## §2. Complex 2-Forms: Cauchy-Pompeiu's Formula

In $\S 3.1$ we have seen $d f=f_{z} d z+f_{\bar{z}} d \bar{z}$, (where $f_{z}=\partial f / \partial z$ and $f_{\bar{z}}=\partial f / \partial \bar{z}$,) for a complex (smooth) function $f$. Recall: $z=x+i y, \bar{z}=x-i y, d z=d x+i d y$ and $d \bar{z}=d x-i d y$. In general, a complex 1-form on a set in the complex plane is an expression of the form $\omega=g d z+h d \bar{z}$, where $g$ and $h$ are complex functions. By using exterior products, we can build complex 2-forms from complex 1-forms.

Example 2.1. Verify the following identities:

$$
\begin{equation*}
d z \wedge d z=0, \quad d \bar{z} \wedge d \bar{z}=0, \quad \text { and } \quad d \bar{z} \wedge d z=-d z \wedge d \bar{z}=2 i d x \wedge d y \tag{2.1}
\end{equation*}
$$

(Suggestion: keep these identities in mind.)
Solution. Here we go:

$$
\begin{aligned}
d z \wedge d z & =(d x+i d y) \wedge(d x+i d y) \\
& =d x \wedge d x+i d y \wedge d x+i d x \wedge d y-d y \wedge d y \\
& =0+i(-d x \wedge d y)+i d x \wedge d y-0=0 \\
d \bar{z} \wedge d z & =(d x-i d y) \wedge(d x+i d y) \\
& =d x \wedge d x-i d y \wedge d x+i d x \wedge d y+d y \wedge d y \\
& =0-i(-d x \wedge d y)+i d x \wedge d y+0=2 i d x \wedge d y
\end{aligned}
$$

In the same way we can show that $d \bar{z} \wedge d \bar{z}=0$ and $-d z \wedge d \bar{z}=2 i d x \wedge d y$.
Given 1-forms $\omega_{1}=f_{1} d z+g_{1} d \bar{z}$ and $\omega_{2}=f_{2} d z+g_{2} d \bar{z}$, using (2.1), we have

$$
\begin{align*}
\omega_{1} \wedge \omega_{2} & =\left(f_{1} d z+g_{1} d \bar{z}\right) \wedge\left(f_{2} d z+g_{2} d \bar{z}\right)=f_{1} g_{2} d z \wedge d \bar{z}+g_{1} f_{2} d \bar{z} \wedge d z \\
& =\left(f_{1} g_{2}-f_{1} h_{2}\right) d z \wedge d \bar{z} \equiv\left|\begin{array}{ll}
f_{1} & g_{1} \\
f_{2} & g_{2}
\end{array}\right| d z \wedge d \bar{z} \tag{2.2}
\end{align*}
$$

The differential $d \omega$ of a complex 1-form $\omega=g d z+h d \bar{z}$ is the complex 2-form given by $d \omega=d g \wedge d z+d h \wedge d \bar{z}$. More explicitly,

$$
\begin{align*}
d \omega & =d g \wedge d z+d h \wedge d \bar{z}=\left(g_{z} d z+g_{\bar{z}} d \bar{z}\right) \wedge d z+\left(h_{z} d z+h_{\bar{z}} d \bar{z}\right) \wedge d \bar{z}  \tag{2.3}\\
& =g_{\bar{z}} d \bar{z} \wedge d z+h_{z} d z \wedge d \bar{z}=\left(g_{\bar{z}}-h_{z}\right) d \bar{z} \wedge d z
\end{align*}
$$

(There is no need to memorize (2.2) and (2.3).)

Example 2.2. Verify that 1 -form $\omega=z d \bar{z}+\bar{z} d z$ is closed, that is, $d \omega=0$.

Solution: $d \omega=d z \wedge d \bar{z}+d \bar{z} \wedge d z=d z \wedge d \bar{z}-d z \wedge d \bar{z}=0$. Alternatively, we can proceed as follows: $d(z d \bar{z}+\bar{z} d z)=d(d(z \bar{z}))=0$, in view of $d^{2}=0$.

It is convenient to introduce the following "partial differentials":

$$
\partial f=\frac{\partial f}{\partial z} d z, \quad \bar{\partial} f=\frac{\partial f}{\partial \bar{z}} d \bar{z}
$$

(Note: in some books $\bar{\partial} f$ stands for $\partial f / \partial \bar{z}$.) Then identity $d f=f_{z} d z+f_{\bar{z}} d \bar{z}$ becomes $d f=\partial f+\bar{\partial} f$. So we simply write $d=\partial+\bar{\partial}$. For a 1-form $\omega=f d z+g d \bar{z}$, we have

$$
\bar{\partial} \omega=\bar{\partial} f \wedge d z+\bar{\partial} g \wedge d \bar{z}=f_{\bar{z}} d \bar{z} \wedge d z+g_{\bar{z}} d \bar{z} \wedge d \bar{z}=f_{\bar{z}} d \bar{z} \wedge d z
$$

Similarly we have $\partial \omega=g_{z} d z \wedge d \bar{z}$. Notice that $\partial^{2} f=\partial(\partial f)=\partial\left(f_{z} d z\right)=0$. Thus we write $\partial^{2}=0$. Similarly we have $\bar{\partial}^{2}=0$. We already know that $d^{2}=0$ (Rule E6). So

$$
0=d^{2}=(\partial+\bar{\partial})(\partial+\bar{\partial})=\partial^{2}+\partial \bar{\partial}+\bar{\partial} \partial+\bar{\partial}^{2}=\partial \bar{\partial}+\bar{\partial} \partial
$$

So we have $\partial \bar{\partial}+\bar{\partial} \partial=0$. This makes us curious about $\bar{\partial} \partial$. Well,

$$
\bar{\partial} \partial f=\bar{\partial}\left(f_{z} d z\right)=\bar{\partial} f_{z} \wedge d z=f_{\bar{z} z} d \bar{z} \wedge d z=\frac{1}{4} \Delta f .2 i d x \wedge d y=\frac{i}{2} \Delta f d x \wedge d y
$$

where $\Delta f=f_{x x}+g_{y y}$, the Laplacian of $f$; (see (1.6) in $\S 3.1$ for $f_{z \bar{z}}=\frac{1}{4} \Delta f$ ).
In $\S 3.1$ we have introduced the complex 1 -form $z^{-1} d z$ to study the complex logarithm $\log z$. Here we entertain a slightly more general expression $\chi=\left(z-z_{0}\right)^{-1} d z$ for some fixed $z_{0}=x_{0}+i y_{0}$, which is defined for all $z$ except $z_{0}$. Notice that

$$
d \chi=d\left(z-z_{0}\right)^{-1} \wedge d z=(-1)\left(z-z_{0}\right)^{-2} d z \wedge d z=0
$$

Thus $\chi$ is a closed 1 -form. However, $\chi$ is not an exact 1 -form, that is, there is no function $f$ defined on the complex plane with the point $z_{0}$ deleted such that $d f=\chi$. We have seen the kind of argument at the end of $\S 1.1$ leading to this conclusion. But let us repeat here using different symbols. Take a loop around $z_{0}$, say $\gamma(t)=z_{0}+r e^{i t}(0 \leq t \leq 2 \pi)$, where $r$ is any positive number and $e^{i t}$ stands for $\cos t+i \sin t$. Then

$$
\int_{\gamma} \chi=\int_{\gamma} \frac{d z}{z-z_{0}}=\int_{0}^{2 \pi} \frac{d\left(z_{0}+r e^{i t}\right)}{\left(z_{0}+r e^{i t}\right)-z_{0}}=\int_{0}^{2 \pi} \frac{r \times i e^{i t} d t}{r e^{i t}}=\int_{0}^{2 \pi} i d t=2 \pi i
$$

Thus we have:

$$
\begin{equation*}
\frac{1}{2 \pi i} \int_{\gamma} \frac{d z}{z-z_{0}} d z=1 \tag{2.4}
\end{equation*}
$$

(Suggestion: keep this identity in mind.) If $\chi$ were exact, say $\chi=d f$, we would have $\int_{\gamma} \chi=\int_{\gamma} d f=f(\gamma(2 \pi))-f(\gamma(0))=f\left(z_{0}+r\right)-f\left(z_{0}+r\right)=0$. So $\chi$ is not exact.

Take any complex function $f$ and differentiate the 1-form $f \chi$ :

$$
\begin{align*}
d(f \chi) & =d f \wedge \chi+f d \chi=d f \wedge \chi \quad[\text { because of the closedness of } \chi] \\
& =\left(f_{z} d z+f_{\bar{z}} d \bar{z}\right) \wedge\left(z-z_{0}\right)^{-1} d z \\
& =f_{\bar{z}}\left(z-z_{0}\right)^{-1} d \bar{z} \wedge d z \quad[\text { because } d z \wedge d z=0]  \tag{2.5}\\
& \equiv \frac{f_{\bar{z}}}{z-z_{0}} d \bar{z} \wedge d z=2 i \frac{f_{\bar{z}}}{z-z_{0}} d x \wedge d y .
\end{align*}
$$

Now we take a domain $D$ in the complex plane and apply Stokes' Theorem to $d(f \chi)$ over $D$. We have to be careful in doing this in view of the singularity of $\chi$ presented at $z_{0}$. More precisely, $D$ is a bounded open set in the complex plane with (positively oriented) boundary $\partial D$ consisting of finitely many smooth curves. Let $f$ be a smooth function defined on its closure $\bar{D}(=D \cup \partial D)$. Take a point $z_{0}$ in $D$ and consider the 1-form $\omega=f \chi \equiv f(z)\left(z-z_{0}\right)^{-1} d z$, which is defined on $\bar{D}$ except at $z_{0}$. Due to the singularity of $\omega$ at $z_{0}$, we remove a tiny disk $D\left(z_{0} ; r\right)$ centered at $z_{0}$ from $D$; here $D\left(z_{0} ; r\right)=\left\{z| | z-z_{0} \mid \leq r\right\}$, the disk centered at $z_{0}$ with radius $r$. Assume that the radius $r$ is so small that $D\left(z_{0} ; r\right)$ is contained in $D$. Let $D_{r}=D \backslash D\left(z_{0} ; r\right)$, which is the remaining part of $D$ when this disk is removed. Notice that $\partial D_{r}$, the boundary of $D_{r}$, consists of the boundary $\partial D$ of $D$, together with the boundary $\partial D\left(z_{0} ; r\right)$ of $D\left(z_{0} ; r\right)$ oriented in the opposite way. Thus we write $\partial D_{r}=\partial D \cup\left(-\partial D\left(z_{0} ; r\right)\right)$. Now $\omega$ has no singularity in $D_{r}$, we may apply Stokes' Theorem to $\omega$ over $D_{r}$ :

$$
\begin{equation*}
\iint_{D_{r}} d \omega=\int_{\partial D_{r}} \omega=\int_{\partial D} \omega-\int_{\partial D\left(z_{0} ; r\right)} \omega . \tag{2.6}
\end{equation*}
$$

Here $\partial D\left(z_{0} ; r\right)$ is a circle $\gamma$ around the point $z_{0}$ we have seen before: it is described by the parametric equation $\gamma(t)=z_{0}+r e^{i t}(1 \leq t \leq 2 \pi)$. Since $f\left(z_{0}+r e^{i t}\right)$ converges to $f\left(z_{0}\right)$ as the radius $r$ shrinks to 0 , the last line integral in (2.6)

$$
\int_{\partial D\left(z_{0} ; r\right)} \omega=\int_{0}^{2 \pi} \frac{f\left(z_{0}+r e^{i t}\right)}{\left(z_{0}+e^{i t}\right)-z_{0}} d e^{i t}=i \int_{0}^{2 \pi} f\left(z_{0}+r e^{i t}\right) d t
$$

converges to $i \int_{0}^{2 \pi} f\left(z_{0}\right) d t=2 \pi i f\left(z_{0}\right)$ as $r$ decreases to 0 . Letting $r \rightarrow 0$ in (2.6), the disk $D\left(z_{0} ; r\right)$ diminishes and $D_{r}$ fills up $D$. Consequently we have

$$
\begin{equation*}
\iint_{D} d \omega=\int_{\partial D} \omega-2 \pi i f\left(z_{0}\right) \tag{2.7}
\end{equation*}
$$

Now $\omega=\left(z-z_{0}\right)^{-1} f(z) d z$ and $d \omega=\left(z-z_{0}\right)^{-1} f_{\bar{z}} d \bar{z} \wedge d z$; see (2.5). Substituting them in (2.7) and rearranging terms, we arrive at the celebrated Cauchy-Pompeiu's formula:

$$
\begin{equation*}
f\left(z_{0}\right)=\frac{1}{2 \pi i} \int_{\partial D} \frac{f(z)}{z-z_{0}} d z-\frac{1}{2 \pi i} \iint_{D} \frac{f_{\bar{z}}}{z-z_{0}} d \bar{z} \wedge d z . \tag{2.8}
\end{equation*}
$$

We discuss two special cases. First, when $f$ is analytic, $f_{\bar{z}}=0$ and hence we have

$$
\begin{equation*}
f\left(z_{0}\right)=\frac{1}{2 \pi i} \int_{\partial D} \frac{f(z)}{z-z_{0}} d z \tag{2.9}
\end{equation*}
$$

which is the well-known Cauchy formula in complex analysis. (Many elementary textbooks in complex analysis set are written in such a way that as though the authors designate half of the time for establishing this formula, and treat the rest as applications of it.) This formula tells us that the value of $f$ at any point $z_{0}$ in $D$ is completely determined by the values of $f$ on the boundary of $D$. Note that Cauchy's formula tells us that (2.4) holds for any closed curve $\gamma$ surrounding $z_{0}$, not just for a circle with $z_{0}$ as its center. Secondly, when $f$ vanishes on the boundary of $D$, that is, $\left.f\right|_{\partial D}=0$, (2.8) becomes

$$
\begin{equation*}
f(\zeta)=\frac{1}{2 \pi i} \iint_{D} \frac{f_{\bar{z}}(z)}{\zeta-z}=\frac{1}{\pi} \iint_{D} \frac{f_{\bar{z}}(z)}{\zeta-z} d x \wedge d y \tag{2.10}
\end{equation*}
$$

where $z_{0}$ is switched to $\zeta$. This identity suggests how to find a solution to the following so-called $\bar{\partial}$-equation (pronounced as "dee bar equation")

$$
\begin{equation*}
\frac{\partial u}{\partial \bar{z}}=g \tag{2.11}
\end{equation*}
$$

(which is a first order elliptic partial differential equation). Here $g$ is a continuous function on the complex plane with a compact support (that is $g(z)=0$ if $|z|$ is large enough.) Take a large circular disk $D$ in the complex plane so that $g$ vanishes outside $D$. Let us restrict our attention to solutions on $D$, that is, those $u$ defined on $D$ so that the identity $\partial u / \partial \bar{z}=g$ holds on $D$. Notice that if $u$ is a solution and $v$ is an analytic function on $D$, then $u+v$ is another solution; indeed $(\partial / \partial \bar{z})(u+v)=(\partial / \partial \bar{z}) u+(\partial / \partial \bar{z}) v=g+0=g$. We can reverse the above argument to show that the general solution to (4.11) is of the form $u+v$, where $v$ is an analytic function. Thus, it is enough to find a special solution to (2.11). Suppose that $u$ is indeed a solution. Then, by Cauchy-Pompeiu's formula, for all $\zeta \in D$,

$$
u(\zeta)=\frac{1}{2 \pi i} \int_{\partial D} \frac{u(z)}{z-\zeta} d z-\frac{1}{\pi} \iint_{D} \frac{u_{\bar{z}}}{z-\zeta} d x \wedge d y=v(\zeta)+\frac{1}{\pi} \iint_{D} \frac{g(z)}{\zeta-z} d x \wedge d y
$$

where $v(\zeta)=\frac{1}{2 \pi i} \int_{\partial D} \frac{u(z)}{z-\zeta} d z$ can be shown to be analytic as a function of $\zeta \in D$ : indeed, by the "punching and kicking" principle, we have

$$
\frac{\partial}{\partial \bar{\zeta}} v(\zeta)=\frac{1}{2 \pi i} \int_{\partial D} f(z) \frac{\partial}{\partial \bar{\zeta}}\left(\frac{1}{z-\zeta}\right) d z=0
$$

Thus, assuming (2.11) has a solution, the function $u_{0}$ given by

$$
\begin{equation*}
u_{0}(\zeta)=\frac{1}{\pi} \iint_{D} \frac{g(z)}{\zeta-z} d x \wedge d y \tag{2.12}
\end{equation*}
$$

is a solution to this $\bar{\partial}$-equation. Note that we have not established the existence of a solution to (2.11) yet. What we have done is that, if solutions do exist, then one of them can be explicitly written down as (2.12). In order to show that (2.11) has a solution, the most natural thing to do is to check directly that $u_{0}$ given by (2.12) is indeed a solution. This needs some delicate work and we do not pursue this endeavor. The function $u_{0}$ in by (2.12) is called the Cauchy transform of $g$ and

$$
C(\zeta, z) \equiv \frac{1}{\zeta-z}
$$

is called the Cauchy kernel of this integral transform.

## Exercises

1. The conjugate of a 1 -form $\omega=f d x+g d y$ is defined to be $\bar{\omega}=\bar{f} d x+\bar{g} d y$. Verify the following identities: $\overline{d z}=d \bar{z}, \overline{d \bar{z}}=d z, \overline{\partial f}=\bar{\partial} \bar{f}, \overline{\bar{\partial} f}=\partial \bar{f}$.
2. (a) Verify that $(\partial-\bar{\partial})^{2}=0$. (b) Verify the product rule: $\bar{\partial}(f \omega)=\bar{\partial} f \wedge \omega+f \bar{\partial} \omega$.
3. Is the identity $\bar{\partial}\left(g^{*} f\right)=g^{*} \bar{\partial} f$ true ?
(Hint: Take any analytic function $f$ and any $g$ so that $f \circ g$ is not analytic.)
4. Let $f$ be an analytic function and write $w=u+i v=f(z)=f(x+i y)$ as usual. Then $d w=f^{\prime}(z) d z$ and $d \bar{w}=\overline{f^{\prime}(z)} d \bar{z}$. (a) Verify $d \bar{w} \wedge d w=\left|f^{\prime}(z)\right|^{2} d \bar{z} \wedge d z$ and compare it to $d u d v=\{\partial(u, v) / \partial(x, y)\} d x d y$ to deduce

$$
\begin{equation*}
\left|f^{\prime}(z)\right|^{2}=\frac{\partial(u, v)}{\partial(x, y)} \tag{2.13}
\end{equation*}
$$

(b) Use the Cauchy-Riemann equation to check (2.13) directly.
5. We extend Hodge's $*$-operator (see (2.1) in $\S 1.2$ ) to complex differential forms $\omega=$ $P d x+Q d y$ (where $P$ and $Q$ are complex-valued functions of $x$ and $y$ ) by putting

$$
\begin{equation*}
* \omega=*(P d x+Q d y)=\bar{P} * d x+\bar{Q} * d y \equiv \bar{P} d y-\bar{Q} d x . \tag{2.14}
\end{equation*}
$$

(a) Check that $* * \omega=-\omega$.
(b) Check that $* d z=i d \bar{z}$ and $* d \bar{z}=-i d z$.
(c) Verify that the Cauchy Riemann equations $u_{x}=v_{y}, u_{y}=-v_{x}$ (see (1.8) in §3.1) can be rewritten as $* d u=d v$.
(d) Check that a real valued function $w=w(x, y)$ is harmonic if and only if $d * d w=0$.
6. We say that a real-valued function $v=v(x, y)$ is a harmonic conjugate of a realvalued function $u=u(x, y)$ if $d v=* d u$. From part (c) of the previous exercise we know that the last condition is equivalent to $u+i v$ being a holomorphic function of $z=x+i y$. In each of the following parts, show that the given function $u=u(x, y)$ is harmonic and find a harmonic conjugate of it.
(a) $u=x^{2}-y^{2}$,
(b) $u=x^{3}-3 x y^{2}$
(c) $u=e^{x} \cos y$
(d) $u=\log \sqrt{x^{2}+y^{2}}$.
7. Let $u=u(x, y)$ be a harmonic function on $\mathbf{D}^{o}=\{z=x+i y \in \mathbf{C}:|z|<1\}$ such that $u(0,0)=0$. Check that $\partial u / \partial z$ is an analytic function and the complex 1-form $(\partial u / \partial z) d z$ is closed. For each $z$ in $\mathbf{D}^{o}$, take any path linking 0 to $z$ and define

$$
f(z)=2 \int_{\gamma} \frac{\partial u}{\partial z} d z
$$

Show that $f$ is a well defined analytic function, with $\partial f / \partial z=2 \partial u / \partial z$. Finally, check that the real part of $f$ is $u$.
8. The punctured plane $\mathbf{C}_{0}$ is defined to be the complex plane with the origin removed, that is, $\mathbf{C}_{0}=\mathbf{C} \backslash\{0\}$. The reflection of the punctured plane (with respect to the unit circle) is defined to be the transformation $R$ of $\mathbf{C}_{0}$ given by $R(z)=1 / \bar{z}$.
(a) Verify that $R$ is idempotent in the sense that $R(R(z))=z$.
(b) Verify that three points $0, z$ and $R(z)$ are collinear, and $|z \| R(z)|=1$. Make a simple sketch (including the unit circle, the origin, the ray from the origin containing $z$ and $R(z))$ to understand the effect of $R$. Also check that $R$ leaving all points on the unit circle fixed, that is, $R(z)=z$ if $|z|=1$.
(c) Check that if we identify the complex variable $z=x+i y$ with $(x, y)$, then

$$
R(x, y)=\left(\frac{x}{x^{2}+y^{2}}, \frac{y}{x^{2}+y^{2}}\right)
$$

(The above identity suggests how to generalize $R$ to higher dimensional spaces.)
(d) Check that $R^{*}\left(z^{-1} d z\right)=-\bar{z}^{-1} d \bar{z}$ and $R^{*}\left(\bar{z}^{-1} d \bar{z}\right)=-z^{-1} d z$; (recall that $R^{*}$ stands for pull back by $R$ ). Deduce that the 2 -form $\Omega \equiv|z|^{-2} d z \wedge d \bar{z}$ is invariant under $R$ in the sense that $R^{*} \Omega=\Omega$ and consequently the 2 -form $\left(x^{2}+y^{2}\right)^{-1} d x \wedge d y$ is also invariant under $R$.
(e) Verify the following remarkable identity the reflection $R$ :

$$
\frac{\left|e^{i \theta}-z\right|}{\left|e^{i \theta}-R(z)\right|}=|z|,
$$

that is, the LHS is independent of which point $e^{i \theta}$ on the unit circle is chosen. (Remark: the image $R(z)$ of $z$ under the reflexion $R$ is often denoted by $z^{*}$. We avoid using this notation because it crashes with pullbacks.)
9. Let $\gamma$ be the elliptic path given by $z(t)=a \cos t+i b \sin t(0 \leq t \leq 2 \pi)$, where $a$ and $b$ are positive constants. Use the identity $\int_{\gamma} d z / z=2 \pi i$ to verify

$$
\int_{0}^{2 \pi} \frac{\cos t \sin t}{a^{2} \cos ^{2} t+b^{2} \sin ^{2} t} d t=0 \quad \text { and } \quad \int_{0}^{2 \pi} \frac{d t}{a^{2} \cos ^{2} t+b^{2} \sin ^{2} t}=\frac{2 \pi}{a b}
$$

(The first identity can also be obtained by changing the variable $t$ to $s=2 \pi-t$.)
10. Let $\mathbf{D}$ be the unit disk and let $m, n$ be integers $\geq 0$. Consider

$$
I=\iint_{\mathbf{D}} \bar{z}^{m} z^{n} d \bar{z} \wedge d z
$$

Compute this double integral by two different methods: first, use Green's theorem (2.2); second, use polar coordinates: $d \bar{z} \wedge d z=2 i d x \wedge d y=2 i r d r d \theta$. (You need to use the identity $\int_{0}^{2 \pi} e^{i(m-n) t} d t=2 \pi \delta_{m n}$, which can be easily checked.)
11. Let $\mathbf{D}$ be the unit disk and let $\chi$ be its characteristic function: $\chi(z)=1$ if $z$ belongs to $\mathbf{D}$ and $=0$ otherwise. Consider $\varphi(\zeta)=(2 \pi i)^{-1} \iint(\zeta-z)^{-1} \chi d \bar{z} \wedge d z$, the Cauchy transform of $\chi$.
(a) Use polar coordinates to compute $\varphi(0)$ directly.
(b) Use Cauchy-Pompeiu's formula to find $\varphi(\zeta)$ for $\zeta$ in the interior of $\mathbf{D}:|\zeta|<1$.
(c) Can you find $\varphi(\zeta)$ for general $\zeta$ ?

