7.5.9) shows that the variational symmetry $u$ is a linear differential operator on the $C^{\infty}(X)$-module $\mathcal{C}^{(0)}$ of ghosts with values into the real space $\mathfrak{g}_{L}$ of variational symmetries. Following Definition 7.2.10 extended to Lagrangian theories of odd variables, we call $u$ (7.5.9) the gauge symmetry of a graded Lagrangian $L$ which is associated to the complete NI (7.4.12).

Remark 7.5.1: In contrast with Definitions 7.2 .10 and 7.2 .12 , gauge symmetries $u(7.5 .9)$ are parameterized by ghosts, but not gauge parameters. Given a gauge symmetry $u(7.2 .28)$ defined as a derivation of the real ring $\mathcal{O}_{\infty}^{0}[Y \times E]$, one can associate to it the gauge symmetry

$$
\begin{equation*}
u=\left(\sum_{0 \leq|\Lambda| \leq m} u_{a}^{\lambda \Lambda}\left(x^{\mu}\right) c_{\Lambda}^{a}\right) \partial_{\lambda}+\left(\sum_{0 \leq|\Lambda| \leq m} u_{a}^{i \Lambda}\left(x^{\mu}, y_{\Sigma}^{j}\right) c_{\Lambda}^{a}\right) \partial_{i} \tag{7.5.12}
\end{equation*}
$$

which is an odd graded derivation of the real ring $\mathcal{S}_{\infty}^{0}[E ; Y]$, and vice versa.

Turn now to the relation (7.5.11). For $k=1$, it takes the form

$$
\sum_{0 \leq|\Sigma|} d_{\Sigma} u^{r} \frac{\partial}{\partial c_{\Sigma}^{r}} u^{A}=\bar{\delta}\left(\alpha^{A}\right)
$$

of a first-stage gauge symmetry condition on-shell which the non-trivial gauge symmetry $u(7.5 .9)$ satisfies. Therefore, one can treat the odd
graded derivation

$$
u^{(1)}=u^{r} \frac{\partial}{\partial c^{r}}, \quad u^{r}=\sum_{0 \leq|\Lambda|} c_{\Lambda}^{r_{1}} \eta\left(\Delta_{r_{1}}^{r}\right)^{\Lambda}
$$

as a first-stage gauge symmetry associated to the complete first-stage NI

$$
\sum_{0 \leq|\Lambda|} \Delta_{r_{1}}^{r, \Lambda} d_{\Lambda}\left(\sum_{0 \leq|\Sigma|} \Delta_{r}^{A, \Sigma} \bar{s}_{\Sigma A}\right)=-\bar{\delta}\left(\sum_{0 \leq|\Sigma|,|\Xi|} h_{r_{1}}^{\left.(B, \Sigma)(A, \Xi)_{\bar{s}} \bar{s}_{\Sigma B} \bar{s}_{\Xi A}\right) . . . . ~ . ~}\right.
$$

Iterating the arguments, one comes to the relation (7.5.11) which provides a $k$-stage gauge symmetry condition which is associated to the complete $k$-stage NI (7.4.24). The odd graded derivation $u_{(k)}$ (7.5.10) is called the $k$-stage gauge symmetry.

In accordance with Theorem 7.5.3, components of the ascent operator $\mathbf{u}(7.5 .8)$ are complete non-trivial gauge and higher-stage gauge symmetries. Therefore, we agree to call this operator the gauge operator.

Being a variational symmetry, a gauge symmetry $u$ (7.5.9) defines the weak conservation law (7.3.30). Let $u$ be an exact Lagrangian symmetry. In this case, the associated symmetry current

$$
\begin{equation*}
\left.\mathcal{J}_{u}=-h_{0}(u\rfloor \Xi_{L}\right) \tag{7.5.13}
\end{equation*}
$$

is conserved. The peculiarity of gauge conservation laws always is that the symmetry current $(7.5 .13)$ is reduced to a superpotential as follows.

Theorem 7.5.4: If $u(7.5 .9)$ is an exact gauge symmetry of a graded Lagrangian $L$, the corresponding conserved symmetry current $\mathcal{J}_{u}$ (7.5.13) takes the form

$$
\begin{equation*}
\mathcal{J}_{u}=W+d_{H} U=\left(W^{\mu}+d_{\nu} U^{\nu \mu}\right) \omega_{\mu} \tag{7.5.14}
\end{equation*}
$$

where the term $W$ vanishes on-shell, and $U$ is a horizontal ( $n-2$ )-form.

## Chapter 8

## Topics on commutative geometry

Several relevant topics on commutative geometry and algebraic topology are compiled in this Chapter [9, 17, 23].

### 8.1 Commutative algebra

An algebra $\mathcal{A}$ is an additive group which is additionally provided with distributive multiplication. All algebras throughout the book are associative, unless they are Lie algebras. A ring is defined to be a unital algebra, i.e., it contains a unit element $\mathbf{1} \neq 0$.

A subset $\mathcal{I}$ of an algebra $\mathcal{A}$ is called a left (resp. right) ideal if it is a subgroup of the additive group $\mathcal{A}$ and $a b \in \mathcal{I}$ (resp. $b a \in \mathcal{I}$ ) for all $a \in \mathcal{A}, b \in \mathcal{I}$. If $\mathcal{I}$ is both a left and right ideal, it is called a two-sided ideal. An ideal is a subalgebra, but a proper ideal (i.e., $\mathcal{I} \neq \mathcal{A}$ ) of a ring is not a subring because it does not contain the unit element.

Let $\mathcal{A}$ be a commutative ring. Of course, its ideals are two-sided. Its proper ideal is said to be maximal if it does not belong to another proper ideal. A commutative ring $\mathcal{A}$ is called local if it has a unique maximal ideal. This ideal consists of all non-invertible elements of $\mathcal{A}$.

Given an ideal $\mathcal{I} \subset \mathcal{A}$, the additive factor group $\mathcal{A} / \mathcal{I}$ is an algebra, called the factor algebra. If $\mathcal{A}$ is a ring, then $\mathcal{A} / \mathcal{I}$ is so. If $\mathcal{I}$ is a maximal ideal, the factor $\operatorname{ring} \mathcal{A} / \mathcal{I}$ is a field.

Given an algebra $\mathcal{A}$, an additive group $P$ is said to be a left (resp. right) $\mathcal{A}$-module if it is provided with distributive multiplication $\mathcal{A} \times P \rightarrow$ $P$ by elements of $\mathcal{A}$ such that $(a b) p=a(b p)$ (resp. $(a b) p=b(a p)$ ) for all $a, b \in \mathcal{A}$ and $p \in P$. If $\mathcal{A}$ is a ring, one additionally assumes that $\mathbf{1} p=p=p \mathbf{1}$ for all $p \in P$. Left and right module structures are usually written by means of left and right multiplications $(a, p) \rightarrow a p$ and $(a, p) \rightarrow p a$, respectively. If $P$ is both a left module over an algebra $\mathcal{A}$ and a right module over an algebra $\mathcal{A}^{\prime}$, it is called an $\left(\mathcal{A}-\mathcal{A}^{\prime}\right)$ bimodule (an $\mathcal{A}$-bimodule if $\mathcal{A}=\mathcal{A}^{\prime}$ ). If $\mathcal{A}$ is a commutative algebra, an $\mathcal{A}$-bimodule $P$ is said to be commutative if $a p=p a$ for all $a \in \mathcal{A}$ and $p \in P$. Any left or right module over a commutative algebra $\mathcal{A}$ can be brought into a commutative bimodule. Therefore, unless otherwise stated, any module over a commutative algebra $\mathcal{A}$ is called an $\mathcal{A}$-module. A module over a field is called a vector space.

If an algebra $\mathcal{A}$ is a module over a commutative ring $\mathcal{K}$, it is said to be a $\mathcal{K}$-algebra.

Hereafter, all associative algebras are assumed to be commutative.
The following are standard constructions of new modules from old ones.

- The direct sum $P_{1} \oplus P_{2}$ of $\mathcal{A}$-modules $P_{1}$ and $P_{2}$ is the additive group $P_{1} \times P_{2}$ provided with the $\mathcal{A}$-module structure

$$
a\left(p_{1}, p_{2}\right)=\left(a p_{1}, a p_{2}\right), \quad p_{1,2} \in P_{1,2}, \quad a \in \mathcal{A}
$$

Let $\left\{P_{i}\right\}_{i \in I}$ be a set of modules. Their direct sum $\oplus P_{i}$ consists of elements $\left(\ldots, p_{i}, \ldots\right)$ of the Cartesian product $\Pi P_{i}$ such that $p_{i} \neq 0$ at most for a finite number of indices $i \in I$.

- Given a submodule $Q$ of an $\mathcal{A}$-module $P$, the quotient $P / Q$ of the additive group $P$ with respect to its subgroup $Q$ also is provided with an $\mathcal{A}$-module structure. It is called a factor module.
- The set $\operatorname{Hom}_{\mathcal{A}}(P, Q)$ of $\mathcal{A}$-linear morphisms of an $\mathcal{A}$-module $P$ to
an $\mathcal{A}$-module $Q$ is naturally an $\mathcal{A}$-module. The $\mathcal{A}$-module

$$
P^{*}=\operatorname{Hom}_{\mathcal{A}}(P, \mathcal{A})
$$

is called the dual of an $\mathcal{A}$-module $P$. There is a monomorphism $P \rightarrow$ $P^{* *}$.

- The tensor product $P \otimes Q$ of $\mathcal{A}$-modules $P$ and $Q$ is an additive group which is generated by elements $p \otimes q, p \in P, q \in Q$, obeying the relations

$$
\begin{aligned}
& \left(p+p^{\prime}\right) \otimes q=p \otimes q+p^{\prime} \otimes q, \\
& p \otimes\left(q+q^{\prime}\right)=p \otimes q+p \otimes q^{\prime}, \\
& p a \otimes q=p \otimes a q, \quad p \in P, \quad q \in Q, \quad a \in \mathcal{A} .
\end{aligned}
$$

It is provided with the $\mathcal{A}$-module structure

$$
a(p \otimes q)=(a p) \otimes q=p \otimes(q a)=(p \otimes q) a .
$$

In particular, we have the following.
(i) If a ring $\mathcal{A}$ is treated as an $\mathcal{A}$-module, the tensor product $\mathcal{A} \otimes_{\mathcal{A}} Q$ is canonically isomorphic to $Q$ via the assignment

$$
\mathcal{A} \otimes_{\mathcal{A}} Q \ni a \otimes q \leftrightarrow a q \in Q .
$$

(ii) The tensor product of Abelian groups $G$ and $G^{\prime}$ is defined as their tensor product $G \otimes G^{\prime}$ as $\mathbb{Z}$-modules.
(iii) The tensor product of commutative algebras $\mathcal{A}$ and $\mathcal{A}^{\prime}$ is defined as their tensor product $\mathcal{A} \otimes \mathcal{A}^{\prime}$ as modules provided with the multiplication

$$
\left(a \otimes a^{\prime}\right)\left(b \otimes b^{\prime}\right)=\left(a a^{\prime}\right) \otimes b b^{\prime} .
$$

An $\mathcal{A}$-module $P$ is called free if it has a basis, i.e., a linearly independent subset $I \subset P$ spanning $P$ such that each element of $P$ has a unique expression as a linear combination of elements of $I$ with a finite number of non-zero coefficients from an algebra $\mathcal{A}$. Any vector space is free. Any
module is isomorphic to a quotient of a free module. A module is said to be finitely generated (or of finite rank) if it is a quotient of a free module with a finite basis.

One says that a module $P$ is projective if it is a direct summand of a free module, i.e., there exists a module $Q$ such that $P \oplus Q$ is a free module. A module $P$ is projective iff $P=\mathbf{p} S$ where $S$ is a free module and $\mathbf{p}$ is a projector of $S$, i.e., $\mathbf{p}^{2}=\mathbf{p}$.
Theorem 8.1.1: Any projective module over a local ring is free.
Now we focus on exact sequences, direct and inverse limits of modules. A composition of module morphisms

$$
P \xrightarrow{i} Q \xrightarrow{j} T
$$

is said to be exact at $Q$ if $\operatorname{Ker} j=\operatorname{Im} i$. A composition of module morphisms

$$
\begin{equation*}
0 \rightarrow P \xrightarrow{i} Q \xrightarrow{j} T \rightarrow 0 \tag{8.1.1}
\end{equation*}
$$

is called a short exact sequence if it is exact at all the terms $P, Q$, and $T$. This condition implies that: (i) $i$ is a monomorphism, (ii) $\operatorname{Ker} j=\operatorname{Im} i$, and (iii) $j$ is an epimorphism onto the quotient $T=Q / P$.

Theorem 8.1.2: Given an exact sequence of modules (8.1.1) and another $\mathcal{A}$-module $R$, the sequence of modules

$$
0 \rightarrow \operatorname{Hom}_{\mathcal{A}}(T, R) \xrightarrow{j^{*}} \operatorname{Hom}_{\mathcal{A}}(Q, R) \xrightarrow{i^{*}} \operatorname{Hom}(P, R)
$$

is exact at the first and second terms, i.e., $j^{*}$ is a monomorphism, but $i^{*}$ need not be an epimorphism.

One says that the exact sequence (8.1.1) is split if there exists a monomorphism $s: T \rightarrow Q$ such that $j \circ s=\operatorname{Id} T$ or, equivalently,

$$
Q=i(P) \oplus s(T) \cong P \oplus T
$$

Theorem 8.1.3: The exact sequence (8.1.1) is always split if $T$ is a projective module.

A directed set $I$ is a set with an order relation $<$ which satisfies the following three conditions: (i) $i<i$, for all $i \in I$; (ii) if $i<j$ and $j<k$, then $i<k$; (iii) for any $i, j \in I$, there exists $k \in I$ such that $i<k$ and $j<k$. It may happen that $i \neq j$, but $i<j$ and $j<i$ simultaneously.

A family of modules $\left\{P_{i}\right\}_{i \in I}$ (over the same algebra), indexed by a directed set $I$, is called a direct system if, for any pair $i<j$, there exists a morphism $r_{j}^{i}: P_{i} \rightarrow P_{j}$ such that

$$
r_{i}^{i}=\operatorname{Id} P_{i}, \quad r_{j}^{i} \circ r_{k}^{j}=r_{k}^{i}, \quad i<j<k .
$$

A direct system of modules admits a direct limit. This is a module $P_{\infty}$ together with morphisms $r_{\infty}^{i}: P_{i} \rightarrow P_{\infty}$ such that $r_{\infty}^{i}=r_{\infty}^{j} \circ r_{j}^{i}$ for all $i<j$. The module $P_{\infty}$ consists of elements of the direct sum $\oplus_{I} P_{i}$ modulo the identification of elements of $P_{i}$ with their images in $P_{j}$ for all $i<j$. An example of a direct system is a direct sequence

$$
\begin{equation*}
P_{0} \longrightarrow P_{1} \longrightarrow \cdots P_{i} \xrightarrow{r_{i+1}^{i}} \cdots, \quad I=\mathbb{N} . \tag{8.1.2}
\end{equation*}
$$

Note that direct limits also exist in the categories of commutative and graded commutative algebras and rings, but not in categories containing non-Abelian groups.

Theorem 8.1.4: Direct limits commute with direct sums and tensor products of modules. Namely, let $\left\{P_{i}\right\}$ and $\left\{Q_{i}\right\}$ be two direct systems of modules over the same algebra which are indexed by the same directed set $I$, and let $P_{\infty}$ and $Q_{\infty}$ be their direct limits. Then the direct limits of the direct systems $\left\{P_{i} \oplus Q_{i}\right\}$ and $\left\{P_{i} \otimes Q_{i}\right\}$ are $P_{\infty} \oplus Q_{\infty}$ and $P_{\infty} \otimes Q_{\infty}$, respectively.

Theorem 8.1.5: A morphism of a direct system $\left\{P_{i}, r_{j}^{i}\right\}_{I}$ to a direct system $\left\{Q_{i^{\prime}}, \rho_{j^{\prime}}^{i^{\prime}}\right\}_{I^{\prime}}$ consists of an order preserving map $f: I \rightarrow I^{\prime}$ and morphisms $F_{i}: P_{i} \rightarrow Q_{f(i)}$ which obey the compatibility conditions

$$
\rho_{f(j)}^{f(i)} \circ F_{i}=F_{j} \circ r_{j}^{i} .
$$

If $P_{\infty}$ and $Q_{\infty}$ are limits of these direct systems, there exists a unique morphism $F_{\infty}: P_{\infty} \rightarrow Q_{\infty}$ such that

$$
\rho_{\infty}^{f(i)} \circ F_{i}=F_{\infty} \circ r_{\infty}^{i}
$$

Theorem 8.1.6: Direct limits preserve monomorphisms and epimorphisms, i.e., if all $F_{i}: P_{i} \rightarrow Q_{f(i)}$ are monomorphisms or epimorphisms, so is $\Phi_{\infty}: P_{\infty} \rightarrow Q_{\infty}$. Let short exact sequences

$$
\begin{equation*}
0 \rightarrow P_{i} \xrightarrow{F_{i}} Q_{i} \xrightarrow{\Phi_{i}} T_{i} \rightarrow 0 \tag{8.1.3}
\end{equation*}
$$

for all $i \in I$ define a short exact sequence of direct systems of modules $\left\{P_{i}\right\}_{I},\left\{Q_{i}\right\}_{I}$, and $\left\{T_{i}\right\}_{I}$ which are indexed by the same directed set $I$. Then their direct limits form a short exact sequence

$$
\begin{equation*}
0 \rightarrow P_{\infty} \xrightarrow{F_{\infty}} Q_{\infty} \xrightarrow{\Phi_{\infty}} T_{\infty} \rightarrow 0 \tag{8.1.4}
\end{equation*}
$$

In particular, the direct limit of factor modules $Q_{i} / P_{i}$ is the factor module $Q_{\infty} / P_{\infty}$. By virtue of Theorem 8.1.4, if all the exact sequences (8.1.3) are split, the exact sequence (8.1.4) is well.

Remark 8.1.1: Let $P$ be an $\mathcal{A}$-module. We denote

$$
P^{\otimes k}=\stackrel{k}{\otimes} P
$$

Let us consider the direct system of $\mathcal{A}$-modules

$$
\mathcal{A} \longrightarrow(\mathcal{A} \oplus P) \longrightarrow \cdots\left(\mathcal{A} \oplus P \oplus \cdots \oplus P^{\otimes k}\right) \longrightarrow \cdots
$$

Its direct limit

$$
\begin{equation*}
\otimes P=\mathcal{A} \oplus P \oplus \cdots \oplus P^{\otimes k} \oplus \cdots \tag{8.1.5}
\end{equation*}
$$

is an $\mathbb{N}$-graded $\mathcal{A}$-algebra with respect to the tensor product $\otimes$. It is called the tensor algebra of a module $P$. Its quotient with respect to the ideal generated by elements

$$
p \otimes p^{\prime}+p^{\prime} \otimes p, \quad p, p^{\prime} \in P
$$

is an $\mathbb{N}$-graded commutative algebra, called the exterior algebra of $P$.
Given an inverse sequences of modules

$$
\begin{equation*}
P^{0} \longleftarrow P^{1} \longleftarrow \cdots P^{i} \stackrel{\pi_{i}^{i+1}}{\leftarrow} \cdots \tag{8.1.6}
\end{equation*}
$$

its inductive limit is a module $P^{\infty}$ together with morphisms $\pi_{i}^{\infty}: P^{\infty} \rightarrow$ $P^{i}$ such that $\pi_{i}^{\infty}=\pi_{i}^{j} \circ \pi_{j}^{\infty}$ for all $i<j$. It consists of elements $\left(\ldots, p^{i}, \ldots\right), p^{i} \in P^{i}$, of the Cartesian product $\Pi P^{i}$ such that $p^{i}=\pi_{i}^{j}\left(p^{j}\right)$ for all $i<j$.

Theorem 8.1.7: Inductive limits preserve monomorphisms, but not epimorphisms. Let exact sequences

$$
0 \rightarrow P^{i} \xrightarrow{F^{i}} Q^{i} \xrightarrow{\Phi^{i}} T^{i}, \quad i \in \mathbb{N},
$$

for all $i \in \mathbb{N}$ define an exact sequence of inverse systems of modules $\left\{P^{i}\right\}$, $\left\{Q^{i}\right\}$ and $\left\{T^{i}\right\}$. Then their inductive limits form an exact sequence

$$
0 \rightarrow P^{\infty} \xrightarrow{F^{\infty}} Q^{\infty} \xrightarrow{\Phi^{\infty}} T^{\infty} .
$$

In contrast with direct limits, the inductive ones exist in the category of groups which are not necessarily commutative.

### 8.2 Differential operators on modules

This Section addresses the notion of a linear differential operator on a module over a commutative ring.

Let $\mathcal{K}$ be a commutative ring and $\mathcal{A}$ a commutative $\mathcal{K}$-ring. Let $P$ and $Q$ be $\mathcal{A}$-modules. The $\mathcal{K}$-module $\operatorname{Hom}_{\mathcal{K}}(P, Q)$ of $\mathcal{K}$-module homomorphisms $\Phi: P \rightarrow Q$ can be endowed with the two different $\mathcal{A}$-module structures

$$
\begin{equation*}
(a \Phi)(p)=a \Phi(p), \quad(\Phi \bullet a)(p)=\Phi(a p), \quad a \in \mathcal{A}, \quad p \in P \tag{8.2.1}
\end{equation*}
$$

For the sake of convenience, we refer to the second one as the $\mathcal{A}^{\bullet}$-module structure. Let us put

$$
\begin{equation*}
\delta_{a} \Phi=a \Phi-\Phi \bullet a, \quad a \in \mathcal{A} \tag{8.2.2}
\end{equation*}
$$

Definition 8.2.1: An element $\Delta \in \operatorname{Hom}_{\mathcal{K}}(P, Q)$ is called a $Q$-valued differential operator of order $s$ on $P$ if

$$
\delta_{a_{0}} \circ \cdots \circ \delta_{a_{s}} \Delta=0
$$

for any tuple of $s+1$ elements $a_{0}, \ldots, a_{s}$ of $\mathcal{A}$. The set $\operatorname{Diff}_{s}(P, Q)$ of these operators inherits the $\mathcal{A}$ - and $\mathcal{A}^{\bullet}$-module structures (8.2.1).

In particular, zero order differential operators obey the condition

$$
\delta_{a} \Delta(p)=a \Delta(p)-\Delta(a p)=0, \quad a \in \mathcal{A}, \quad p \in P
$$

and, consequently, they coincide with $\mathcal{A}$-module morphisms $P \rightarrow Q$. A first order differential operator $\Delta$ satisfies the condition

$$
\delta_{b} \circ \delta_{a} \Delta(p)=b a \Delta(p)-b \Delta(a p)-a \Delta(b p)+\Delta(a b p)=0, \quad a, b \in \mathcal{A}
$$

The following fact reduces the study of $Q$-valued differential operators on an $\mathcal{A}$-module $P$ to that of $Q$-valued differential operators on the ring $\mathcal{A}$.

Theorem 8.2.2: Let us consider the $\mathcal{A}$-module morphism

$$
\begin{equation*}
h_{s}: \operatorname{Diff}_{s}(\mathcal{A}, Q) \rightarrow Q, \quad h_{s}(\Delta)=\Delta(\mathbf{1}) \tag{8.2.3}
\end{equation*}
$$

Any $Q$-valued $s$-order differential operator $\Delta \in \operatorname{Diff}_{s}(P, Q)$ on $P$ uniquely factorizes as

$$
\begin{equation*}
\Delta: P \xrightarrow{\mathrm{f} \Delta} \operatorname{Diff}_{s}(\mathcal{A}, Q) \xrightarrow{h_{s}} Q \tag{8.2.4}
\end{equation*}
$$

through the morphism $h_{s}$ (8.2.3) and some homomorphism

$$
\begin{equation*}
\mathfrak{f}_{\Delta}: P \rightarrow \operatorname{Diff}_{s}(\mathcal{A}, Q), \quad\left(\mathfrak{f}_{\Delta} p\right)(a)=\Delta(a p), \quad a \in \mathcal{A} \tag{8.2.5}
\end{equation*}
$$

of the $\mathcal{A}$-module $P$ to the $\mathcal{A}^{\bullet}$-module $\operatorname{Diff}{ }_{s}(\mathcal{A}, Q)$. The assignment $\Delta \rightarrow$ $f_{\Delta}$ defines the isomorphism

$$
\begin{equation*}
\operatorname{Diff}_{s}(P, Q)=\operatorname{Hom}_{\mathcal{A}-\mathcal{A}} \bullet\left(P, \operatorname{Diff}_{s}(\mathcal{A}, Q)\right) \tag{8.2.6}
\end{equation*}
$$

Let $P=\mathcal{A}$. Any zero order $Q$-valued differential operator $\Delta$ on $\mathcal{A}$ is defined by its value $\Delta(\mathbf{1})$. Then there is an isomorphism

$$
\operatorname{Diff}_{0}(\mathcal{A}, Q)=Q
$$

via the association

$$
Q \ni q \rightarrow \Delta_{q} \in \operatorname{Diff}_{0}(\mathcal{A}, Q),
$$

where $\Delta_{q}$ is given by the equality $\Delta_{q}(\mathbf{1})=q$. A first order $Q$-valued differential operator $\Delta$ on $\mathcal{A}$ fulfils the condition

$$
\Delta(a b)=b \Delta(a)+a \Delta(b)-b a \Delta(\mathbf{1}), \quad a, b \in \mathcal{A}
$$

It is called a $Q$-valued derivation of $\mathcal{A}$ if $\Delta(\mathbf{1})=0$, i.e., the Leibniz rule

$$
\begin{equation*}
\Delta(a b)=\Delta(a) b+a \Delta(b), \quad a, b \in \mathcal{A} \tag{8.2.7}
\end{equation*}
$$

holds. One obtains at once that any first order differential operator on $\mathcal{A}$ falls into the sum

$$
\Delta(a)=a \Delta(\mathbf{1})+[\Delta(a)-a \Delta(\mathbf{1})]
$$

of the zero order differential operator $a \Delta(\mathbf{1})$ and the derivation $\Delta(a)-$ $a \Delta(\mathbf{1})$. If $\partial$ is a $Q$-valued derivation of $\mathcal{A}$, then $a \partial$ is well for any $a \in \mathcal{A}$. Hence, $Q$-valued derivations of $\mathcal{A}$ constitute an $\mathcal{A}$-module $\mathfrak{d}(\mathcal{A}, Q)$, called the derivation module. There is the $\mathcal{A}$-module decomposition

$$
\begin{equation*}
\operatorname{Diff}_{1}(\mathcal{A}, Q)=Q \oplus \mathfrak{d}(\mathcal{A}, Q) \tag{8.2.8}
\end{equation*}
$$

If $P=Q=\mathcal{A}$, the derivation module $\mathfrak{d} \mathcal{A}$ of $\mathcal{A}$ also is a Lie $\mathcal{K}$-algebra with respect to the Lie bracket

$$
\begin{equation*}
\left[u, u^{\prime}\right]=u \circ u^{\prime}-u^{\prime} \circ u, \quad u, u^{\prime} \in \mathcal{A} \tag{8.2.9}
\end{equation*}
$$

Accordingly, the decomposition (8.2.8) takes the form

$$
\begin{equation*}
\operatorname{Diff}_{1}(\mathcal{A})=\mathcal{A} \oplus \mathfrak{d} \mathcal{A} \tag{8.2.10}
\end{equation*}
$$

Definition 8.2.3: A connection on an $\mathcal{A}$-module $P$ is an $\mathcal{A}$-module morphism

$$
\begin{equation*}
\mathfrak{d} \mathcal{A} \ni u \rightarrow \nabla_{u} \in \operatorname{Diff}_{1}(P, P) \tag{8.2.11}
\end{equation*}
$$

such that the first order differential operators $\nabla_{u}$ obey the Leibniz rule

$$
\begin{equation*}
\nabla_{u}(a p)=u(a) p+a \nabla_{u}(p), \quad a \in \mathcal{A}, \quad p \in P \tag{8.2.12}
\end{equation*}
$$

Though $\nabla_{u}$ (8.2.11) is called a connection, it in fact is a covariant differential on a module $P$.

Let $P$ be a commutative $\mathcal{A}$-ring and $\mathfrak{o} P$ the derivation module of $P$ as a $\mathcal{K}$-ring. The $\mathfrak{d} P$ is both a $P$ - and $\mathcal{A}$-module. Then Definition 8.2.3 is modified as follows.

Definition 8.2.4: A connection on an $\mathcal{A}$-ring $P$ is an $\mathcal{A}$-module morphism

$$
\begin{equation*}
\mathfrak{d} \mathcal{A} \ni u \rightarrow \nabla_{u} \in \mathfrak{o} P \subset \operatorname{Diff}_{1}(P, P) \tag{8.2.13}
\end{equation*}
$$

which is a connection on $P$ as an $\mathcal{A}$-module.

### 8.3 Homology and cohomology of complexes

This Section summarizes the relevant basics on complexes of modules over a commutative ring.

Let $\mathcal{K}$ be a commutative ring. A sequence

$$
\begin{equation*}
0 \leftarrow B_{0} \stackrel{\partial_{1}}{\longleftarrow} B_{1} \stackrel{\partial_{2}}{\longleftarrow} \cdots B_{p} \stackrel{\partial_{p+1}}{\rightleftarrows} \cdots \tag{8.3.1}
\end{equation*}
$$

of $\mathcal{K}$-modules $B_{p}$ and homomorphisms $\partial_{p}$ is said to be a chain complex if

$$
\partial_{p} \circ \partial_{p+1}=0, \quad p \in \mathbb{N}
$$

i.e., $\operatorname{Im} \partial_{p+1} \subset \operatorname{Ker} \partial_{p}$. Homomorphisms $\partial_{p}$ are called boundary operators. Elements of a module $B_{p}$, its submodules $\operatorname{Ker} \partial_{p} \subset B_{p}$ and $\operatorname{Im} \partial_{p+1} \subset \operatorname{Ker} \partial_{p}$ are called $p$-chains, $p$-cycles and $p$-boundaries, respectively. The $p$-th homology group of the chain complex $B_{*}$ (8.3.1) is the factor module

$$
H_{p}\left(B_{*}\right)=\operatorname{Ker} \partial_{p} / \operatorname{Im} \partial_{p+1} .
$$

It is a $\mathcal{K}$-module. In particular, we have $H_{0}\left(B_{*}\right)=B_{0} / \operatorname{Im} \partial_{1}$. The chain complex (8.3.1) is exact at a term $B_{p}$ iff $H_{p}\left(B_{*}\right)=0$. This complex is said to be $k$-exact if its homology groups $H_{p \leq k}\left(B_{*}\right)$ are trivial. It is called exact if all its homology groups are trivial, i.e., it is an exact sequence.

A sequence

$$
\begin{equation*}
0 \rightarrow B^{0} \xrightarrow{\delta^{0}} B^{1} \xrightarrow{\delta^{1}} \cdots B^{p} \xrightarrow{\delta^{p}} \cdots \tag{8.3.2}
\end{equation*}
$$

of modules $B^{p}$ and their homomorphisms $\delta^{p}$ is said to be a cochain complex (or, simply, a complex) if

$$
\delta^{p+1} \circ \delta^{p}=0, \quad p \in \mathbb{N},
$$

i.e., $\operatorname{Im} \delta^{p} \subset \operatorname{Ker} \delta^{p+1}$. The homomorphisms $\delta^{p}$ are called coboundary operators. Elements of a module $B^{p}$, its submodules $\operatorname{Ker} \delta^{p} \subset B^{p}$ and $\operatorname{Im} \delta^{p-1}$ are called $p$-cochains, $p$-cocycles and $p$-coboundaries, respectively. The $p$-th cohomology group of the complex $B^{*}$ (8.3.2) is the factor module

$$
H^{p}\left(B^{*}\right)=\operatorname{Ker} \delta^{p} / \operatorname{Im} \delta^{p-1} .
$$

It is a $\mathcal{K}$-module. In particular, $H^{0}\left(B^{*}\right)=\operatorname{Ker} \delta^{0}$. The complex (8.3.2) is exact at a term $B^{p}$ iff $H^{p}\left(B^{*}\right)=0$. This complex is an exact sequence if all its cohomology groups are trivial.

A complex $\left(B^{*}, \delta^{*}\right)$ is called acyclic if its cohomology groups $H^{0<p}\left(B^{*}\right)$ are trivial. A complex $\left(B^{*}, \delta^{*}\right)$ is said to be a resolution of a module $B$ if it is acyclic and $H^{0}\left(B^{*}\right)=\operatorname{Ker} \delta^{0}=B$.

The following are the standard constructions of new complexes from old ones.

- Given complexes $\left(B_{1}^{*}, \delta_{1}^{*}\right)$ and $\left(B_{2}^{*}, \delta_{2}^{*}\right)$, their direct sum $B_{1}^{*} \oplus B_{2}^{*}$ is a complex of modules

$$
\left(B_{1}^{*} \oplus B_{2}^{*}\right)^{p}=B_{1}^{p} \oplus B_{2}^{p}
$$

with respect to the coboundary operators

$$
\delta_{\oplus}^{p}\left(b_{1}^{p}+b_{2}^{p}\right)=\delta_{1}^{p} b_{1}^{p}+\delta_{2}^{p} b_{2}^{p} .
$$

- Given a subcomplex $\left(C^{*}, \delta^{*}\right)$ of a complex $\left(B^{*}, \delta^{*}\right)$, the factor complex $B^{*} / C^{*}$ is defined as a complex of factor modules $B^{p} / C^{p}$ provided with the coboundary operators $\delta^{p}\left[b^{p}\right]=\left[\delta^{p} b^{p}\right]$, where $\left[b^{p}\right] \in B^{p} / C^{p}$ denotes the coset of an element $b^{p}$.
- Given complexes $\left(B_{1}^{*}, \delta_{1}^{*}\right)$ and $\left(B_{2}^{*}, \delta_{2}^{*}\right)$, their tensor product $B_{1}^{*} \otimes B_{2}^{*}$ is a complex of modules

$$
\left(B_{1}^{*} \otimes B_{2}^{*}\right)^{p}=\underset{k+r=p}{\oplus} B_{1}^{k} \otimes B_{2}^{r}
$$

with respect to the coboundary operators

$$
\delta_{\otimes}^{p}\left(b_{1}^{k} \otimes b_{2}^{r}\right)=\left(\delta_{1}^{k} b_{1}^{k}\right) \otimes b_{2}^{r}+(-1)^{k} b_{1}^{k} \otimes\left(\delta_{2}^{r} b_{2}^{r}\right) .
$$

A cochain morphism of complexes

$$
\begin{equation*}
\gamma: B_{1}^{*} \rightarrow B_{2}^{*} \tag{8.3.3}
\end{equation*}
$$

is defined as a family of degree-preserving homomorphisms

$$
\gamma^{p}: B_{1}^{p} \rightarrow B_{2}^{p}, \quad p \in \mathbb{N},
$$

such that

$$
\delta_{2}^{p} \circ \gamma^{p}=\gamma^{p+1} \circ \delta_{1}^{p}, \quad p \in \mathbb{N}
$$

It follows that if $b^{p} \in B_{1}^{p}$ is a cocycle or a coboundary, then $\gamma^{p}\left(b^{p}\right) \in B_{2}^{p}$ is so. Therefore, the cochain morphism of complexes (8.3.3) yields an induced homomorphism of their cohomology groups

$$
[\gamma]^{*}: H^{*}\left(B_{1}^{*}\right) \rightarrow H^{*}\left(B_{2}^{*}\right) .
$$

Let short exact sequences

$$
0 \rightarrow C^{p} \xrightarrow{\gamma_{p}} B^{p} \xrightarrow{\zeta_{p}} F^{p} \rightarrow 0
$$

for all $p \in \mathbb{N}$ define a short exact sequence of complexes

$$
\begin{equation*}
0 \rightarrow C^{*} \xrightarrow{\gamma} B^{*} \xrightarrow{\zeta} F^{*} \rightarrow 0 \tag{8.3.4}
\end{equation*}
$$

where $\gamma$ is a cochain monomorphism and $\zeta$ is a cochain epimorphism onto the quotient $F^{*}=B^{*} / C^{*}$.

Theorem 8.3.1: The short exact sequence of complexes (8.3.4) yields the long exact sequence of their cohomology groups

$$
\begin{aligned}
& 0 \rightarrow H^{0}\left(C^{*}\right) \xrightarrow{[\gamma]^{0}} H^{0}\left(B^{*}\right) \xrightarrow{[\zeta]^{0}} H^{0}\left(F^{*}\right) \xrightarrow{\tau^{0}} H^{1}\left(C^{*}\right) \longrightarrow \cdots \\
&\left.\longrightarrow H^{p}\left(C^{*}\right) \xrightarrow{[\gamma]^{p}} H^{p}\left(B^{*}\right) \xrightarrow{[\zeta]^{p}} H^{p}\left(F^{*}\right) \xrightarrow{\tau^{p}} H^{p+1}\left(C^{*}\right) \longrightarrow \cdots .3\right) \\
& \hline
\end{aligned}
$$

Theorem 8.3.2: A direct sequence of complexes

$$
\begin{equation*}
B_{0}^{*} \longrightarrow B_{1}^{*} \longrightarrow \cdots B_{k}^{*} \xrightarrow{\gamma_{k+1}^{k}} B_{k+1}^{*} \longrightarrow \cdots \tag{8.3.6}
\end{equation*}
$$

admits a direct limit $B_{\infty}^{*}$ which is a complex whose cohomology $H^{*}\left(B_{\infty}^{*}\right)$ is a direct limit of the direct sequence of cohomology groups

$$
H^{*}\left(B_{0}^{*}\right) \longrightarrow H^{*}\left(B_{1}^{*}\right) \longrightarrow \cdots H^{*}\left(B_{k}^{*}\right) \xrightarrow{\left[l_{k+1}^{k}\right]} H^{*}\left(B_{k+1}^{*}\right) \longrightarrow \cdots .
$$

### 8.4 Differential calculus over a commutative ring

Let $\mathfrak{g}$ be a Lie algebra over a commutative ring $\mathcal{K}$. Let $\mathfrak{g}$ act on a $\mathcal{K}$ module $P$ on the left such that

$$
\left[\varepsilon, \varepsilon^{\prime}\right] p=\left(\varepsilon \circ \varepsilon^{\prime}-\varepsilon^{\prime} \circ \varepsilon\right) p, \quad \varepsilon, \varepsilon^{\prime} \in \mathfrak{g}
$$

Then one calls $P$ the Lie algebra $\mathfrak{g}$-module. Let us consider $\mathcal{K}$-multilinear skew-symmetric maps

$$
c^{k}: \stackrel{k}{\times} \mathfrak{g} \rightarrow P
$$

They form a $\mathfrak{g}$-module $C^{k}[\mathfrak{g} ; P]$. Let us put $C^{0}[\mathfrak{g} ; P]=P$. We obtain the cochain complex

$$
\begin{equation*}
0 \rightarrow P \xrightarrow{\delta^{0}} C^{1}[\mathfrak{g} ; P] \xrightarrow{\delta^{1}} \cdots C^{k}[\mathfrak{g} ; P] \xrightarrow{\delta^{k}} \cdots \tag{8.4.1}
\end{equation*}
$$

with respect to the Chevalley-Eilenberg coboundary operators

$$
\begin{align*}
& \delta^{k} c^{k}\left(\varepsilon_{0}, \ldots, \varepsilon_{k}\right)=\sum_{i=0}^{k}(-1)^{i} \varepsilon_{i} c^{k}\left(\varepsilon_{0}, \ldots, \widehat{\varepsilon}_{i}, \ldots, \varepsilon_{k}\right)+  \tag{8.4.2}\\
& \sum_{1 \leq i<j \leq k}(-1)^{i+j} c^{k}\left(\left[\varepsilon_{i}, \varepsilon_{j}\right], \varepsilon_{0}, \ldots, \widehat{\varepsilon}_{i}, \ldots, \widehat{\varepsilon}_{j}, \ldots, \varepsilon_{k}\right)
\end{align*}
$$

where the caret ${ }^{\text {- denotes omission. For instance, we have }}$

$$
\begin{align*}
& \delta^{0} p\left(\varepsilon_{0}\right)=\varepsilon_{0} p  \tag{8.4.3}\\
& \delta^{1} c^{1}\left(\varepsilon_{0}, \varepsilon_{1}\right)=\varepsilon_{0} c^{1}\left(\varepsilon_{1}\right)-\varepsilon_{1} c^{1}\left(\varepsilon_{0}\right)-c^{1}\left(\left[\varepsilon_{0}, \varepsilon_{1}\right]\right) \tag{8.4.4}
\end{align*}
$$

The complex (8.4.1) is called the Chevalley-Eilenberg complex, and its cohomology $H^{*}(\mathfrak{g}, P)$ is the Chevalley-Eilenberg cohomology of a Lie algebra $\mathfrak{g}$ with coefficients in $P$.

Let $\mathcal{A}$ be a commutative $\mathcal{K}$-ring. Since the derivation module $\mathfrak{d} \mathcal{A}$ of $\mathcal{A}$ is a Lie $\mathcal{K}$-algebra, one can associate to $\mathcal{A}$ the Chevalley-Eilenberg complex $C^{*}[\mathfrak{d} \mathcal{A} ; \mathcal{A}]$. Its subcomplex of $\mathcal{A}$-multilinear maps is a DGA, also called the differential calculus over $\mathcal{A}$. By a gradation throughout this Section is meant the $\mathbb{N}$-gradation.

A graded algebra $\Omega^{*}$ over a commutative ring $\mathcal{K}$ is defined as a direct $\operatorname{sum} \Omega^{*}=\underset{k}{\oplus} \Omega^{k}$ of $\mathcal{K}$-modules $\Omega^{k}$, provided with an associative multiplication law $\alpha \cdot \beta, \alpha, \beta \in \Omega^{*}$, such that $\alpha \cdot \beta \in \Omega^{|\alpha|+|\beta|}$, where $|\alpha|$ denotes the degree of an element $\alpha \in \Omega^{|\alpha|}$. In particular, it follows that $\Omega^{0}$ is a (non-commutative) $\mathcal{K}$-algebra $\mathcal{A}$, while $\Omega^{k>0}$ are $\mathcal{A}$-bimodules and $\Omega^{*}$ is an $(\mathcal{A}-\mathcal{A})$-algebra. A graded algebra is said to be graded commutative if

$$
\alpha \cdot \beta=(-1)^{|\alpha||\beta|} \beta \cdot \alpha, \quad \alpha, \beta \in \Omega^{*}
$$

A graded algebra $\Omega^{*}$ is called the differential graded algebra (DGA) or the differential calculus over $\mathcal{A}$ if it is a cochain complex of $\mathcal{K}$-modules

$$
\begin{equation*}
0 \rightarrow \mathcal{K} \longrightarrow \mathcal{A} \xrightarrow{\delta} \Omega^{1} \xrightarrow{\delta} \cdots \Omega^{k} \xrightarrow{\delta} \cdots \tag{8.4.5}
\end{equation*}
$$

relative to a coboundary operator $\delta$ which obeys the graded Leibniz rule

$$
\begin{equation*}
\delta(\alpha \cdot \beta)=\delta \alpha \cdot \beta+(-1)^{|\alpha|} \alpha \cdot \delta \beta . \tag{8.4.6}
\end{equation*}
$$

In particular, $\delta: \mathcal{A} \rightarrow \Omega^{1}$ is a $\Omega^{1}$-valued derivation of a $\mathcal{K}$-algebra $\mathcal{A}$. The cochain complex (8.4.5) is said to be the abstract de Rham complex of the DGA $\left(\Omega^{*}, \delta\right)$. Cohomology $H^{*}\left(\Omega^{*}\right)$ of the complex (8.4.5) is called the abstract de Rham cohomology.

A morphism $\gamma$ between two DGAs $\left(\Omega^{*}, \delta\right)$ and $\left(\Omega^{\prime *}, \delta^{\prime}\right)$ is defined as a cochain morphism, i.e., $\gamma \circ \delta=\gamma \circ \delta^{\prime}$. It yields the corresponding morphism of the abstract de Rham cohomology groups of these algebras.

One considers the minimal differential graded subalgebra $\Omega^{*} \mathcal{A}$ of the DGA $\Omega^{*}$ which contains $\mathcal{A}$. Seen as an $(\mathcal{A}-\mathcal{A})$-algebra, it is generated by the elements $\delta a, a \in \mathcal{A}$, and consists of monomials

$$
\alpha=a_{0} \delta a_{1} \cdots \delta a_{k}, \quad a_{i} \in \mathcal{A}
$$

whose product obeys the juxtaposition rule

$$
\left(a_{0} \delta a_{1}\right) \cdot\left(b_{0} \delta b_{1}\right)=a_{0} \delta\left(a_{1} b_{0}\right) \cdot \delta b_{1}-a_{0} a_{1} \delta b_{0} \cdot \delta b_{1}
$$

in accordance with the equality (8.4.6). The $\operatorname{DGA}\left(\Omega^{*} \mathcal{A}, \delta\right)$ is called the minimal differential calculus over $\mathcal{A}$.

Let now $\mathcal{A}$ be a commutative $\mathcal{K}$-ring possessing a non-trivial Lie algebra $\mathfrak{o} \mathcal{A}$ of derivations. We consider the extended Chevalley-Eilenberg complex

$$
0 \rightarrow \mathcal{K} \xrightarrow{\text { in }} C^{*}[\mathfrak{o A} ; \mathcal{A}]
$$

of the Lie algebra $\mathfrak{O} \mathcal{A}$ with coefficients in the $\operatorname{ring} \mathcal{A}$, regarded as a $\mathfrak{O} \mathcal{A}-$ module. It is easily justified that this complex contains a subcomplex $\mathcal{O}^{*}[\mathfrak{O} \mathcal{A}]$ of $\mathcal{A}$-multilinear skew-symmetric maps

$$
\begin{equation*}
\phi^{k}: \stackrel{k}{\times} \mathfrak{d} \mathcal{A} \rightarrow \mathcal{A} \tag{8.4.7}
\end{equation*}
$$

with respect to the Chevalley-Eilenberg coboundary operator

$$
\begin{gather*}
d \phi\left(u_{0}, \ldots, u_{k}\right)=\sum_{i=0}^{k}(-1)^{i} u_{i}\left(\phi\left(u_{0}, \ldots, \widehat{u_{i}}, \ldots, u_{k}\right)\right)+  \tag{8.4.8}\\
\sum_{i<j}(-1)^{i+j} \phi\left(\left[u_{i}, u_{j}\right], u_{0}, \ldots, \widehat{u}_{i}, \ldots, \widehat{u}_{j}, \ldots, u_{k}\right)
\end{gather*}
$$

In particular, we have

$$
\begin{aligned}
& (d a)(u)=u(a), \quad a \in \mathcal{A}, \quad u \in \mathfrak{d} \mathcal{A} \\
& (d \phi)\left(u_{0}, u_{1}\right)=u_{0}\left(\phi\left(u_{1}\right)\right)-u_{1}\left(\phi\left(u_{0}\right)\right)-\phi\left(\left[u_{0}, u_{1}\right]\right), \quad \phi \in \mathcal{O}^{1}[\mathfrak{d} \mathcal{A}] \\
& \mathcal{O}^{0}[\mathfrak{d} \mathcal{A}]=\mathcal{A}, \\
& \mathcal{O}^{1}[\mathfrak{d} \mathcal{A}]=\operatorname{Hom}_{\mathcal{A}}(\mathfrak{d} \mathcal{A}, \mathcal{A})=\mathfrak{d} \mathcal{A}^{*}
\end{aligned}
$$

It follows that $d(\mathbf{1})=0$ and $d$ is a $\mathcal{O}^{1}[\mathfrak{o} \mathcal{A}]$-valued derivation of $\mathcal{A}$.
The graded module $\mathcal{O}^{*}[\mathfrak{D} \mathcal{A}]$ is provided with the structure of a graded $\mathcal{A}$-algebra with respect to the exterior product

$$
\begin{align*}
& \phi \wedge \phi^{\prime}\left(u_{1}, \ldots, u_{r+s}\right)=  \tag{8.4.9}\\
& \quad \sum_{\substack{i_{1}<\cdots<i_{r} ; j_{1}<\cdots<j_{s}}}^{\operatorname{sgn}_{1 \cdots r+s}^{i_{1} \cdots i_{r} j_{1} \cdots j_{s}} \phi\left(u_{i_{1}}, \ldots, u_{i_{r}}\right) \phi^{\prime}\left(u_{j_{1}}, \ldots, u_{j_{s}}\right),} \\
& \phi \in \mathcal{O}^{r}\left[\mathfrak{d A}, \quad \phi^{\prime} \in \mathcal{O}^{s}[\mathfrak{d} \mathcal{A}], \quad u_{k} \in \mathfrak{d} \mathcal{A},\right.
\end{align*}
$$

where sgn... is the sign of a permutation. This product obeys the relations

$$
\begin{align*}
& d\left(\phi \wedge \phi^{\prime}\right)=d(\phi) \wedge \phi^{\prime}+(-1)^{|\phi|} \phi \wedge d\left(\phi^{\prime}\right), \quad \phi, \phi^{\prime} \in \mathcal{O}^{*}[\mathcal{O A}] \\
& \phi \wedge \phi^{\prime}=(-1)^{|\phi|\left|\phi^{\prime}\right|} \phi^{\prime} \wedge \phi \tag{8.4.10}
\end{align*}
$$

By virtue of the first one, $\mathcal{O}^{*}[\mathfrak{d} \mathcal{A}]$ is a differential graded $\mathcal{K}$-algebra, called the Chevalley-Eilenberg differential calculus over a $\mathcal{K}$-ring $\mathcal{A}$. The relation (8.4.10) shows that $\mathcal{O}^{*}[\mathfrak{o} \mathcal{A}]$ is a graded commutative algebra.

The minimal Chevalley-Eilenberg differential calculus $\mathcal{O}^{*} \mathcal{A}$ over a $\operatorname{ring} \mathcal{A}$ consists of the monomials

$$
a_{0} d a_{1} \wedge \cdots \wedge d a_{k}, \quad a_{i} \in \mathcal{A}
$$

Its complex

$$
\begin{equation*}
0 \rightarrow \mathcal{K} \longrightarrow \mathcal{A} \xrightarrow{d} \mathcal{O}^{1} \mathcal{A} \xrightarrow{d} \cdots \mathcal{O}^{k} \mathcal{A} \xrightarrow{d} \cdots \tag{8.4.11}
\end{equation*}
$$

is said to be the de Rham complex of a $\mathcal{K}$-ring $\mathcal{A}$, and its cohomology $H^{*}(\mathcal{A})$ is called the de Rham cohomology of $\mathcal{A}$.

### 8.5 Sheaf cohomology

A sheaf on a topological space $X$ is a continuous fibre bundle $\pi: S \rightarrow X$ in modules over a commutative ring $\mathcal{K}$, where the surjection $\pi$ is a local homeomorphism and fibres $S_{x}, x \in X$, called the stalks, are provided with the discrete topology. Global sections of a sheaf $S$ make up a $\mathcal{K}$ module $S(X)$, called the structure module of $S$.

Any sheaf is generated by a presheaf. A presheaf $S_{\{U\}}$ on a topological space $X$ is defined if a module $S_{U}$ over a commutative ring $\mathcal{K}$ is assigned to every open subset $U \subset X\left(S_{\emptyset}=0\right)$ and if, for any pair of open subsets $V \subset U$, there exists the restriction morphism $r_{V}^{U}: S_{U} \rightarrow S_{V}$ such that

$$
r_{U}^{U}=\operatorname{Id} S_{U}, \quad r_{W}^{U}=r_{W}^{V} r_{V}^{U}, \quad W \subset V \subset U
$$

Every presheaf $S_{\{U\}}$ on a topological space $X$ yields a sheaf on $X$ whose stalk $S_{x}$ at a point $x \in X$ is the direct limit of the modules $S_{U}, x \in U$, with respect to the restriction morphisms $r_{V}^{U}$. It means that, for each open neighborhood $U$ of a point $x$, every element $s \in S_{U}$ determines an element $s_{x} \in S_{x}$, called the germ of $s$ at $x$. Two elements $s \in S_{U}$ and $s^{\prime} \in S_{V}$ belong to the same germ at $x$ iff there exists an open neighborhood $W \subset U \cap V$ of $x$ such that $r_{W}^{U} s=r_{W}^{V} s^{\prime}$.

Example 8.5.1: Let $C_{\{U\}}^{0}$ be the presheaf of continuous real functions on a topological space $X$. Two such functions $s$ and $s^{\prime}$ define the same germ $s_{x}$ if they coincide on an open neighborhood of $x$. Hence, we obtain the sheaf $C_{X}^{0}$ of continuous functions on $X$. Similarly, the sheaf $C_{X}^{\infty}$ of smooth functions on a smooth manifold $X$ is defined. Let us also mention the presheaf of real functions which are constant on connected open subsets of $X$. It generates the constant sheaf on $X$ denoted by $\mathbb{R}$.

Different presheaves may generate the same sheaf. Conversely, every sheaf $S$ defines a presheaf $S(\{U\})$ of modules $S(U)$ of its local sections. It is called the canonical presheaf of the sheaf $S$. If a sheaf $S$ is constructed
from a presheaf $S_{\{U\}}$, there are natural module morphisms

$$
S_{U} \ni s \rightarrow s(U) \in S(U), \quad s(x)=s_{x}, \quad x \in U,
$$

which are neither monomorphisms nor epimorphisms in general. For instance, it may happen that a non-zero presheaf defines a zero sheaf. The sheaf generated by the canonical presheaf of a sheaf $S$ coincides with $S$.

A direct sum and a tensor product of presheaves (as families of modules) and sheaves (as fibre bundles in modules) are naturally defined. By virtue of Theorem 8.1.4, a direct sum (resp. a tensor product) of presheaves generates a direct sum (resp. a tensor product) of the corresponding sheaves.

Remark 8.5.2: In a different terminology, a sheaf is introduced as a presheaf which satisfies the following additional axioms.
(S1) Suppose that $U \subset X$ is an open subset and $\left\{U_{\alpha}\right\}$ is its open cover. If $s, s^{\prime} \in S_{U}$ obey the condition

$$
r_{U_{\alpha}}^{U}(s)=r_{U_{\alpha}}^{U}\left(s^{\prime}\right)
$$

for all $U_{\alpha}$, then $s=s^{\prime}$.
(S2) Let $U$ and $\left\{U_{\alpha}\right\}$ be as in previous item. Suppose that we are given a family of presheaf elements $\left\{s_{\alpha} \in S_{U_{\alpha}}\right\}$ such that

$$
r_{U_{\alpha} \cap U_{\lambda}}^{U_{\alpha}}\left(s_{\alpha}\right)=r_{U_{\alpha} \cap U_{\lambda}}^{U_{\lambda}}\left(s_{\lambda}\right)
$$

for all $U_{\alpha}, U_{\lambda}$. Then there exists a presheaf element $s \in S_{U}$ such that $s_{\alpha}=r_{U_{\alpha}}^{U}(s)$.
Canonical presheaves are in one-to-one correspondence with presheaves obeying these axioms. For instance, presheaves of continuous, smooth and locally constant functions in Example 8.5.1 satisfy the axioms (S1) - (S2).

Remark 8.5.3: The notion of a sheaf can be extended to sets, but not to non-commutative groups. One can consider a presheaf of such groups,
but it generates a sheaf of sets because a direct limit of non-commutative groups need not be a group. The first (but not higher) cohomology of $X$ with coefficients in this sheaf is defined.

A morphism of a presheaf $S_{\{U\}}$ to a presheaf $S_{\{U\}}^{\prime}$ on the same topological space $X$ is defined as a set of module morphisms $\gamma_{U}: S_{U} \rightarrow S_{U}^{\prime}$ which commute with restriction morphisms. A morphism of presheaves yields a morphism of sheaves generated by these presheaves. This is a bundle morphism over $X$ such that $\gamma_{x}: S_{x} \rightarrow S_{x}^{\prime}$ is the direct limit of morphisms $\gamma_{U}, x \in U$. Conversely, any morphism of sheaves $S \rightarrow S^{\prime}$ on a topological space $X$ yields a morphism of canonical presheaves of local sections of these sheaves. Let $\operatorname{Hom}\left(\left.S\right|_{U},\left.S^{\prime}\right|_{U}\right)$ be the commutative group of sheaf morphisms $\left.\left.S\right|_{U} \rightarrow S^{\prime}\right|_{U}$ for any open subset $U \subset X$. These groups are assembled into a presheaf, and define the sheaf $\operatorname{Hom}\left(S, S^{\prime}\right)$ on $X$. There is a monomorphism

$$
\begin{equation*}
\operatorname{Hom}\left(S, S^{\prime}\right)(U) \rightarrow \operatorname{Hom}\left(S(U), S^{\prime}(U)\right) \tag{8.5.1}
\end{equation*}
$$

which need not be an isomorphism.
By virtue of Theorem 8.1.6, if a presheaf morphism is a monomorphism or an epimorphism, so is the corresponding sheaf morphism. Furthermore, the following holds.

Theorem 8.5.1: A short exact sequence

$$
\begin{equation*}
0 \rightarrow S_{\{U\}}^{\prime} \rightarrow S_{\{U\}} \rightarrow S_{\{U\}}^{\prime \prime} \rightarrow 0 \tag{8.5.2}
\end{equation*}
$$

of presheaves on the same topological space yields the short exact sequence of sheaves generated by these presheaves

$$
\begin{equation*}
0 \rightarrow S^{\prime} \rightarrow S \rightarrow S^{\prime \prime} \rightarrow 0 \tag{8.5.3}
\end{equation*}
$$

where the factor sheaf $S^{\prime \prime}=S / S^{\prime}$ is isomorphic to that generated by the factor presheaf $S_{\{U\}}^{\prime \prime}=S_{\{U\}} / S_{\{U\}}^{\prime}$. If the exact sequence of presheaves (8.5.2) is split, i.e.,

$$
S_{\{U\}} \cong S_{\{U\}}^{\prime} \oplus S_{\{U\}}^{\prime \prime}
$$

the corresponding splitting

$$
S \cong S^{\prime} \oplus S^{\prime \prime}
$$

of the exact sequence of sheaves (8.5.3) holds.
The converse is more intricate. A sheaf morphism induces a morphism of the corresponding canonical presheaves. If $S \rightarrow S^{\prime}$ is a monomorphism,

$$
S(\{U\}) \rightarrow S^{\prime}(\{U\})
$$

also is a monomorphism. However, if $S \rightarrow S^{\prime}$ is an epimorphism,

$$
S(\{U\}) \rightarrow S^{\prime}(\{U\})
$$

need not be so. Therefore, the short exact sequence (8.5.3) of sheaves yields the exact sequence of the canonical presheaves

$$
\begin{equation*}
0 \rightarrow S^{\prime}(\{U\}) \rightarrow S(\{U\}) \rightarrow S^{\prime \prime}(\{U\}) \tag{8.5.4}
\end{equation*}
$$

where $S(\{U\}) \rightarrow S^{\prime \prime}(\{U\})$ is not necessarily an epimorphism. At the same time, there is the short exact sequence of presheaves

$$
\begin{equation*}
0 \rightarrow S^{\prime}(\{U\}) \rightarrow S(\{U\}) \rightarrow S_{\{U\}}^{\prime \prime} \rightarrow 0 \tag{8.5.5}
\end{equation*}
$$

where the factor presheaf

$$
S_{\{U\}}^{\prime \prime}=S(\{U\}) / S^{\prime}(\{U\})
$$

generates the factor sheaf $S^{\prime \prime}=S / S^{\prime}$, but need not be its canonical presheaf.

Let us turn now to sheaf cohomology. Note that only proper covers are considered.

Let $S_{\{U\}}$ be a presheaf of modules on a topological space $X$, and let $\mathfrak{U}=\left\{U_{i}\right\}_{i \in I}$ be an open cover of $X$. One constructs a cochain complex where a $p$-cochain is defined as a function $s^{p}$ which associates an element

$$
\begin{equation*}
s^{p}\left(i_{0}, \ldots, i_{p}\right) \in S_{U_{i_{0}} \cap \cdots \cap U_{i_{p}}} \tag{8.5.6}
\end{equation*}
$$

to each $(p+1)$-tuple $\left(i_{0}, \ldots, i_{p}\right)$ of indices in $I$. These $p$-cochains are assembled into a module $C^{p}\left(\mathfrak{U}, S_{\{U\}}\right)$. Let us introduce the coboundary operator

$$
\begin{align*}
& \delta^{p}: C^{p}\left(\mathfrak{U}, S_{\{U\}}\right) \rightarrow C^{p+1}\left(\mathfrak{U}, S_{\{U\}}\right), \\
& \delta^{p} s^{p}\left(i_{0}, \ldots, i_{p+1}\right)=\sum_{k=0}^{p+1}(-1)^{k} r_{W}^{W_{k}} s^{p}\left(i_{0}, \ldots, \widehat{i}_{k}, \ldots, i_{p+1}\right),  \tag{8.5.7}\\
& W=U_{i_{0}} \cap \ldots \cap U_{i_{p+1}}, \quad W_{k}=U_{i_{0}} \cap \cdots \cap \widehat{U}_{i_{k}} \cap \cdots \cap U_{i_{p+1}} .
\end{align*}
$$

One can easily check that $\delta^{p+1} \circ \delta^{p}=0$. Thus, we obtain the cochain complex of modules

$$
\begin{equation*}
0 \rightarrow C^{0}\left(\mathfrak{U}, S_{\{U\}}\right) \xrightarrow{\delta^{0}} \cdots C^{p}\left(\mathfrak{U}, S_{\{U\}}\right) \xrightarrow{\delta^{p}} C^{p+1}\left(\mathfrak{U}, S_{\{U\}}\right) \cdots . \tag{8.5.8}
\end{equation*}
$$

Its cohomology groups

$$
H^{p}\left(\mathfrak{U} ; S_{\{U\}}\right)=\operatorname{Ker} \delta^{p} / \operatorname{Im} \delta^{p-1}
$$

are modules. Of course, they depend on an open cover $\mathfrak{U}$ of $X$.
Let $\mathfrak{U}^{\prime}$ be a refinement of the cover $\mathfrak{U}$. Then there is a morphism of cohomology groups

$$
\begin{equation*}
H^{*}\left(\mathfrak{U} ; S_{\{U\}}\right) \rightarrow H^{*}\left(\mathfrak{U}^{\prime} ; S_{\{U\}}\right) \tag{8.5.9}
\end{equation*}
$$

Let us take the direct limit of cohomology groups $H^{*}\left(\mathfrak{U} ; S_{\{U\}}\right)$ relative to these morphisms, where $\mathfrak{U}$ runs through all open covers of $X$. This limit $H^{*}\left(X ; S_{\{U\}}\right)$ is called the cohomology of $X$ with coefficients in the presheaf $S_{\{U\}}$.

Let $S$ be a sheaf on a topological space $X$. Cohomology of $X$ with coefficients in $S$ or, simply, sheaf cohomology of $X$ is defined as cohomology

$$
H^{*}(X ; S)=H^{*}(X ; S(\{U\}))
$$

with coefficients in the canonical presheaf $S(\{U\})$ of the sheaf $S$.
In this case, a $p$-cochain $s^{p} \in C^{p}(\mathfrak{U}, S(\{U\}))$ is a collection

$$
s^{p}=\left\{s^{p}\left(i_{0}, \ldots, i_{p}\right)\right\}
$$

of local sections $s^{p}\left(i_{0}, \ldots, i_{p}\right)$ of the sheaf $S$ over $U_{i_{0}} \cap \cdots \cap U_{i_{p}}$ for each $(p+1)$-tuple $\left(U_{i_{0}}, \ldots, U_{i_{p}}\right)$ of elements of the cover $\mathfrak{U}$. The coboundary operator (8.5.7) reads

$$
\delta^{p} s^{p}\left(i_{0}, \ldots, i_{p+1}\right)=\left.\sum_{k=0}^{p+1}(-1)^{k} s^{p}\left(i_{0}, \ldots, \widehat{i}_{k}, \ldots, i_{p+1}\right)\right|_{U_{i_{0}} \cap \cdots \cap U_{i_{p+1}}}
$$

For instance, we have

$$
\begin{align*}
& \delta^{0} s^{0}(i, j)=\left.\left[s^{0}(j)-s^{0}(i)\right]\right|_{U_{i} \cap U_{j}}  \tag{8.5.10}\\
& \delta^{1} s^{1}(i, j, k)=\left.\left[s^{1}(j, k)-s^{1}(i, k)+s^{1}(i, j)\right]\right|_{U_{i} \cap U_{j} \cap U_{k}} \tag{8.5.11}
\end{align*}
$$

A glance at the expression (8.5.10) shows that a zero-cocycle is a collection $s=\{s(i)\}_{I}$ of local sections of the sheaf $S$ over $U_{i} \in \mathfrak{U}$ such that $s(i)=s(j)$ on $U_{i} \cap U_{j}$. It follows from the axiom (S2) in Remark 8.5.2 that $s$ is a global section of the sheaf $S$, while each $s(i)$ is its restriction $\left.s\right|_{U_{i}}$ to $U_{i}$. Consequently, the cohomology group $H^{0}(\mathfrak{U} ; S(\{U\}))$ is isomorphic to the structure module $S(X)$ of global sections of the sheaf $S$. A one-cocycle is a collection $\{s(i, j)\}$ of local sections of the sheaf $S$ over overlaps $U_{i} \cap U_{j}$ which satisfy the cocycle condition

$$
\begin{equation*}
\left.[s(j, k)-s(i, k)+s(i, j)]\right|_{U_{i} \cap U_{j} \cap U_{k}}=0 \tag{8.5.12}
\end{equation*}
$$

If $X$ is a paracompact space, the study of its sheaf cohomology is essentially simplified due to the following fact.

Theorem 8.5.2: Cohomology of a paracompact space $X$ with coefficients in a sheaf $S$ coincides with cohomology of $X$ with coefficients in any presheaf generating the sheaf $S$.

Remark 8.5.4: We follow the definition of a paracompact topological space as a Hausdorff space such that any its open cover admits a locally finite open refinement, i.e., any point has an open neighborhood which intersects only a finite number of elements of this refinement. A topological space $X$ is paracompact iff any cover $\left\{U_{\xi}\right\}$ of $X$ admits a subordinate partition of unity $\left\{f_{\xi}\right\}$, i.e.:
(i) $f_{\xi}$ are real positive continuous functions on $X$;
(ii) $\operatorname{supp} f_{\xi} \subset U_{\xi}$;
(iii) each point $x \in X$ has an open neighborhood which intersects only a finite number of the sets $\operatorname{supp} f_{\xi}$;
(iv) $\sum_{\xi} f_{\xi}(x)=1$ for all $x \in X$.

The key point of the analysis of sheaf cohomology is that short exact sequences of sheaves yield long exact sequences of their cohomology groups.

Let $S_{\{U\}}$ and $S_{\{U\}}^{\prime}$ be presheaves on the same topological space $X$. It is readily observed that, given an open cover $\mathfrak{U}$ of $X$, any morphism $S_{\{U\}} \rightarrow S_{\{U\}}^{\prime}$ yields a cochain morphism of complexes

$$
C^{*}\left(\mathfrak{U}, S_{\{U\}}\right) \rightarrow C^{*}\left(\mathfrak{U}, S_{\{U\}}^{\prime}\right)
$$

and the corresponding morphism

$$
H^{*}\left(\mathfrak{U} ; S_{\{U\}}\right) \rightarrow H^{*}\left(\mathfrak{U} ; S_{\{U\}}^{\prime}\right)
$$

of cohomology groups of these complexes. Passing to the direct limit through all refinements of $\mathfrak{U}$, we come to a morphism of cohomology groups

$$
H^{*}\left(X ; S_{\{U\}}\right) \rightarrow H^{*}\left(X ; S_{\{U\}}^{\prime}\right)
$$

of $X$ with coefficients in the presheaves $S_{\{U\}}$ and $S_{\{U\}}^{\prime}$. In particular, any sheaf morphism $S \rightarrow S^{\prime}$ yields a morphism of canonical presheaves

$$
S(\{U\}) \rightarrow S^{\prime}(\{U\})
$$

and the corresponding cohomology morphism

$$
H^{*}(X ; S) \rightarrow H^{*}\left(X ; S^{\prime}\right)
$$

By virtue of Theorems 8.3.1 and 8.3.2, every short exact sequence

$$
\begin{equation*}
0 \rightarrow S_{\{U\}}^{\prime} \longrightarrow S_{\{U\}} \longrightarrow S_{\{U\}}^{\prime \prime} \rightarrow 0 \tag{8.5.13}
\end{equation*}
$$

of presheaves on the same topological space $X$ and the corresponding exact sequence of complexes (8.5.8) yield the long exact sequence

$$
\begin{align*}
& 0 \rightarrow H^{0}\left(X ; S_{\{U\}}^{\prime}\right) \longrightarrow H^{0}\left(X ; S_{\{U\}}\right) \longrightarrow H^{0}\left(X ; S_{\{U\}}^{\prime \prime}\right) \longrightarrow(8  \tag{8.5.14}\\
& H^{1}\left(X ; S_{\{U\}}^{\prime}\right) \longrightarrow \cdots H^{p}\left(X ; S_{\{U\}}^{\prime}\right) \longrightarrow H^{p}\left(X ; S_{\{U\}}\right) \longrightarrow \\
& H^{p}\left(X ; S_{\{U\}}^{\prime \prime}\right) \longrightarrow H^{p+1}\left(X ; S_{\{U\}}^{\prime}\right) \longrightarrow \cdots
\end{align*}
$$

of the cohomology groups of $X$ with coefficients in these presheaves. This result is extended to the exact sequence of sheaves. Let

$$
\begin{equation*}
0 \rightarrow S^{\prime} \longrightarrow S \longrightarrow S^{\prime \prime} \rightarrow 0 \tag{8.5.15}
\end{equation*}
$$

be a short exact sequence of sheaves on $X$. It yields the short exact sequence of presheaves (8.5.5) where the presheaf $S_{\{U\}}^{\prime \prime}$ generates the sheaf $S^{\prime \prime}$. If $X$ is paracompact,

$$
H^{*}\left(X ; S_{\{U\}}^{\prime \prime}\right)=H^{*}\left(X ; S^{\prime \prime}\right)
$$

in accordance with Theorem 8.5.2, and we have the exact sequence of sheaf cohomology

$$
\begin{align*}
0 \rightarrow & H^{0}\left(X ; S^{\prime}\right)  \tag{8.5.16}\\
H^{1}\left(X ; S^{\prime}\right) & \left.\longrightarrow H^{0}(X ; S) \longrightarrow H^{p}\left(X ; S^{\prime}\right) \longrightarrow S^{\prime \prime}\right) \longrightarrow \\
& \left.H^{p}\left(X ; S^{\prime \prime}\right) \longrightarrow H^{p+1}(X ; S) \longrightarrow S^{\prime}\right) \longrightarrow \cdots .
\end{align*}
$$

Let us turn now to the abstract de Rham theorem which provides a powerful tool of studying algebraic systems on paracompact spaces.

Let us consider an exact sequence of sheaves

$$
\begin{equation*}
0 \rightarrow S \xrightarrow{h} S_{0} \xrightarrow{h^{0}} S_{1} \xrightarrow{h^{1}} \cdots S_{p} \xrightarrow{h^{p}} \cdots . \tag{8.5.17}
\end{equation*}
$$

It is said to be a resolution of the sheaf $S$ if each sheaf $S_{p \geq 0}$ is acyclic, i.e., its cohomology groups $H^{k>0}\left(X ; S_{p}\right)$ vanish.

Any exact sequence of sheaves (8.5.17) yields the sequence of their structure modules

$$
\begin{equation*}
0 \rightarrow S(X) \xrightarrow{h_{*}} S_{0}(X) \xrightarrow{h_{*}^{0}} S_{1}(X) \xrightarrow{h_{*}^{1}} \cdots S_{p}(X) \xrightarrow{h_{*}^{p}} \cdots \tag{8.5.18}
\end{equation*}
$$

which is always exact at terms $S(X)$ and $S_{0}(X)$ (see the exact sequence (8.5.4)). The sequence (8.5.18) is a cochain complex because

$$
h_{*}^{p+1} \circ h_{*}^{p}=0 .
$$

If $X$ is a paracompact space and the exact sequence (8.5.17) is a resolution of $S$, the forthcoming abstract de Rham theorem establishes an isomorphism of cohomology of the complex (8.5.18) to cohomology of $X$ with coefficients in the sheaf $S$.

Theorem 8.5.3: Given a resolution (8.5.17) of a sheaf $S$ on a paracompact topological space $X$ and the induced complex (8.5.18), there are isomorphisms

$$
H^{0}(X ; S)=\operatorname{Ker} h_{*}^{0}, \quad H^{q}(X ; S)=\operatorname{Ker} h_{*}^{q} / \operatorname{Im} h_{*}^{q-1}, \quad q>0
$$

A sheaf $S$ on a paracompact space $X$ is called fine if, for each locally finite open cover $\mathfrak{U}=\left\{U_{i}\right\}_{i \in I}$ of $X$, there exists a system $\left\{h_{i}\right\}$ of endomorphisms $h_{i}: S \rightarrow S$ such that:
(i) there is a closed subset $V_{i} \subset U_{i}$ and $h_{i}\left(S_{x}\right)=0$ if $x \notin V_{i}$,
(ii) $\sum_{i \in I} h_{i}$ is the identity map of $S$.

Theorem 8.5.4: A fine sheaf on a paracompact space is acyclic.
There is the following important example of fine sheaves.
Theorem 8.5.5: Let $X$ be a paracompact topological space which admits a partition of unity performed by elements of the structure module $\mathfrak{A}(X)$ of some sheaf $\mathfrak{A}$ of real functions on $X$. Then any sheaf $S$ of $\mathfrak{A}$-modules on $X$, including $\mathfrak{A}$ itself, is fine.

In particular, the sheaf $C_{X}^{0}$ of continuous functions on a paracompact topological space is fine, and so is any sheaf of $C_{X}^{0}$-modules.

### 8.6 Local-ringed spaces

Local-ringed spaces are sheafs of local rings. For instance, smooth manifolds, represented by sheaves of real smooth functions, make up a subcategory of the category of local-ringed spaces.

A sheaf $\Re$ on a topological space $X$ is said to be a ringed space if its stalk $\Re_{x}$ at each point $x \in X$ is a real commutative ring. A ringed space is often denoted by a pair $(X, \mathfrak{R})$ of a topological space $X$ and a sheaf $\mathfrak{R}$ of rings on $X$. They are called the body and the structure sheaf of a ringed space, respectively.

A ringed space is said to be a local-ringed space (a geometric space) if it is a sheaf of local rings.

For instance, the sheaf $C_{X}^{0}$ of continuous real functions on a topological space $X$ is a local-ringed space. Its stalk $C_{x}^{0}, x \in X$, contains the unique maximal ideal of germs of functions vanishing at $x$.

Morphisms of local-ringed spaces are defined as those of sheaves on different topological spaces as follows.

Let $\varphi: X \rightarrow X^{\prime}$ be a continuous map. Given a sheaf $S$ on $X$, its direct image $\varphi_{*} S$ on $X^{\prime}$ is generated by the presheaf of assignments

$$
X^{\prime} \supset U^{\prime} \rightarrow S\left(\varphi^{-1}\left(U^{\prime}\right)\right)
$$

for any open subset $U^{\prime} \subset X^{\prime}$. Conversely, given a sheaf $S^{\prime}$ on $X^{\prime}$, its inverse image $\varphi^{*} S^{\prime}$ on $X$ is defined as the pull-back onto $X$ of the continuous fibre bundle $S^{\prime}$ over $X^{\prime}$, i.e., $\varphi^{*} S_{x}^{\prime}=S_{\varphi(x)}$. This sheaf is generated by the presheaf which associates to any open $V \subset X$ the direct limit of modules $S^{\prime}(U)$ over all open subsets $U \subset X^{\prime}$ such that $V \subset f^{-1}(U)$.

Remark 8.6.1: Let $i: X \rightarrow X^{\prime}$ be a closed subspace of $X^{\prime}$. Then $i_{*} S$ is a unique sheaf on $X^{\prime}$ such that

$$
\left.i_{*} S\right|_{X}=S,\left.\quad i_{*} S\right|_{X^{\prime} \backslash X}=0
$$

Indeed, if $x^{\prime} \in X \subset X^{\prime}$, then

$$
i_{*} S\left(U^{\prime}\right)=S\left(U^{\prime} \cap X\right)
$$

for any open neighborhood $U$ of this point. If $x^{\prime} \notin X$, there exists its neighborhood $U^{\prime}$ such that $U^{\prime} \cap X$ is empty, i.e., $i_{*} S\left(U^{\prime}\right)=0$. The sheaf $i_{*} S$ is called the trivial extension of the sheaf $S$.

By a morphism of ringed spaces

$$
(X, \mathfrak{R}) \rightarrow\left(X^{\prime}, \mathfrak{R}^{\prime}\right)
$$

is meant a pair $(\varphi, \hat{\varphi})$ of a continuous map $\varphi: X \rightarrow X^{\prime}$ and a sheaf morphism $\hat{\varphi}: \mathfrak{R}^{\prime} \rightarrow \varphi_{*} \Re$ or, equivalently, a sheaf morphisms $\varphi^{*} \Re^{\prime} \rightarrow \mathfrak{R}$. Restricted to each stalk, a sheaf morphism $\Phi$ is assumed to be a ring homomorphism. A morphism of ringed spaces is said to be:

- a monomorphism if $\varphi$ is an injection and $\Phi$ is an epimorphism,
- an epimorphism if $\varphi$ is a surjection, while $\Phi$ is a monomorphism.

Let $(X, \mathfrak{R})$ be a local-ringed space. By a sheaf $\mathfrak{O} \mathfrak{R}$ of derivations of the sheaf $\mathfrak{R}$ is meant a subsheaf of endomorphisms of $\mathfrak{R}$ such that any section $u$ of $\mathfrak{d} \mathfrak{R}$ over an open subset $U \subset X$ is a derivation of the real ring $\mathfrak{R}(U)$. It should be emphasized that, since the monomorphism (8.5.1) is not necessarily an isomorphism, a derivation of the ring $\mathfrak{R}(U)$ need not be a section of the sheaf $\left.\mathfrak{d} \mathfrak{R}\right|_{U}$. Namely, it may happen that, given open sets $U^{\prime} \subset U$, there is no restriction morphism

$$
\mathfrak{d}(\mathfrak{R}(U)) \rightarrow \mathfrak{d}\left(\mathfrak{R}\left(U^{\prime}\right)\right) .
$$

Given a local-ringed space $(X, \mathfrak{R})$, a sheaf $P$ on $X$ is called a sheaf of $\mathfrak{R}$-modules if every stalk $P_{x}, x \in X$, is an $\Re_{x}$-module or, equivalently, if $P(U)$ is an $\mathfrak{R}(U)$-module for any open subset $U \subset X$. A sheaf of $\mathfrak{R}$ modules $P$ is said to be locally free if there exists an open neighborhood $U$ of every point $x \in X$ such that $P(U)$ is a free $\mathfrak{R}(U)$-module. If all these free modules are of finite rank (resp. of the same finite rank), one says that $P$ is of finite type (resp. of constant rank). The structure module of a locally free sheaf is called a locally free module.

The following is a generalization of Theorem 8.5.5.

Theorem 8.6.1: Let $X$ be a paracompact space which admits a partition of unity by elements of the structure module $S(X)$ of some sheaf $S$ of real functions on $X$. Let $P$ be a sheaf of $S$-modules. Then $P$ is fine and, consequently, acyclic.

Assumed to be paracompact, a smooth manifold $X$ admits a partition of unity performed by smooth real functions. It follows that the sheaf $C_{X}^{\infty}$ of smooth real functions on $X$ is fine, and so is any sheaf of $C_{X^{-}}^{\infty}$ modules, e.g., the sheaves of sections of smooth vector bundles over $X$.

Similarly to the sheaf $C_{X}^{0}$ of continuous functions, the sheaf $C_{X}^{\infty}$ of smooth real functions on a smooth manifold $X$ is a local-ringed spaces. Its stalk $C_{x}^{\infty}$ at a point $x \in X$ has a unique maximal ideal $\mu_{x}$ of germs of smooth functions vanishing at $x$. Though the sheaf $C_{X}^{\infty}$ is defined on a topological space $X$, it fixes a unique smooth manifold structure on $X$ as follows.

Theorem 8.6.2: Let $X$ be a paracompact topological space and ( $X, \mathfrak{R}$ ) a local-ringed space. Let $X$ admit an open cover $\left\{U_{i}\right\}$ such that the sheaf $\mathfrak{R}$ restricted to each $U_{i}$ is isomorphic to the local-ringed space $\left(\mathbb{R}^{n}, C_{R^{n}}^{\infty}\right)$. Then $X$ is an $n$-dimensional smooth manifold together with a natural isomorphism of local-ringed spaces $(X, \mathfrak{R})$ and $\left(X, C_{X}^{\infty}\right)$.

One can think of this result as being an alternative definition of smooth real manifolds in terms of local-ringed spaces. In particular, there is one-to-one correspondence between smooth manifold morphisms $X \rightarrow X^{\prime}$ and the $\mathbb{R}$-ring morphisms $C^{\infty}\left(X^{\prime}\right) \rightarrow C^{\infty}(X)$.

For instance, let $Y \rightarrow X$ be a smooth vector bundle. The germs of its sections make up a sheaf of $C_{X}^{\infty}$-modules, called the structure sheaf $S_{Y}$ of a vector bundle $Y \rightarrow X$. The sheaf $S_{Y}$ is fine. The structure module of this sheaf coincides with the structure module $Y(X)$ of global sections of a vector bundle $Y \rightarrow X$. The following Serre-Swan theorem shows that these modules exhaust all projective modules of finite rank over $C^{\infty}(X)$. Originally proved for bundles over a compact base $X$, this
theorem has been extended to an arbitrary $X$.
Theorem 8.6.3: Let $X$ be a smooth manifold. A $C^{\infty}(X)$-module $P$ is isomorphic to the structure module of a smooth vector bundle over $X$ iff it is a projective module of finite rank.

This theorem states the categorial equivalence between the vector bundles over a smooth manifold $X$ and projective modules of finite rank over the ring $C^{\infty}(X)$ of smooth real functions on $X$. The following are corollaries of this equivalence

- The structure module $Y^{*}(X)$ of the dual $Y^{*} \rightarrow X$ of a vector bundle $Y \rightarrow X$ is the $C^{\infty}(X)$-dual $Y(X)^{*}$ of the structure module $Y(X)$ of $Y \rightarrow X$.
- Any exact sequence of vector bundles

$$
\begin{equation*}
0 \rightarrow Y \longrightarrow Y^{\prime} \longrightarrow Y^{\prime \prime} \rightarrow 0 \tag{8.6.1}
\end{equation*}
$$

over the same base $X$ yields the exact sequence

$$
\begin{equation*}
0 \rightarrow Y(X) \longrightarrow Y^{\prime}(X) \longrightarrow Y^{\prime \prime}(X) \rightarrow 0 \tag{8.6.2}
\end{equation*}
$$

of their structure modules, and vice versa. In accordance with Theorem 1.2.2, the exact sequence (8.6.1) is always split. Every its splitting defines that of the exact sequence (8.6.2), and vice versa.

- The derivation module of the real ring $C^{\infty}(X)$ coincides with the $C^{\infty}(X)$-module $\mathcal{T}(X)$ of vector fields on $X$, i.e., with the structure module of the tangent bundle $T X$ of $X$. Hence, it is a projective $C^{\infty}(X)$ module of finite rank. It is the $C^{\infty}(X)$-dual $\mathcal{T}(X)=\mathcal{O}^{1}(X)^{*}$ of the structure module $\mathcal{O}^{1}(X)$ of the cotangent bundle $T^{*} X$ of $X$ which is the module of differential one-forms on $X$ and, conversely,

$$
\mathcal{O}^{1}(X)=\mathcal{T}(X)^{*}
$$

- Therefore, if $P$ is a $C^{\infty}(X)$-module, one can reformulate Definition 8.2.3 of a connection on $P$ as follows. A connection on a $C^{\infty}(X)$-module $P$ is a $C^{\infty}(X)$-module morphism

$$
\begin{equation*}
\nabla: P \rightarrow \mathcal{O}^{1}(X) \otimes P \tag{8.6.3}
\end{equation*}
$$

which satisfies the Leibniz rule

$$
\nabla(f p)=d f \otimes p+f \nabla(p), \quad f \in C^{\infty}(X), \quad p \in P
$$

It associates to any vector field $\tau \in \mathcal{T}(X)$ on $X$ a first order differential operator $\nabla_{\tau}$ on $P$ which obeys the Leibniz rule

$$
\begin{equation*}
\left.\nabla_{\tau}(f p)=(\tau\rfloor d f\right) p+f \nabla_{\tau} p . \tag{8.6.4}
\end{equation*}
$$

In particular, let $Y \rightarrow X$ be a vector bundle and $Y(X)$ its structure module. The notion of a connection on the structure module $Y(X)$ is equivalent to the standard geometric notion of a connection on a vector bundle $Y \rightarrow X$.

Since the derivation module of the real ring $C^{\infty}(X)$ is the $C^{\infty}(X)$ module $\mathcal{T}(X)$ of vector fields on $X$ and

$$
\mathcal{O}^{1}(X)=\mathcal{T}(X)^{*},
$$

the Chevalley-Eilenberg differential calculus over the real ring $C^{\infty}(X)$ is exactly the DGA $\left(\mathcal{O}^{*}(X), d\right)$ of exterior forms on $X$, where the ChevalleyEilenberg coboundary operator $d$ (8.4.8) coincides with the exterior differential. Moreover, one can show that $\left(\mathcal{O}^{*}(X), d\right)$ is a minimal differential calculus, i.e., the $C^{\infty}(X)$-module $\mathcal{O}^{1}(X)$ is generated by elements $d f, f \in C^{\infty}(X)$. Therefore, the de Rham complex (8.4.11) of the real ring $C^{\infty}(X)$ is the de Rham complex

$$
\begin{equation*}
0 \rightarrow \mathbb{R} \longrightarrow C^{\infty}(X) \xrightarrow{d} \mathcal{O}^{1}(X) \xrightarrow{d} \cdots \mathcal{O}^{k}(X) \xrightarrow{d} \cdots \tag{8.6.5}
\end{equation*}
$$

of exterior forms on a manifold $X$.
The de Rham cohomology of the complex (8.6.5) is called the de Rham cohomology $H_{\mathrm{DR}}^{*}(X)$ of $X$. To describe them, let us consider the de Rham complex

$$
\begin{equation*}
0 \rightarrow \mathbb{R} \longrightarrow C_{X}^{\infty} \xrightarrow{d} \mathcal{O}_{X}^{1} \xrightarrow{d} \cdots \mathcal{O}_{X}^{k} \xrightarrow{d} \cdots \tag{8.6.6}
\end{equation*}
$$

of sheaves $\mathcal{O}_{X}^{k}, k \in \mathbb{N}_{+}$, of germs of exterior forms on $X$. These sheaves are fine. Due to the Poincaré lemma, the complex (8.6.6) is exact and,
thereby, is a fine resolution of the constant sheaf $\mathbb{R}$ on a manifold $X$. Then a corollary of Theorem 8.5.3 is the classical de Rham theorem.

Theorem 8.6.4: There is an isomorphism

$$
\begin{equation*}
H_{\mathrm{DR}}^{k}(X)=H^{k}(X ; \mathbb{R}) \tag{8.6.7}
\end{equation*}
$$

of the de Rham cohomology $H_{\mathrm{DR}}^{*}(X)$ of a manifold $X$ to cohomology of $X$ with coefficients in the constant sheaf $\mathbb{R}$.

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## Index

$B^{n, \mu} 117$
$C_{X}^{\infty}, 179$
$C_{\mathrm{W}}, 103$
$D_{\Gamma}, 56$
$G$-bundle, 74
principal, 77
smooth, 76
$G L_{n}, 101$
HY, 52
$H^{1}\left(X ; G_{X}^{0}\right), 74$
$J^{1} Y, 33$
$J^{1} \Phi, 35$
$J_{\Sigma}^{1} Y, 64$
$J^{1} s, 34$
$J^{1} u, 35$
$J^{\infty} Y, 45$
$J^{k} u, 39$
$J^{r} Y, 36$
$J^{r} \Phi, 38$
$J^{r} s, 38$
LX, 101
$L_{G}, 70$
$L_{g}, 70$
$P^{G}, 89$
$P^{\otimes k}, 168$
$P_{\Sigma}, 96$
$R_{G}, 70$
$R_{g}, 70$
$R_{G P}, 77$
$R_{g P}, 77$
$T Z, 7$
$T^{*} Z, 7$
$T_{G} P, 79$
Tf, 7
VY, 19
VГ, 67
$V \Phi, 19$
$V^{*} Y, 19$
$V^{*} \Gamma, 67$
$V_{G} P, 79$
$V_{\Sigma} Y, 65$
$V_{\Sigma}^{*} Y, 65$
$Y(X), 16$
$Y^{h}, 15$
[A], 145
$\Gamma \tau, 52$
$\Gamma^{*}, 59$
$\mathbf{L}_{u}, 27$
1, 69
u, 160
$\mathcal{A}_{E}, 123$
$\mathcal{E}_{L}, 133$
$\mathcal{E}_{i}, 134$
$\mathfrak{o} \mathcal{A}_{E}, 126$
$\mathcal{F}, 87$
$\mathfrak{g} l, 70$
$\mathfrak{g}_{r}, 70$
$\mathcal{J}_{v}, 139$
$\mathcal{O}^{*}(Z), 25$
$\mathcal{O}^{*}[\mathfrak{d} \mathcal{A}], 177$
$\mathcal{O}^{*} \mathcal{A}, 178$
$\mathcal{O}_{\infty}^{*} Y, 46$
$\mathcal{O}_{\infty}^{*}, 46$
$\mathcal{O}_{\infty}^{k, m}, 48$
$\mathcal{O}_{k}^{*}, 38$
$\mathcal{P}_{\infty}^{*}, 136$
$\mathcal{P}_{\infty}^{*}[E ; Y], 153$
$\mathcal{Q}_{\infty}^{*}, 47$
$\mathcal{S}^{*}[E ; Z], 128$
$\mathcal{S}_{\infty}^{*}[F ; X], 143$
$\mathcal{S}_{\infty}^{*}[F ; Y], 144$
$\mathcal{S}_{r}^{*}[F ; Y], 143$
$\mathcal{T}(Z), 22$
$\mathcal{V}_{E}, 126$
$\mathcal{G}(X), 89$
$\mathcal{G}_{L}, 150$
$\delta_{N}, 158$
$\delta_{\mathrm{KT}}, 157$
$\dot{\partial}_{\lambda}, 23$
$\epsilon_{m}, 71$
$\mathfrak{A}(Z), 122$
$\mathfrak{A}_{E}, 123$
$\mathfrak{o} \mathcal{A}, 171$
$\nabla^{\Gamma}, 56$
$\bar{E}, 153$
$\bar{d} y^{i}, 20$
$\bar{\delta}, 154$
$\omega, 27$
$\omega_{\lambda}, 27$
$\overleftarrow{v}, 151$
$\otimes P, 168$
$\pi^{1}, 34$
$\pi_{0}^{1}, 34$
$\pi_{r}^{\infty}, 45$
$\pi_{k}^{r}, 36$
$\pi_{Y \Sigma}, 15$
$\pi_{\Sigma X}, 15$
$\psi_{\xi}, 10$
$\theta^{i}, 34$
$\theta_{\Lambda}^{i}, 40$
$\theta_{Z}, 28$
$\theta_{L X}, 102$
$\varepsilon_{m}, 71$
$\varrho_{\xi \zeta}, 10$
$\wedge Y, 17$
$\widehat{0}, 16$
$\widetilde{D}, 66$
$\widetilde{\tau}, 23$
$\left\{\lambda^{\nu}{ }_{\mu}\right\}, 61$
$d_{H}, 48$
$d_{V}, 48$
$d_{\lambda}, 34$
$f^{*} Y, 14$
$f^{*} \Gamma, 53$
$f^{*} \phi, 26$
$h_{0}, 41$
$u\rfloor \phi, 27$
$u_{\xi}, 90$
$\mathrm{Ad}_{g}, 72$
action of a group, 69
effective, 70
free, 70
on the left, 69
on the right, 69
transitive, 70
action of a structure group
on $J^{1} P, 81$
on $P, 77$
on $T P, 79$
adjoint representation
of a Lie algebra, 72
of a Lie group, 72
affine bundle, 20
morphism, 21
algebra, 163
$\mathbb{Z}_{2}$-graded, 114
commutative, 114
$\mathcal{O}_{\infty}^{*} Y, 46$
$\mathcal{O}_{\infty}^{*}, 46$
$\mathcal{P}_{\infty}^{*}[E ; Y], 153$
$\mathcal{S}_{\infty}^{*}[F ; Y], 144$
differential bigraded, 120
differential graded, 176
graded, 176
commutative, 176
unital, 163
annihilator of a distribution, 23
antifield, 154
$k$-stage, 157
Noether, 155
associated bundles, 75
automorphism
associated, 93
principal, 88
autoparallel, 106
base of a fibred manifold, 9
basic form, 27
basis
for a graded manifold, 124
for a module, 165
generating, 145
Batchelor theorem, 123
Bianchi identity
first, 58
second, 57
bigraded
de Rham complex, 121
exterior algebra, 115
bimodule, 164
commutative, 164
graded, 114
body
of a graded manifold, 122
of a ringed space, 188
body map, 116
boundary, 173
boundary operator, 173
bundle
$P$-associated, 92
affine, 20
associated, 95
canonically, 95
atlas, 11
associated, 92
holonomic, 18
of constant local trivializations, 63
automorphism, 13
vertical, 13
composite, 15
continuous, 11
locally trivial, 11
coordinates, 12
affine, 20
linear, 16
cotangent, 7
density-dual, 153
epimorphism, 13
exterior, 17
gauge natural, 103
isomorphism, 13
lift, 95
locally trivial, 11
monomorphism, 13
morphism, 13
affine, 21
linear, 16
of principal bundles, 78
natural, 100
of principal connections, 83
of world connections, 103
principal, 77
product, 14
smooth, 11
tangent, 7
affine, 21
vertical, 19
with a structure group, 74
principal, 77
smooth, 76
canonical principal connection, 87
Cartan connection, 62
chain, 173
Chevalley-Eilenberg
coboundary operator, 176
graded, 120
cohomology, 176
complex, 176
differential calculus, 178
Grassmann-graded, 120
minimal, 178
Christoffel symbols, 61
classical solution, 42
closed map, 10
coboundary, 173
coboundary operator, 173
Chevalley-Eilenberg, 176
cochain, 173
cochain morphism, 174
cocycle, 173
cocycle condition, 10
for a sheaf, 184
codistribution, 23
coframe, 7
cohomology, 173
Chevalley-Eilenberg, 176
de Rham
abstract, 177
of a manifold, 192
with coefficients in a sheaf, 183
complex, 173
$k$-exact, 173
acyclic, 173
chain, 172
Chevalley-Eilenberg, 176
cochain, 173
de Rham
abstract, 177
exact, 173
variational, 132
component of a connection, 51
composite bundle, 15
composite connection, 65
connection, 51
affine, 61
composite, 65
covertical, 68
dual, 59
flat, 63
linear, 58
world, 103
on a graded commutative ring, 122
on a graded manifold, 127
on a graded module, 121
on a module, 172
on a ring, 172
principal, 83
associated, 94
canonical, 87
projectable, 65
reducible, 57
vertical, 67
connection form, 52
of a principal connection, 84
local, 84
vertical, 53
conservation law
weak, 138
contact derivation, 136
graded, 149
projectable, 137
vertical, 137
contact form, 41
graded, 145
local, 34
of higher jet order, 40
contraction, 27
cotangent bundle, 7
vertical, 19
covariant derivative, 56
covariant differential, 56
on a module, 172
vertical, 66
curvature, 57
of a principal connection, 85
of a world connection, 60
of an associated principal connection, 94
soldered, 58
curvature-free connection, 63
curve integral, 22
cycle, 173

DBGA, 120
de Rham cohomology
abstract, 177
of a graded manifold, 129
of a manifold, 192
of a ring, 178
de Rham complex
abstract, 177
bigraded, 121
of a ring, 178
of exterior forms, 192
of sheaves, 192
de Rham theorem, 193
abstract, 187
density, 27
density-dual bundle, 153
density-dual vector bundle, 153
graded, 153
derivation, 171
contact, 136
projectable, 137
vertical, 137
graded, 118
contact, 149
nilpotent, 151
right, 151
derivation module, 171
graded, 119
DGA, 176
differential
covariant, 56
vertical, 66
exterior, 26
total, 48
graded, 146
vertical, 48
graded, 146
differential calculus, 176
Chevalley-Eilenberg, 178
minimal, 178
minimal, 177
differential equation, 42
associated to a differential operator, 43
differential form, 46
graded, 145
differential ideal, 24
differential operator
as a morphism, 43
as a section, 42
graded, 118
of standard form, 43
on a module, 170
direct image of a sheaf, 188
direct limit, 167
direct sequence, 167
direct sum
of complexes, 174
of modules, 164
direct system of modules, 167
directed set, 167
distribution, 23
horizontal, 52
involutive, 23
domain, 12
dual module, 165
dual vector bundle, 17

Ehresmann connection, 54
equivalent $G$-bundle atlases, 74
equivalent $G$-bundles, 74
equivalent bundle atlases, 11
equivariant
automorphism, 88
connection, 83
function, 88
Euler-Lagrange operator, 133
Euler-Lagrange-type operator, 134
even element, 114
even morphism, 115
exact sequence
of modules, 166
short, 166
split, 166
of vector bundles, 18
short, 18
split, 18
exterior algebra, 169
bigraded, 115
exterior bundle, 17
exterior differential, 26
exterior form, 25
basic, 27
graded, 129
horizontal, 27
exterior product, 26
graded, 115
of vector bundles, 17
factor
algebra, 163
bundle, 18
complex, 174
module, 164
sheaf, 181
fibration, 9
fibre, 9
fibre bundle, 10
fibre coordinates, 12
fibred coordinates, 9
fibred manifold, 9
fibrewise morphism, 13
first Noether theorem, 138
for a graded Lagrangian, 151
first variational formula, 137
for a graded Lagrangian, 149
flow, 22
foliated manifold, 25
foliation, 24
horizontal, 63
simple, 25
Frölicher-Nijenhuis bracket, 29
frame, 16
holonomic, 7
frame field, 101
gauge
algebra, 81
algebra bundle, 81
group, 89
operator, 162
parameters, 140
transformation, 88
infinitesimal, 89
gauge natural bundle, 103
gauge symmetry, 140
$k$-stage, 162
first-stage, 162
of a graded Lagrangian, 161
reducible, 142
general covariant transformation, 100
infinitesimal, 100
generalized vector field, 136
graded, 149
generating basis, 145
geodesic, 107
geodesic equation, 107
geometric space, 188
ghost, 160
graded
algebra, 176
bimodule, 114
commutative algebra, 176
commutative ring, 116
Banach, 116
real, 116
connection, 127
derivation, 118
of a field system algebra, 151
derivation module, 119
differential form, 145
differential operator, 118
exterior differential, 129
exterior form, 129
exterior product, 115
function, 122
interior product, 121
Leibniz rule, 177
manifold, 122
composite, 143
simple, 123
module, 114
free, 114
dual, 115
morphism, 115
even, 115
odd, 115
ring, 114
vector field, 125
generalized, 149
vector space, 114
( $n, m$ )-dimensional, 114
graded-homogeneous element, 114
grading automorphism, 114
Grassmann algebra, 116
group bundle, 88

Heimholtz condition, 134
holonomic
atlas, 18
of the frame bundle, 101
automorphisms, 102
coframe, 7
coordinates, 7
frame, 7
homogeneous space, 70
homology, 173
homology regularity condition, 158
horizontal
distribution, 52
foliation, 63
form, 27
graded, 145
lift
of a path, 54
of a vector field, 52
projection, 41
splitting, 52
canonical, 54
vector field, 52
ideal, 163
maximal, 163
of nilpotents, 116
proper, 163
image of a sheaf
direct, 188
inverse, 188
imbedded submanifold, 8
imbedding, 8
immersion, 8
induced coordinates, 18
inductive limit, 169
infinitesimal generator, 22
infinitesimal transformation
of a Lagrangian system, 136
Grassmann-graded, 148
integral curve, 22
integral manifold, 24
maximal, 24
integral section of a connection, 56
interior product, 27
graded, 121
of vector bundles, 17
inverse image of a sheaf, 188
inverse sequence, 169
jet
first order, 33
higher order, 36
infinite order, 45
jet bundle, 34
affine, 34
jet coordinates, 34
jet manifold, 34
higher order, 36
infinite order, 45
jet prolongation
functor, 37
of a morphism, 35
higher order, 38
of a section, 34
higher order, 38
of a structure group action, 81 of a vector field, 35
higher order, 39
juxtaposition rule, 177
kernel
of a bundle morphism, 13
of a differential operator, 43
of a vector bundle morphism, 16

Koszul-Tate complex, 158
Koszul-Tate operator, 158
Lagrangian, 133
degenerate, 154
graded, 147
variationally trivial, 134
Lagrangian system, 134
$N$-stage reducible, 157
finitely degenerate, 155
Grassmann-graded, 147
irreducible, 156
reducible, 156
leaf, 24
left-invariant form, 73
canonical, 73
Leibniz rule, 171
for a connection, 172
Grassmann-graded, 121
graded, 177
Grassmann-graded, 118
Lepage equivalent
of a graded Lagrangian, 148
Levi-Civita connection, 61
Levi-Civita symbol, 27
Lie algebra
left, 70
right, 70
Lie algebra bundle, 79
Lie bracket, 22
Lie derivative
graded, 121
of a tangent-valued form, 29
of an exterior form, 27
Lie superalgebra, 117
Lie superbracket, 117
lift of a bundle, 95
lift of a vector field
canonical, 23
functorial, 100
horizontal, 52
linear derivative of an affine morphism, 21
linear frame bundle, 101
local diffeomerphism, 8
local ring, 163
local-ringed space, 188
locally finite cover, 184
manifold, 5
fibred, 9
flat, 106
parallelizable, 106
matrix group, 72
module, 164
dual, 165
finitely generated, 166
free, 165
graded, 114
locally free, 189
of finite rank, 166
over a Lie algebra, 175
over a Lie superalgebra, 117
projective, 166
morphism
of fibre bundles, 13
of graded manifolds, 125
of presheaves, 181
of ringed spaces, 189
of sheaves, 181
Mourer-Cartan equation, 73
multi-index, 35
natural bundle, 100
NI, 152
Nijenhuis differential, 29
nilpotent derivation, 151
Noether identities, 154
complete, 155
first stage, 156
non-trivial, 156
trivial, 156
first-stage
complete, 157
higher-stage, 158
non-trivial, 154
Noether theorem
first, 138
second
direct, 142
inverse, 159
odd element, 114
odd morphism, 115
on-shell, 138
open map, 8
paracompact space, 184
partition of unity, 184
path, 54
Poincaré lemma, 192
presheaf, 179
canonical, 179
principal
automorphism, 88
of a connection bundle, 89
of an associated bundle, 93
bundle, 77
continuous, 77
connection, 83
associated, 94
canonical, 87
conjugate, 86
vector field, 89
vertical, 89
product connection, 56
proper cover, 12
proper map, 9
pull-back
bundle, 14
connection, 53
form, 26
section, 14
vertical-valued form, 31
rank of a morphism, 8
reduced structure, 95
reduced subbundle, 96
reducible connection, 57
representation of a Lie algebra, 72
resolution, 173
of a sheaf, 186
restriction of a bundle, 14
Ricci tensor, 105
right derivation, 151
right structure constants, 71
right-invariant form, 73
ring, 163
graded, 114
local, 163
ringed space, 188
section, 9
global, 9
integral, 56
local, 9
of a jet bundle, 34
integrable, 34
zero-valued, 16
Serre-Swan theorem, 190
for graded manifolds, 123
sheaf, 179
acyclic, 186
constant, 179
fine, 187
locally free, 189
of constant rank, 189
of finite type, 189
of continuous functions, 179
of derivations, 189
of graded derivations, 125
of modules, 189
of smooth functions, 179
sheaf cohomology, 183
smooth manifold, 5
soldered curvature, 58
soldering form, 30
basic, 30
soul map, 116
splitting domain, 122
stalk, 179
strength, 85
canonical, 87
form, 87
of a linear connection, 94
structure group, 74
action, 77
reduction, 95
structure module
of a sheaf, 179
of a vector bundle, 16
structure ring of a graded manifold, 122
structure sheaf
of a graded manifold, 122
of a ringed space, 188
of a vector bundle, 190
subbundle, 13
submanifold, 8
submersion, 8
continuous, 11
superspace, 117
supersymmetry, 150
supervector space, 117
symmetry, 138
classical, 139
exact, 138
gauge, 140
generalized, 139
variational, 137
symmetry current, 139
tangent bundle, 7
affine, 21
vertical, 19
tangent morphism, 7
vertical, 19
tangent prolongation
of a group action, 72
of a structure group action, 79
tangent-valued form, 28
canonical, 28
horizontal, 30
projectable, 30
tensor algebra, 168
tensor bundle, 19
tensor product
of Abelian groups, 165
of commutative algebras, 165
of complexes, 174
of graded modules, 115
of modules, 165
of vector bundles, 17
tensor product connection, 59
torsion form, 58
of a world connection, 60
soldering, 105
total derivative, 34
graded, 146
higher order, 36
infinite order, 48
total space, 9
transition functions, 10
$G$-valued, 74
trivial extension of a sheaf, 189
trivialization chart, 11
trivialization morphism, 10
tubular neighborhood, 25
typical fibre, 10
variational
bicomplex, 132
graded, 146
complex, 132 graded, 146 short, 136
derivative, 134
formula, 134
operator, 132
graded, 146
symmetry, 137 classical, 139
of a graded Lagrangian, 150
vector bundle, 15
characteristic, 123
dual, 17
graded, 153
vector field, 22
complete, 22
fundamental, 79
generalized, 136
graded, 125
holonomic, 107
horizontal, 52
standard, 107
integrable, 39
left-invariant, 70
parallel, 106
principal, 89
projectable, 22
on a jet manifold, 39
right-invariant, 70
subordinate to a distribution, 23
vertical, 23
vector space, 164
graded, 114
vector-valued form, 31
vertical automorphism, 13
vertical splitting, 19
of a vector bundle, 19
of an affine bundle, 21
vertical-valued form, 30
weak conservation law, 138
Whitney sum
of vector bundles, 17
world connection, 59
affine, 108
linear, 103
on a tensor bundle, 60
on the cotangent bundle, 59
symmetric, 60
world metric, 61

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