

# Finite length Solenoid potential and field

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The surface current density is (Jackson, 1998):

$$\vec{K} = \frac{I}{L} \delta(\rho - a), \quad z \in \left(-\frac{L}{2}, \frac{L}{2}\right)$$

The vector potential is:

$$\vec{A} = A_\phi \hat{\phi} = \hat{\phi} \frac{\mu_0 I}{4\pi L} \int_{-\frac{L}{2}}^{\frac{L}{2}} \int_0^{2\pi} \int_0^\infty \frac{\delta(\rho' - a) \cos \phi'}{\sqrt{\rho^2 + \rho'^2 + (z - z')^2 - 2a\rho \cos \phi'}} \rho' d\rho' d\phi' dz'$$

$$A_\phi = \frac{\mu_0 I a}{2\pi L} \int_{-\frac{L}{2}}^{\frac{L}{2}} \int_0^\pi \frac{\cos \phi'}{\sqrt{\rho^2 + a^2 + (z - z')^2 - 2a\rho \cos \phi'}} d\phi' dz'$$

Simplify the form by setting  $\zeta = (z - z')$  and the integration of  $\zeta$  is a log function (Edmund E. Callaghan, 1960):

$$\int_0^\pi \cos \phi' \left[ \ln \left( \zeta + \sqrt{\zeta^2 + \rho^2 + a^2 - 2\rho a \cos \phi'} \right) \right]_{\zeta_-}^{\zeta_+} d\phi' = \left[ \int_0^\pi \cos \phi' \ln(\zeta + \alpha(\zeta)) d\phi' \right]_{\zeta_-}^{\zeta_+}$$

$$\alpha(\zeta) = \sqrt{\zeta^2 + \rho^2 + a^2 - 2\rho a \cos \phi'}, \quad \zeta_\pm = z \mp \frac{L}{2}$$

Integration by path:

$$\int_0^\pi \cos \phi' \ln(\zeta + \alpha(\zeta)) d\phi' = \sin \phi' \ln(\zeta + \alpha(\zeta)) \Big|_0^{2\pi} - \int_0^\pi \sin \phi' d(\ln(\zeta + \alpha(\zeta)))$$

The first term is zero, and the derivative of  $\ln(\zeta + \alpha(\zeta))$  is:

$$\frac{d \ln(\zeta + \alpha(\zeta))}{d\phi'} = \frac{\rho a \sin \phi'}{(\alpha(\zeta) + \zeta)\alpha(\zeta)}$$

Multiple by  $(\alpha(\zeta) + \zeta)/(\alpha(\zeta) - \zeta)$

$$= \frac{(\alpha(\zeta) - \zeta)\rho a \sin \phi'}{(\alpha^2(\zeta) + \zeta^2)\alpha(\zeta)} = \frac{\rho a \sin \phi'}{(\rho^2 + a^2 - 2\rho a \cos \phi')} - \frac{\zeta \rho a \sin \phi'}{(\rho^2 + a^2 - 2\rho a \cos \phi')\alpha(\zeta)}$$

The first term on the right side appeared on both  $\alpha_\pm$ , then,

$$\int_0^\pi \sin \phi' d(\ln(\zeta + \alpha(\zeta))) = - \int_0^\pi \frac{\zeta \rho a \sin^2 \phi'}{(\rho^2 + a^2 - 2\rho a \cos \phi')\alpha(\zeta)} d\phi'$$

$$A_\phi = \frac{\mu_0 I a^2 \rho}{2\pi L} \left[ \zeta \int_0^\pi \frac{\sin^2 \phi'}{(\rho^2 + a^2 - 2\rho a \cos \phi')\sqrt{\zeta^2 + \rho^2 + a^2 - 2\rho a \cos \phi'}} d\phi' \right]_{\zeta_-}^{\zeta_+}$$

By using the change of integration interval

$$\begin{aligned}
 & \int_0^{\frac{\pi}{2}} \frac{\sin^2(2\theta)}{((a + \rho)^2 - 4\rho a \sin^2 \theta)\sqrt{\zeta^2 + (a + \rho)^2 - 4\rho a \sin^2 \theta}} d\theta \\
 &= \frac{kh^2}{8(\sqrt{a\rho})^3} \int_0^{\frac{\pi}{2}} \frac{\sin^2(2\theta)}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
 &= \frac{kh^2}{2(\sqrt{a\rho})^3} \int_0^{\frac{\pi}{2}} \frac{\sin^2 \theta - \sin^4 \theta}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
 & \quad h^2 = \frac{4a\rho}{(a + \rho)^2} \\
 & \quad k^2 = \frac{4a\rho}{(a + \rho)^2 + \zeta^2}
 \end{aligned}$$

The integral can be spitted into 2 parts, the first part is (Milton Abramowitz, 1965) (NIST Digital Library of Mathematical Functions) :

$$\begin{aligned}
 & \int_0^{\frac{\pi}{2}} \frac{\sin^2 \theta}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
 &= -\frac{1}{h^2} \int_0^{\frac{\pi}{2}} \frac{1 - h^2 \sin^2 \theta - 1}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
 &= -\frac{1}{h^2} \int_0^{\frac{\pi}{2}} \frac{1}{\sqrt{1 - k^2 \sin^2 \theta}} - \frac{1}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
 &= \frac{1}{h^2} (\Pi(h^2, k^2) - K(k^2))
 \end{aligned}$$

Here

$$\Pi(n, m) = \int_0^{\frac{\pi}{2}} \frac{1}{(1 - n \sin^2 \theta)\sqrt{1 - m \sin^2 \theta}} d\theta$$

It is the elliptic integral of 3<sup>rd</sup> kind.

The 2<sup>nd</sup> part is:

$$\int_0^{\frac{\pi}{2}} \frac{-\sin^4 \theta}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta$$

$$\begin{aligned}
&= \frac{1}{h^4} \int_0^{\frac{\pi}{2}} \frac{1 - h^4 \sin^4 \theta - 1}{(1 - h^2 \sin^2 \theta) \sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
&= \frac{1}{h^4} \int_0^{\frac{\pi}{2}} \frac{1 + h^2 \sin^4 \theta}{\sqrt{1 - k^2 \sin^2 \theta}} - \frac{1}{(1 - h^2 \sin^2 \theta) \sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
&= \frac{1}{h^4} \left( K(k^2) + \frac{h^2}{k^2} (K(k^2) - E(k^2)) - \Pi(h^2, k^2) \right)
\end{aligned}$$

Thus, combine it and we have:

$$A_\phi = \frac{\mu_0 I}{4\pi L} \sqrt{\frac{a}{\rho}} \left[ \zeta k \left( \frac{k^2 + h^2 - h^2 k^2}{h^2 k^2} K(k^2) - \frac{1}{k^2} E(k^2) + \frac{h^2 - 1}{h^2} \Pi(h^2, k^2) \right) \right]_{\zeta_-}^{\zeta_+}$$

The magnetic field is the curl

$$B_\rho = [\nabla \times A_\phi]_\rho = -\frac{\partial}{\partial z} (A_\phi) = -\frac{\partial}{\partial \zeta} (A_\phi)$$

$$B_z = [\nabla \times A_\phi]_z = \frac{1}{\rho} \frac{\partial}{\partial \rho} (\rho A_\phi) = \frac{1}{\rho} A_\phi + \frac{\partial A_\phi}{\partial \rho}$$

Using the derivative formulae for elliptic integral:

$$\frac{d}{dx} K(x) = -\frac{1}{2x} K(x) + \frac{1}{2x(1-x)} E(k)$$

$$\frac{d}{dx} E(x) = -\frac{1}{2x} K(x) + \frac{1}{2k} E(k)$$

$$\frac{d}{dx} \Pi(v, x) = \frac{1}{2(x-1)(v-x)} E(k) + \frac{1}{2(v-x)} \Pi(v, x)$$

Thus:

$$\frac{d}{dk} \left( k \left( \frac{k^2 + h^2 - h^2 k^2}{h^2 k^2} K(k^2) - \frac{1}{k^2} E(k^2) + \frac{h^2 - 1}{h^2} \Pi(h^2, k^2) \right) \right)$$

$$= -\frac{1}{k^2} K(k^2) + \frac{h^2}{k^2(h^2 - k^2)} E(k^2) + \frac{h^2 - 1}{(h^2 - k^2)} \Pi(h^2, k^2)$$

$$\frac{dk}{dz} = -\frac{k^3 \left( z \pm \frac{L}{2} \right)}{4a\rho}, \quad \frac{dk}{d\rho} = \frac{k}{2\rho} - \frac{k^3(a + \rho)}{4a\rho}$$

$$\frac{d}{dz} (zf(z)) = f(z) + z \frac{df(z)}{dz}, \quad \frac{d}{d\rho} (\sqrt{\rho} f(\rho)) = \frac{f(\rho)}{2\sqrt{\rho}} + \sqrt{\rho} \frac{df(\rho)}{d\rho}$$

Then

$$B_\rho = \frac{\mu_0 I}{2\pi L} \sqrt{\frac{a}{\rho}} \left[ \left( \frac{k^2 - 2}{2k} K(k^2) + \frac{1}{k} E(k^2) \right) \right]_{\zeta_-}^{\zeta_+}$$

Or, by using integration identity

$$\frac{d}{dx} \int_a^b f(x) dx = f(b) - f(a)$$

$$B_\rho = \frac{\mu_0 I a}{2\pi L} \left[ \int_0^\pi \frac{\cos \phi'}{\sqrt{\zeta^2 + \rho^2 + a^2 - 2\rho a \cos \phi'}} d\phi' \right]_{\zeta_-}^{\zeta_+}$$

Since the angle integration is elliptic integral.

$$B_\rho = \frac{\mu_0 I}{2\pi L} \sqrt{\frac{a}{\rho}} \left[ \left( \frac{k^2 - 2}{2k} K(k^2) + \frac{1}{k} E(k^2) \right) \right]_{\zeta_-}^{\zeta_+}, \quad k^2 = \frac{4a\rho}{(a + \rho)^2 + \zeta^2}$$

Thus, the result is double verified. We should notices that the radial component of the magnetic field is just like 2 coil separated by distance  $L$  vertically.

The z-component is :

$$B_z = -\frac{\mu_0 I}{4\pi L 2\sqrt{a\rho}} \left[ \zeta k \left( K(k^2) + \frac{a - \rho}{a + \rho} \Pi(h^2, k^2) \right) \right]_{\zeta_-}^{\zeta_+}$$

Or we can compute  $\frac{\partial A_\phi}{\partial \rho}$ , we use

$$A_\phi = \frac{\mu_0 I a}{2\pi L} \left[ \int_0^\pi \cos \phi' \ln(\zeta + \alpha(\zeta)) d\phi' \right]_{\zeta_-}^{\zeta_+}$$

By

$$\frac{\partial}{\partial \rho} \ln(\zeta + \alpha(\zeta)) = \frac{\rho - a \cos \phi'}{\alpha(\zeta)(\alpha(\zeta) + \zeta)}$$

Using the same trick

$$\begin{aligned} \frac{\rho - a \cos \phi'}{\alpha(\zeta)(\alpha(\zeta) + \zeta)} &= \frac{(\rho - a \cos \phi')(\alpha(\zeta) - \zeta)}{\alpha(\zeta)(\alpha^2(\zeta) + \zeta^2)} \\ &= \frac{(\rho - a \cos \phi')(\alpha(\zeta) - \zeta)}{\alpha(\zeta)(\rho^2 + a^2 - 2\rho a \cos \phi')} = \frac{(\rho - a \cos \phi')}{(\rho^2 + a^2 - 2\rho a \cos \phi')} - \frac{(\rho - a \cos \phi')\zeta}{\alpha(\zeta)(\rho^2 + a^2 - 2\rho a \cos \phi')} \end{aligned}$$

Therefore:

$$\frac{\partial A_\phi}{\partial \rho} = -\frac{\mu_0 I a}{2\pi L} \left[ \int_0^\pi \frac{\zeta \rho \cos \phi' - a \zeta \cos^2 \phi'}{\alpha(\zeta)(\rho^2 + a^2 - 2\rho a \cos \phi')} d\phi' \right]_{\zeta_-}^{\zeta_+}$$

Combined with

$$\frac{1}{\rho} A_\phi = \frac{\mu_0 I a}{2\pi L} \left[ \int_0^\pi \frac{a \zeta \sin^2 \phi'}{\alpha(\zeta)(\rho^2 + a^2 - 2\rho a \cos \phi')} d\phi' \right]_{\zeta_-}^{\zeta_+}$$

Then the magnetic field is

$$B_z = \frac{\mu_0 I a}{2\pi L} \left[ \int_0^\pi \frac{\zeta(a - \rho \cos \phi')}{\alpha(\zeta)(\rho^2 + a^2 - 2\rho a \cos \phi')} d\phi' \right]_{\zeta_-}^{\zeta_+}$$

With change the interval

$$\begin{aligned} & \int_0^\pi \frac{\zeta(a - \rho \cos \phi')}{(\rho^2 + a^2 - 2\rho a \cos \phi') \sqrt{\zeta^2 + \rho^2 + a^2 - 2\rho a \cos \phi'}} d\phi' \\ &= \int_0^{\frac{\pi}{2}} \frac{a\zeta + \rho\zeta \cos(2\theta)}{((a + \rho)^2 - 4\rho a \sin^2 \theta) \sqrt{\zeta^2 + (a + \rho)^2 - 4\rho a \sin^2 \theta}} d\theta \\ &= \frac{kh^2}{8(\sqrt{a\rho})^3} \int_0^{\frac{\pi}{2}} \frac{(a + \rho)\zeta - 2\rho\zeta \sin^2 \theta}{(1 - h^2 \sin^2 \theta) \sqrt{1 - k^2 \sin^2 \theta}} d\theta \\ &= \frac{kh^2}{8(\sqrt{a\rho})^3} \left( (a + \rho)\zeta \Pi(h^2, k^2) - \frac{2\rho\zeta}{h^2} (\Pi(h^2, k^2) - K(k^2)) \right) \\ &= \frac{k\zeta}{8(\sqrt{a\rho})^3} ((h^2(a + \rho) - 2\rho)\Pi(h^2, k^2) + 2\rho K(k^2)) \\ &= -\frac{k\zeta}{8(\sqrt{a\rho})^3} \left( 2\rho K(k^2) + \frac{2\rho(a - \rho)}{a + \rho} \Pi(h^2, k^2) \right) \\ &= -\frac{k\zeta}{4a\sqrt{a\rho}} \left( K(k^2) + \frac{(a - \rho)}{a + \rho} \Pi(h^2, k^2) \right) \end{aligned}$$

Thus we get the same result.

$$B_z = -\frac{\mu_0 I}{4\pi L 2\sqrt{a\rho}} \left[ \zeta k \left( K(k^2) + \frac{a - \rho}{a + \rho} \Pi(h^2, k^2) \right) \right]_{\zeta_-}^{\zeta_+}$$

In conclusion, the field is defined by:

$$A_\phi = \frac{\mu_0 I}{4\pi L} \frac{1}{\sqrt{\rho}} \left[ \zeta k \left( \frac{k^2 + h^2 - h^2 k^2}{h^2 k^2} K(k^2) - \frac{1}{k^2} E(k^2) + \frac{h^2 - 1}{h^2} \Pi(h^2, k^2) \right) \right]_{\zeta_-}^{\zeta_+}$$

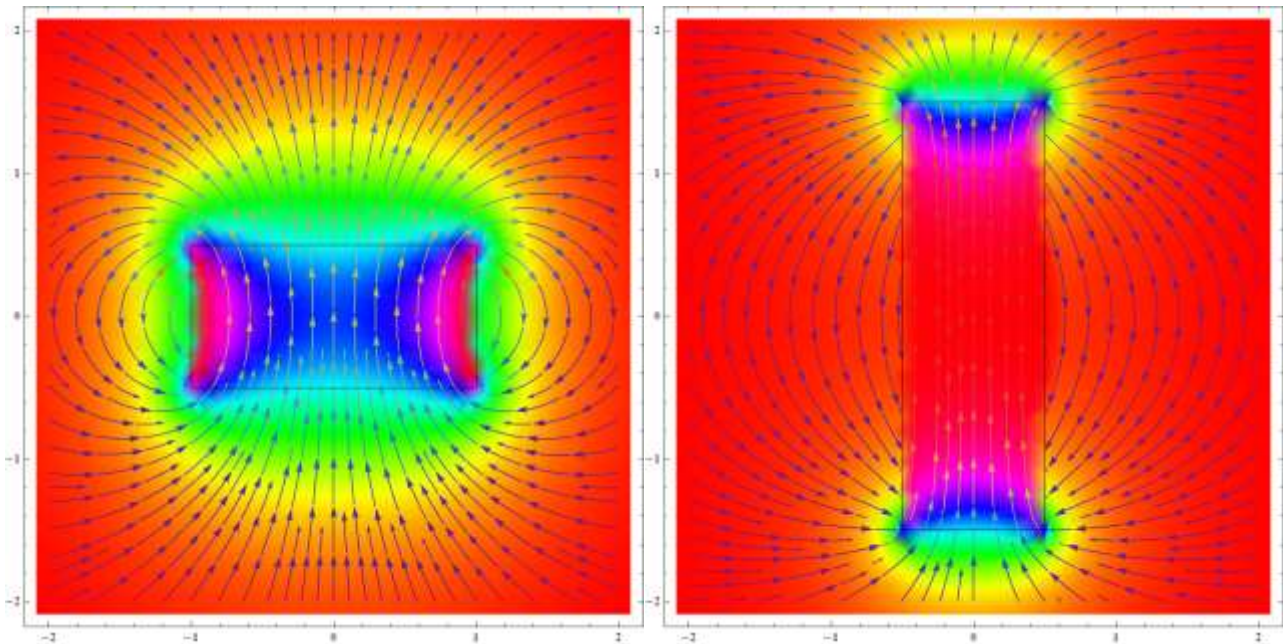
$$B_\rho = \frac{\mu_0 I}{4\pi L} \frac{1}{\sqrt{\rho}} \left[ \left( \frac{k^2 - 2}{k} K(k^2) + \frac{2}{k} E(k^2) \right) \right]_{\zeta_-}^{\zeta_+}$$

$$B_z = -\frac{\mu_0 I}{4\pi} \frac{1}{2L\sqrt{a\rho}} \left[ \zeta k \left( K(k^2) + \frac{a - \rho}{a + \rho} \Pi(h^2, k^2) \right) \right]_{\zeta_-}^{\zeta_+}$$

With

$$h^2 = \frac{4a\rho}{(a + \rho)^2}, \quad k^2 = \frac{4a\rho}{(a + \rho)^2 + \zeta^2}$$

Here is some plot



## Works Cited

- Edmund E. Callaghan, S. H. (1960). *The Magnetic Field of a Finite Solenoid (Technical note D-465)*. Washington, USA: Nation Aeronautics and Space Administration.
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