# Titanium oxide film for the bottom antireflective layer in deep ultraviolet lithography

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Titanium oxide thin film, fabricated with tetraisopropyltitanate and oxygen by electron cyclotron resonance—plasma-enhanced chemical vapor deposition, is investigated as a potential candidate for the antireflective layer in KrF excimer laser (248-nm) lithography. The oxygen flow-rate dependence of the optical properties such as the refractive index (n) and the extinction coefficient (k) of the film at the 248-nm wavelength has been characterized, and the films with the expected combinations of n and k values for the antireflective layer have been deposited. Simulation results indicate that reflectance values of less than 4% and as low as 1.2% can be reached at the interface between the photoresist and the film postulating the structures of the photoresist/300-Å  $TiO_x$  film/c-Si substrate and the W-Si substrate, respectively, by selected proper combinations of n and k values. Moreover the reflectance can be further reduced to almost zero by changing the film thickness. Thus it is found that titanium oxide thin films can be used as the bottom antireflective layer in KrF excimer laser lithography. © 1997 Optical Society of America

Key words: Bottom antirefractive layer, titanium oxide thin film, reflectance, KrF excimer laser lithography.

### 1. Introduction

In the semiconductor lithographic process, antireflective films not only permit better linewidth control but also permit the realization of designs that were previously impossible to print. New high-performance antireflective layers (ARL's) for lithography with the i-line (365 nm), the KrF (248 nm), and even the ArF (193 nm) excimer lasers have been developed for nearly a decade in the semiconductor industry.<sup>1–4</sup>

As the linewidths in semiconductor manufacturing decrease, the use of shorter-wavelength light in the projection tools is indispensable. For the 0.35- $\mu m$  and 0.25- $\mu m$  feature sizes, the wavelengths of 365 nm and 248 nm are to be used, respectively. However, the reflectivity at the interface between the photore-

sist (PR) and the substrate increases as the wavelength of the light decreases. For example, the reflectance increases from 24% at 436 nm to 53% at 248 nm for a Si substrate and from 31% to 35% for a W-Si substrate. The increase in the reflectivity at the interface between the PR and the substrate causes a critical dimension variation that is due to the multiple interference effects as well as the reflection from the substrate topography. Thus, the notching and swing effects are significantly enhanced in the lithographic process.

A simple analytical expression for the reflectivity swing ratio  $(S, R_{\text{max}}/R_{\text{min}})$  that is due to thin-film interference is as follows<sup>5</sup>:

$$S = 4\sqrt{R_1R_2} \exp(-\alpha D), \tag{1}$$

where  $R_1$  is the reflectivity at the air-photoresist interface,  $R_2$  is the reflectivity at the photoresist-substrate interface, and  $\alpha$  and D are the absorption coefficient and the average thickness of the photoresist, respectively. A top antireflective layer (TARL)<sup>6,7</sup> and a bottom antireflective layer (BARL) can reduce  $R_1$  and  $R_2$ , respectively. The schematic structures for the TARL and BARL are shown in Fig. 1. The use of the TARL significantly reduces the

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Received 19 April 1996; revised manuscript received 18 July 1996.

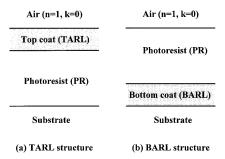


Fig. 1. Schematic structures for the (a) TARL and (b) BARL structures.

swing ratio by reducing  $R_1$ . Assuming a nonabsorbing film and normally incident radiation,  $R_1$  is reduced to zero when  $n_{\rm top\,coat} = \sqrt{n_{\rm resist}}$  (n is the refractive index) and the optical thickness of the coating is a quarter wavelength ( $\lambda/4n_{\rm top\,coat}$ ). The TARL can reduce only the swing effect without reducing the notching problem. In contrast, a BARL, which lies between the photoresist and substrate, can reduce the reflective notching by significantly reducing  $R_2$ . Under the same assumptions as the TARL,  $R_2$  is reduced to zero when  $n_{\rm bottom\,coat} = \sqrt{n_{\rm resist}}\,n_{\rm substrate}$  and  $d_{\rm bottom\,coat} = \lambda/4n_{\rm bottom\,coat}$ . Thus a BARL could eliminate both the swing and notching problems in the lithographic process, so it is the most complete solution to obtaining a high resolution in deep-ultraviolet lithography.

Recently, amorphous carbon<sup>1,8</sup> and inorganic materials such as SiO<sub>x</sub>N<sub>x</sub>:H<sup>2</sup> and SiC<sup>9</sup> have been investigated as candidates for the BARL in deep-ultraviolet lithography. In this study, titanium oxide film was chosen as the BARL material in KrF excimer laser lithography, because it is expected that the optical constants (refractive index n and extinction coefficient k) of titanium oxide film at the wavelength of 248 nm are fitted to the optimum ARL condition that is discussed here.<sup>10</sup> In this study, the optical characterizations of titanium oxide films prepared by electron cyclotron resonance-plasma-enhanced chemical vapor deposition (ECR-PECVD)<sup>11,12</sup> are carried out. Moreover, an optical design simulation process is employed to find the optimum condition of zero reflectance between the PR and film and to calculate the reflectance in the fabricated films. Based on these results, the applicability of amorphous titanium oxide films as the BARL in KrF excimer laser lithography is discussed.

# 2. Experiment

The titanium oxide film was deposited in an ECR–PECVD system that could fabricate the dense and high-quality thin film without substrate heating and damage. The schematic diagram for the ECR–PECVD system used in the deposition of titanium oxide films is shown in Fig. 2. The titanium oxide films on (100) p-type Si substrates were deposited by using  $\text{Ti}(\text{O-}i\text{-}\text{C}_3\text{H}_7)_4$  (tetraisopropyltitanate, or TIPT) and  $\text{O}_2$  gas. TIPT is a suitable titanium precursor because it has an adequate vapor pressure even at room temperature.  $^{13,14}$   $\text{O}_2$  gas (5 N) to form the ECR

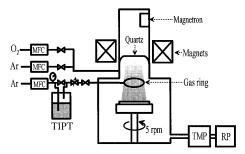


Fig. 2. Schematic diagram of the ECR-PECVD system used in the deposition of titanium oxide films: MFC, mass flow controller; TMP, turbomolecular pump; RP, rotary pump.

layer was introduced into the plasma chamber, and TIPT was introduced into the reaction chamber by means of a heated gas ring to protect condensation. The TIPT maintained to 65 °C was bubbled by Ar gas (5 N) at 10 sccm (where sccm denotes cubic centimeter per minute at STP), and a constant flow rate was achieved by controlling conductance with a metering valve. These excited oxygen and TIPT gases formed the titanium oxide film on the Si substrate located on a susceptor that could be rotated and moved up or down and heated up to maximum 700 °C by a resistant heater. The vacuum pump consisted of a turbomolecular pump and rotary pump. Pressures were measured in the reaction chamber and in the front of the rotary pump by an ionization gauge and a convectron gauge, respectively. Gases were introduced after evacuating the chamber and lines to the base pressure of  $10^{-5}$  Torr, and the deposition was done at a pressure of  $3 \times 10^{-3}$  Torr. The deposition characteristics were controlled by various O2 flow rates (5-30 sccm).

We performed a thin-film x-ray diffraction (Rikagu Inc., with a Cu target and Ni filter) analysis to investigate the crystallinity of the films under the conditions of tube voltage, 30 kV; tube current, 60 mA; incident beam angle, 3°; and scanning speed, 4°/min. The thickness, refractive index, and extinction coefficient at the 248-nm wavelength were evaluated by a spectroscopic ellipsometer (SOPRA SE ESVG). The reflectance at the interface between the photoresist and  ${\rm TiO}_x$  (BARL) film under the structure of PR/300-Å  ${\rm TiO}_x$ /Si or the W-Si substrate at the 248-nm wavelength with the matrix approach assuming a normal incident beam and a uniform and isotropic titanium oxide thin film has been calculated.

### 3. Results and Discussion

# A. Deposition of Titanium Oxide Film

The titanium oxide films were deposited with a range of  $O_2$  flow rates from 5 to 30 sccm at an ECR power of 300 W, a deposition temperature of room temperature, a TIPT bubbler temperature of 65 °C, and a working pressure of  $3 \times 10^{-3}$  Torr.

The x-ray diffraction analysis in Fig. 3 shows the films to be amorphous, which is important for optical homogeneity. All the films showed an almost con-

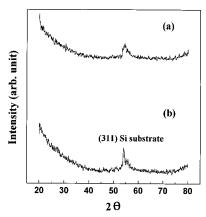


Fig. 3. X-ray diffraction analyses of  ${\rm TiO}_x$  films deposited at 300 W at room temperature, with an  ${\rm O}_2$  flow rate of (a) 10 sccm, (b) 20 sccm.

stant deposition rate of approximately 30 Å/min, irrespective of the variation of  $\mathrm{O}_2$  flow rate.

Figure 4 shows the variation of the optical constants (n and k) of the films at the 248-nm wavelength as a function of  $O_2$  flow rate in the range 5–30 sccm. The refractive index and extinction coefficient of the film, which are influenced by microstructure, compositional stoichiometry, and impurities in the

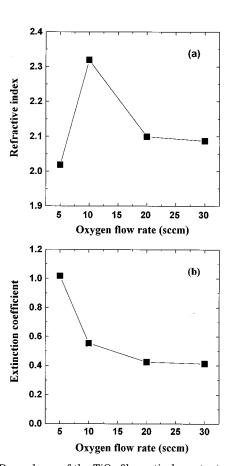


Fig. 4. Dependence of the  ${\rm TiO}_x$  film optical constants on the  ${\rm O}_2$  flow rate at the 248-nm wavelength (deposition condition is 300 W, room temperature): (a) refractive index, (b) extinction coefficient.

film, vary from 2.09 to 2.32 and from 1.02 to 0.41, respectively, depending on the  $O_2$  flow rate. The refractive index of the film deposited at the  $O_2$  flow rate of 10 sccm shows a maximum value of 2.32, and as the  $O_2$  flow rate increases to 30 sccm, it decreases because of the increase of oxygen incorporation in the film. The maximum point of the refractive index may be due to the mixing effects of plasma-forming efficiency variation by partial pressure variations of Ar and  $O_2$  gases and optimum  $O_2$  flow required to form the film. In contrast, the extinction coefficient decreases continuously as the  $O_2$  flow rate increases from 5 to 30 sccm, and the film deposited at an  $O_2$  flow rate of 5 sccm shows the highest extinction coefficient with the lowest refractive index.

# B. Calculation of Reflectance

The reflectance (R) at the interface between the PR and ARL for the exposure wavelength of 248 nm could be calculated by using the following  $2 \times 2$  characteristic matrix equation, assuming a normal incident beam on a uniform and isotropic thin film.<sup>15</sup>

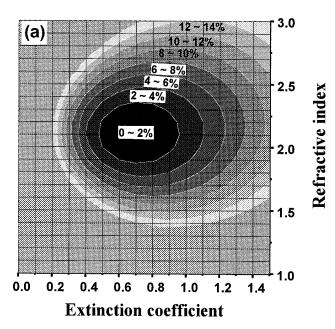
$$R = \left| \frac{(n_{\rm PR} - ik_{\rm PR})B - C}{(n_{\rm PR} - ik_{\rm PR})B + C} \right|^2, \tag{2}$$

$$\begin{bmatrix} B \\ C \end{bmatrix} = \begin{bmatrix} \cos \delta & \frac{i \sin \delta}{n_f - ik_f} \\ i(n_f - ik_f) \sin \delta & \cos \delta \end{bmatrix} \begin{bmatrix} 1 \\ (n_{\text{sub}} - ik_{\text{sub}}) \end{bmatrix}, \quad (3)$$

where  $n_{\rm PR}$  and  $k_{\rm PR}$  are optical constants of the PR, and  $n_f$  and  $k_f$  are optical constants of the ARL,  $n_{\rm sub}$  and  $k_{\rm sub}$  are optical constants of the substrate, B is the standardized electrical field at the incident boundary plane, C is the standardized magnetic field at the incident boundary plane, and  $\delta$  is the optical phase thickness equal to  $[2\pi(n_f-ik_f)d/\lambda]$ , where d is the physical film thickness and  $\lambda$  is 248 nm.

We considered the structures of the PR ( $\times$  P89131;  $n_{\rm PR}=1.8,\,k_{\rm PR}=0.011$  at 248 nm)/300-Å BARL/crystalline Si (c-Si) substrate ( $n_{\rm Si}=1.57,\,k_{\rm Si}=3.565$  at 248 nm) or W-Si substrate ( $n_{\rm W-Si}=1.763,\,k_{\rm W-Si}=1.763,\,k_{\rm W-Si}=1.763,\,k_{\rm W-Si}=1.763$ 2.546 at 248 nm) for the calculation of the reflectance between the PR and ARL. The result of a threedimensional simulated reflectance contour for the c-Si substrate is shown in Fig. 5(a). The reflectance goes to nearly zero at 248 nm when the ARL has approximately a refractive index of 2.11 and an extinction coefficient of 0.68, supposing a film thickness of 300 Å. Also, the result of the three-dimensional simulated reflectance contour for the W-Si substrate is presented in Fig. 5(b). The optimum condition of the ARL for the almost zero reflectance is when the refractive index and extinction coefficient are 2.05 and 0.59, respectively, at a film thickness of 300 A. These simulation results are consistent with other simulation studies.<sup>2,9</sup> Thus, the optical constants of the titanium oxide films prepared in this study are nearly matched with the optimum condition for the antireflective coating in KrF excimer laser lithography.

The calculated reflectance at the interface between



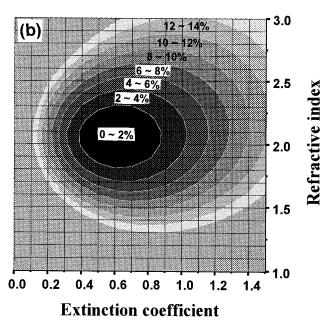


Fig. 5. Simulated reflectance contours at 248 nm: (a) PR( $n_{\rm PR}=1.8,~k_{\rm PR}=0.011$ )/300-Å ARL/c-Si ( $n_{\rm Si}=1.57,~k_{\rm Si}=3.565$ ); (b) PR( $n_{\rm PR}=1.8,~k_{\rm PR}=0.011$ )/300-Å ARL/W-Si ( $n_{\rm W-Si}=1.763,~k_{\rm W-Si}=2.546$ ).

the PR and the film for the structures of PR/300-Å  ${\rm TiO}_x$  film/c-Si and PR/300-Å  ${\rm TiO}_x$  film/W-Si, depending on the  ${\rm O}_2$  flow rate based on the optical parameters displayed in Fig. 4, is shown in Fig. 6. All the films in both structures exhibit a reflectance of less than 4%, which can be used as the antireflective coating in the practical lithographic process. Specifically, the reflectance can be reduced as low as 1.2% for the W-Si substrate when  ${\rm O}_2$  flow rate of 20 sccm is used.

The reflectance of the film can also be controlled by changing the film thickness. The simulated reflec-

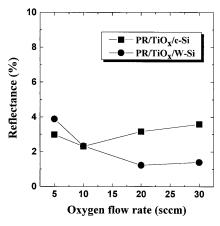


Fig. 6. Effect of  $O_2$  flow rate on the reflectance at 248 nm for two structures of  ${\rm TiO}_x$  films deposited at 300 W, room temperature (assuming a film thickness of 300 Å).

tance variations as a function of film thickness in both  $PR/TiO_x$  film/c-Si and  $PR/TiO_x$  film/W-Si structures, when it is supposed that optical constants are constant regardless of film thickness, are shown in Figs. 7(a) and 7(b), respectively. The reflectance of the film is strongly influenced by the film thickness. For the film thickness of less than 500 Å, which is preferred in the lithographic process, 0.9% reflectance of the film with the  $O_2$  flow rate of 10 sccm is obtained at the film thickness of 250 Å for the Si substrate, and 0.1% reflectance of the film with the  $O_2$  flow rate of 10 sccm is calculated at the film thickness of 225 Å for the W-Si substrate. Therefore the titanium oxide films prepared by ECR–PECVD can

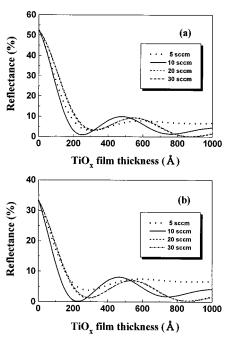


Fig. 7. Effect of the  $\mathrm{TiO}_x$  film thickness on the reflectance at 248 nm with various  $\mathrm{O}_2$  flow rates (deposition condition is 300 W, room temperature): (a)  $\mathrm{PR/TiO}_x/c\text{-Si}$  structure, (b)  $\mathrm{PR/TiO}_x/W\text{-Si}$  structure.

be used as the antireflective layer in KrF excimer laser lithography by controlling the film thickness as well as the process parameters.

# 4. Conclusions

The uniform amorphous titanium oxide thin films were deposited by using an ECR-PECVD system at different O<sub>2</sub> flow rates for use as the antireflective layer in KrF excimer laser lithography. The refractive indices and the extinction coefficients of the films at the 248-nm wavelength lay in the range from 2.09 to 2.32 and from 1.02 to 0.41, respectively. Supposing a film thickness of 300 A at the 248-nm wavelength, the optimum conditions for nearly zero reflectance are when there are an  $n_f$  of 2.11,  $k_f$  of 0.68 and an  $n_f$  of 2.05,  $k_f$  of 0.59 for Si and W-Si substrates, respectively. Thus these values are included in the range of optical constants of titanium oxide films with the selected O<sub>2</sub> flow rates. The calculated reflectance values at the interface between the PR and the film postulating the structures of PR/300-A TiO<sub>x</sub> film/c-Si and PR/300-Å TiO<sub>x</sub> film/W-Si were less than 4% and as low as 1.2%, respectively. Also, the reflectance is very sensitive to the film thickness, so it could be reduced to almost zero by controlling the film thickness. Therefore, the titanium oxide films prepared by ECR–PECVD are optically applicable to the bottom antireflective layer in KrF excimer laser lithography.

B-H. Jun thanks C-K. Hwangbo of Inha University for valuable discussion.

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