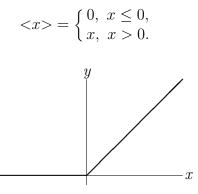
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Singularity functions

Ramp function

The Macauley bracket $<\cdot>$ defines the ramp function



Obviously.

$$\langle x - a \rangle = \begin{cases} 0, & x \le a, \\ x - a, & x > a, \end{cases}$$

 $\langle x - a \rangle^n = \begin{cases} 0, & x \le a, \\ (x - a)^n, & x > a \end{cases}$

for n > 0, and

$$\frac{d}{dx} < x - a >^{(n+1)} = (n+1) < x - a >^n,$$

so that

$$\int \langle x - a \rangle^n \, dx = \frac{1}{n+1} \langle x - a \rangle^{(n+1)}.$$

For example, in a beam loaded as shown above, the bending moment is given by

$$M(x) = \begin{cases} \frac{Fbx}{L}, & x < a \\ \frac{Fa(L-x)}{L}, & x > a \end{cases}$$

But

$$\frac{Fa(L-x)}{L} = \frac{Fbx}{L} - F(x-a);$$

therefore the moment can be given by the single expression

$$M(x) = \frac{Fbx}{L} - F < x - a >.$$

This can now be integrated twice to give a single expression for the deflection:

$$EIv(x) = \frac{Fbx^3}{6L} - \frac{F}{6} < x - a >^3 + C_1 x + C_2.$$

The boundary condition v(0) = 0 yields $C_2 = 0$, while v(L) = 0 gives

$$C_1 = -\frac{Fb}{6L}(L^2 - b^2),$$

leading to

$$v(x) = \frac{F}{6EI} \left\{ \frac{b}{L} [x^3 - (L^2 - b^2)x] - \langle x - a \rangle^3 \right\}.$$

Step function

The derivative of the ramp function is

$$\frac{d}{dx} < x > = \begin{cases} 0, & x < 0, \\ 1, & x > 0, \end{cases}$$

known as the **Heaviside step function**. There are various notations for it: H(x), U(x), 1(x) and others (the notation $\langle x \rangle^0$ is ambiguous because it gives 0^0 for negative x). Let us use H(x). Note that H(0) is not defined; it is connventional to give it the value $\frac{1}{2}$.

The shear in the preceding example is given by

$$V(x) = \frac{F}{L} \left[b - LH(x - a) \right],$$

that is, Fb/L for x < a and -Fa/L for x > a.

The step function can be used to represent loading that is distributed over a part of a beam. For example, a uniform load of intensity w distributed over the middle third of the beam is given by

$$q(x) = -w[H(x - L/3) - H(x - 2L/3)].$$

The deflection can be obtained by integrating four times, resulting in the single expression

$$v(x) - \frac{w}{24EI}[\langle x - L/3 \rangle^4 - \langle x - 2L/3 \rangle^4] + C_1 \frac{x^3}{6} + C_2 \frac{x^2}{2} + C_3 x + C_4,$$

and the constants C_1, \ldots, C_4 can be determined from the boundary conditions at x = 0 and x = L.

Delta function

In order to represent the loading in the case of a concentrated load, we must be able to differentiate the step function: if the shear is given by

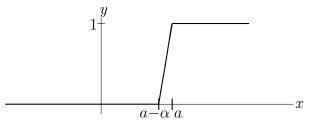
$$V(x) = \frac{F}{L} \left[b - LH(x - a) \right],$$

then the loading is

$$q(x) = \frac{d}{dx}V(x) = -F\frac{d}{dx}H(x)$$

Clearly the derivative of H(x) is zero for $x \neq 0$, but at x = 0 it is infinite, so it is not really a function (mathematicians call it a *distribution*).

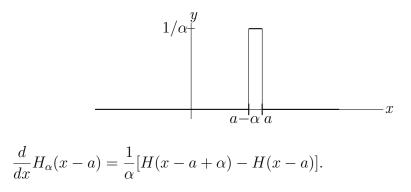
Let's consider the following approximation to the step function, say $H_{\alpha}(x)$:



This is given by

$$H_{\alpha}(x-a) = \frac{1}{\alpha}[-].$$

Note that, in the limit as $\alpha \to 0$, this is just the derivative of $\langle x - a \rangle$. But let's differentiate $H_{\alpha}(x-a)$ before taking the limit:



Formally, this becomes the derivative of H(x - a) in the limit as $\alpha \to 0$. What happens is that the rectangle in the picture gets narrower and taller, but its area remains 1. This limit is known as the **Dirac delta function** and is almost universally denoted $\delta(x-a)$, except that in Mechanics of Materials textbooks the strange notation $\langle x - a \rangle_*^{-1}$ is used; the subscript asterisk is there to tell us that the superscipt -1 is not really the power -1.

The loading on the beam with a concentrated load at x = a can now be written as

$$q(x) = -F\delta(x-a).$$

Note that it doesn't matter what the delta function "really" means; all that matters is that when we integrate it we get the step function, and so on.

By the same token we can define the derivative of the delta function, $\delta'(x-a)$ (denoted $\langle x-a \rangle_*^{-2}$ in the book). Suppose we have a downward concentrated force of magnitude C/α at x = a, and an upward one of the same magnitude at $x = a - \alpha$, forming the clockwise couple C, independent of α ; in the limit as $\alpha \to 0$, this becomes the **concentrated** couple C. Now, mathematically, the loading is given by

$$q(x) = \frac{C}{\alpha} [\delta(x - a + \alpha) - \delta(x - a)],$$

and in the limit as $\alpha \to 0$ this **formally** becomes C times the derivative of $\delta(x-a)$, which we can write as $C\delta'(x-a)$.