

# Stars and Watermelons

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The  **$p$ -vicious walker model** of length  $2n$  consists of  $p$  lattice paths  $W_1, W_2, \dots, W_p$  in  $\mathbb{Z}^2$  where

- $W_k$  starts at the point  $(0, a_k)$  and ends at the point  $(2n, b_k)$  for  $k = 1, \dots, p$
- all steps are directed northeast or southeast (that is, from  $(i, j)$  to  $(i + 1, j + 1)$  or to  $(i + 1, j - 1)$ )
- if  $k \neq \ell$ , then  $W_k$  and  $W_\ell$  never intersect (hence  $a_k \neq a_\ell$  and  $b_k \neq b_\ell$ , for instance).

In  **$p$ -star** configurations,  $a_k = 2k - 2$  for each  $k$  (with no constraint on  $b_k$ ); in  **$p$ -watermelon** configurations,  $b_k = 2k - 2$  as well [1, 2]. We often think of the horizontal axis as time and the vertical axis as space, writing  $W_k(0) = a_k$  and  $W_k(2n) = b_k$ . A  $p$ -watermelon with a **wall** has the additional property that

- $W_k(i) \geq 0$  for all  $0 \leq i \leq 2n$ , for all  $k$ .

Gillet [3] demonstrated that  $\lim_{n \rightarrow \infty} W_k(\lfloor 2nt \rfloor) / \sqrt{2n}$  tends to a family of  $p$  nonintersecting Brownian excursions,  $0 \leq t \leq 1$ , as an extension of a principle given in [4].

The **height** of a path  $W_k$  in a  $p$ -watermelon with wall is the maximum value of  $W_k(i)$  over all  $i$ . The **area** under a path  $W_k$  is the area of the polygonal region determined by the curve  $j = W_k(i)$ , the horizontal line  $j = 0$ , and the vertical lines  $i = 0, i = 2n$ . In the case  $p = 2$ , we will refer to the upper height and upper area (corresponding to  $W_2$ ) and the lower height and lower area (corresponding to  $W_1$ ).

Counting all 1-watermelons with wall (or Dyck paths) and 2-watermelons with wall give

$$\frac{(2n)!}{n!(n+1)!}, \quad \frac{6(2n)!(2n+2)!}{n!(n+1)!(n+2)!(n+3)!}$$

possible configurations of length  $2n$ , respectively. (The former is the  $n^{\text{th}}$  Catalan number.) The average height  $H_1(n)$  for 1-watermelons with wall satisfies [5, 6]

$$H_1(n) \sim \sqrt{\pi n}$$

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as  $n \rightarrow \infty$  and the average area  $A_1(n)$  satisfies [7, 8]

$$A_1(n) \sim \sqrt{\pi} n^{3/2}.$$

To go to the average  $L_\infty$ -norm of Brownian excursion, divide the  $H_1$  result by  $\sqrt{2n}$  (space dimension only), yielding  $\sqrt{\pi/2}$ . To go to the average  $L_1$ -norm, divide the  $A_1$  result by  $(2n)^{3/2}$  (both time and space considered), yielding  $\sqrt{\pi/8}$ . Exact formulas for  $H_1(n)$  and  $A_1(n)$  are also available [9].

The average upper height  $H_2(n)$  for 2-watermelons with wall satisfies

$$H_2(n) \sim (2.57758\dots)\sqrt{n} \sim (1.822625\dots)\sqrt{2n},$$

a new result due to Fulmek [6]. The coefficient can be expressed as a linear combination of several complicated integrals of theta functions; a certain double Dirichlet series also plays a role in the proof. Numerical results for  $3 \leq p \leq 5$  and for higher moments were obtained by Feierl [10]. A different method was proposed in [11]. To go to the average upper  $L_\infty$ -norm of Brownian excursion, divide the  $H_2$  result by  $\sqrt{2n}$ . An exact formula for  $H_2(n)$  is also available [12]. Similar information about the lower height is not known.

An exact formula for  $A_2(n)$  seems to be an open problem. Interestingly, we have both average upper/lower  $L_1$ -norm results for Brownian excursion:

$$\frac{5}{8}(\sqrt{2}-1)\sqrt{\pi}, \quad \frac{5}{8}\sqrt{\pi}$$

due to Tracy & Widom [13]. Multiplying each constant by  $(2n)^{3/2}$  therefore provides the main asymptotic terms for average upper/lower areas under 2-watermelons with wall. Numerical results in [13] also apply for  $3 \leq p \leq 9$ . In a study of average upper  $L_1$ -norms as  $p \rightarrow \infty$ , the constant 1.7710868074... arises [14, 15] and thus random matrix theory lurks nearby.

Counting all 1-watermelons without wall (or bilateral Dyck paths) and 2-watermelons without wall give [16]

$$\frac{(2n)!}{(n!)^2}, \quad \frac{(2n)!(2n+1)!}{(n!)^2((n+1)!)^2}$$

possible configurations of length  $2n$ , respectively. (The former is the  $n^{\text{th}}$  central binomial coefficient.) These tend to Brownian bridges as  $n \rightarrow \infty$  [3, 17]. In the same way,  $p$ -stars with wall tend to Brownian meanders and  $p$ -stars without wall tend to Brownian motions. Corresponding questions about average heights and average areas (suitably generalized) for  $p \geq 2$  seem to be unanswered.

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