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Correlation Theorem for Two-Sided

Quaternion Fourier Transform

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Abstract

In this paper we establish correlation theorem for the two-sided quaternion Fourier transform (QFT). A consequence of the theorem is also investigated.

Keywords: quaternion Fourier transform, correlation

I. Introduction

The quaternion Fourier transform (QFT) is a nontrivial generalization of the classical Fourier transform (FT) using quaternion algebra. A number of already known and useful properties of this extended transform are generalizations of the corresponding properties of the FT with some modifications (see, for example, [1, 2, 3, 4, 5]). *One* of the *most* powerful *properties* of the QFT is the convolution theorem. Recently, in [6] authors proposed the convolution theorem for the two-sided QFT, which describes the relationship between the two-sided QFT and convolution of two quaternion function.

Therefore, the main objective of the present paper is to establish the correlation for the two-sided QFT, which is a generalization of correlation theorem of the classical FT. We find that the correlation theorem does not work well for the right-sided quaternion Fourier transform and left-sided quaternion Fourier transform.

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The quaternion algebra over \mathbb{R} , denoted by

$$\mathbb{H} = \{ q = q_0 + iq_1 + jq_2 + kq_3 \}, \qquad q_0, q_1, q_2, q_3 \in \mathbb{R},$$
 (1)

is an associative non-commutative four-dimensional algebra, which the quaternion units i, j, and k obey the following multiplication rules:

$$i^2 = j^2 = k^2 = -1$$
, $ij = -ji$, $ik = -ki$, $jk = -kj$, and $ijk = -1$. (2)

The quaternion conjugate is defined by

$$\bar{q} = q_0 - \boldsymbol{i}q_1 - \boldsymbol{j}q_2 - \boldsymbol{k}q_3, \tag{3}$$

which is an anti-involution, that is,

$$\bar{p} = p$$
, $\overline{p+q} = \bar{p} + \bar{q}$, $\overline{pq} = \bar{q}\bar{p}$. (4)

The norm of a quaternion is defined as

$$|q| = \sqrt{q\bar{q}} = \sqrt{q_0^2 + q_1^2 + q_2^2 + q_3^2}.$$
 (5)

It is not difficult to check that

$$|qp| = |p||q|, \quad \forall \ p, q \in \mathbb{H}. \tag{6}$$

We further get the inverse

$$q^{-1} = \frac{\overline{q}}{|q|^2}.$$

This fact shows that H is a skew field, that means, every nonzero element has a multiplicative inverse.

For the sake of further simplicity, we will use the real vector notations

$$\mathbf{x} = (x_1, x_2) \in \mathbb{R}^2$$
, $f(\mathbf{x}) = f(x_1, x_2)$, $f(\mathbf{\omega}) = f(\omega_1, \omega_2)$, and so on when there is confusion.

2. Main Results

In this section we introduce the definition of the two-sided QFT and establish the correlation of two quaternion-valued functions associated with the two-sided QFT.

Definition 2.1 (Two-sided QFT) Let f be in $\in L^2(\mathbb{R}^2; \mathbb{H})$. The two-sided QFT of the quaternion function f is the transform given by the integral

$$\mathcal{F}_{q}\lbrace f\rbrace(\boldsymbol{\omega}) = \int_{\mathbb{R}^{2}} e^{-i\omega_{1}x_{1}} f(\boldsymbol{x}) e^{-j\omega_{2}x_{2}} d^{2}\boldsymbol{x}. \tag{7}$$

Theorem 2.1 (Inverse two-sided QFT) Suppose that $f \in L^2(\mathbb{R}^2; \mathbb{H})$ and $\mathcal{F}_q\{f\} \in L^1(\mathbb{R}^2; \mathbb{H})$. Then the two-sided QFT is an invertible transform and its inverse is given by

$$f(\mathbf{x}) = \frac{1}{(2\pi)^2} \int_{\mathbb{R}^2} e^{i\omega_1 x_1} \mathcal{F}_q\{f\}(\boldsymbol{\omega}) e^{j\omega_2 x_2} d^2 \boldsymbol{\omega}.$$
 (8)

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Definition 2.2 (Quaternion Correlation) Suppose f and g are quaternion functions in $L^2(\mathbb{R}^2; \mathbb{H})$. The quaternion correlation of the functions is given by

$$(f \otimes g)(\mathbf{x}) = \int_{\mathbb{R}^2} f(\mathbf{x} + \mathbf{y}) \overline{g(\mathbf{y})} d^2 \mathbf{y}.$$
 (9)

The following theorem is the main result of this paper, which describes how the two-sided QFT behaves under correlations.

Theorem 2.2 Let $f, g \in L^2(\mathbb{R}^2; \mathbb{H})$ be two quaternion-valued functions. The two-sided QFT of correlation of f and g takes the form $\mathcal{F}_a\{f \otimes g\}(\omega)$

$$\begin{split} &= \Big(\mathcal{F}_q\{f_0\}(\boldsymbol{\omega}) + \boldsymbol{i}\mathcal{F}_q\{f_1\}(\boldsymbol{\omega})\Big) \Big(\mathcal{F}_q\{g_0\}(-\boldsymbol{\omega}) - \boldsymbol{j}\mathcal{F}\{g_2\}(-\omega_1,\omega_2)\Big) \\ &\quad + \Big(\mathcal{F}_q\{f_0\}(\omega_1,-\omega_2) + \boldsymbol{i}\mathcal{F}_q\{f_1\}(\omega_1,-\omega_2)\Big) \Big(-\boldsymbol{i}\mathcal{F}_q\{g_1\}(-\boldsymbol{\omega}) - \boldsymbol{k}\mathcal{F}\{g_3\}(\omega_1,-\omega_2)\Big) \\ &\quad + \Big(\boldsymbol{j}\mathcal{F}_q\{f_2\}(-\omega_1,\omega_2) + \boldsymbol{k}\mathcal{F}_q\{f_3\}(-\omega_1,\omega_2)\Big) \Big(\mathcal{F}_q\{g_0\}(\omega_1,-\omega_2) - \boldsymbol{j}\mathcal{F}\{g_2\}(-\boldsymbol{\omega})\Big) \\ &\quad + \Big(\boldsymbol{j}\mathcal{F}_q\{f_2\}(-\boldsymbol{\omega}) + \boldsymbol{k}\mathcal{F}_q\{f_3\}(-\boldsymbol{\omega})\Big) \Big(-\boldsymbol{i}\mathcal{F}_q\{g_1\}(\omega_1,-\omega_2) - \boldsymbol{k}\mathcal{F}\{g_3\}(-\boldsymbol{\omega})\Big). \end{split}$$

Proof. Applying the two-sided QFT definition gives $\mathcal{F}_a\{f \otimes g\}(\boldsymbol{\omega})$

$$= \int_{\mathbb{R}^2} e^{-i\omega_1 x_1} \left(\int_{\mathbb{R}^2} f(\boldsymbol{x} + \boldsymbol{y}) \, \overline{g(\boldsymbol{y})} \, d^2 \boldsymbol{y} \right) e^{-j\omega_2 x_2} \, d^2 \boldsymbol{x}.$$

By inserting the change of variables z = x + y to the above expression we immediately obtain

$$\begin{split} & \mathcal{F}_{q}\{f \otimes g\}(\boldsymbol{\omega}) \\ & = \int_{\mathbb{R}^{2}} e^{-i\omega_{1}(z_{1}-y_{1})} \left(\int_{\mathbb{R}^{2}} f(\mathbf{z}) \, \overline{g(\mathbf{y})} \, d^{2}\mathbf{y} \right) e^{-j\omega_{2}(z_{2}-y_{2})} \, d^{2}\mathbf{z} \\ & = \int_{\mathbb{R}^{2}} \int_{\mathbb{R}^{2}} e^{-i\omega_{1}(z_{1}-y_{1})} \left(\left(f_{0}(\mathbf{z}) + if_{1}(\mathbf{z}) \right) + \left(jf_{2}(\mathbf{z}) + kf_{3}(\mathbf{z}) \right) \right) \\ & \left(\left(g_{0}(\mathbf{y}) - \mathbf{j}g_{2}(\mathbf{y}) \right) + \left(-ig_{1}(\mathbf{y}) - kg_{3}(\mathbf{y}) \right) \right) e^{-j\omega_{2}(z_{2}-y_{2})} \, d^{2}\mathbf{y} \, d^{2}\mathbf{z} \\ & = \int_{\mathbb{R}^{2}} \int_{\mathbb{R}^{2}} e^{-i\omega_{1}(z_{1}-y_{1})} \left(f_{0}(\mathbf{z}) + if_{1}(\mathbf{z}) \right) \left(g_{0}(\mathbf{y}) - \mathbf{j}g_{2}(\mathbf{y}) \right) e^{-j\omega_{2}(z_{2}-y_{2})} \, d^{2}\mathbf{y} \, d^{2}\mathbf{z} \\ & + \int_{\mathbb{R}^{2}} \int_{\mathbb{R}^{2}} e^{-i\omega_{1}(z_{1}-y_{1})} \left(f_{0}(\mathbf{z}) + if_{1}(\mathbf{z}) \right) \left(-ig_{1}(\mathbf{y}) - kg_{3}(\mathbf{y}) \right) e^{-j\omega_{2}(z_{2}-y_{2})} \, d^{2}\mathbf{y} \, d^{2}\mathbf{z} \\ & + \int_{\mathbb{R}^{2}} \int_{\mathbb{R}^{2}} e^{-i\omega_{1}(z_{1}-y_{1})} \left(\mathbf{j}f_{2}(\mathbf{z}) + kf_{3}(\mathbf{z}) \right) \left(g_{0}(\mathbf{y}) - \mathbf{j}g_{2}(\mathbf{y}) \right) e^{-j\omega_{2}(z_{2}-y_{2})} d^{2}\mathbf{y} \, d^{2}\mathbf{z} \end{split}$$

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$$\begin{split} &+ \int\limits_{\mathbb{R}^{2}} \int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}(z_{1}-y_{1})} \left(\left(jf_{2}(\mathbf{z}) + kf_{3}(\mathbf{z}) \right) \left(-ig_{1}(\mathbf{y}) - kg_{3}(\mathbf{y}) \right) \right) e^{-j\omega_{2}(z_{2}-y_{2})} \ d^{2}\mathbf{y} \ d^{2}\mathbf{z} \\ &= \int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}z_{1}} \left(f_{0}(\mathbf{z}) + if_{1}(\mathbf{z}) \right) \\ &\left(\int\limits_{\mathbb{R}^{2}} e^{i\omega_{1}y_{1}} \ g_{0}(\mathbf{y}) e^{j\omega_{2}y_{2}} \ d^{2}\mathbf{y} - \mathbf{j} \int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}y_{1}} g_{2}(\mathbf{y}) e^{j\omega_{2}y_{2}} \ d^{2}\mathbf{y} \right) e^{-j\omega_{2}z_{2}} d^{2}\mathbf{z} \\ &\left(-i \int\limits_{\mathbb{R}^{2}} e^{i\omega_{1}y_{1}} g_{1}(\mathbf{y}) e^{j\omega_{2}y_{2}} \ d^{2}\mathbf{y} - \mathbf{k} \int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}y_{1}} g_{3}(\mathbf{y}) e^{j\omega_{2}y_{2}} \ d^{2}\mathbf{y} \right) e^{-j\omega_{2}z_{2}} d^{2}\mathbf{z} \\ &+ \int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}z_{1}} \left(jf_{2}(\mathbf{z}) + kf_{3}(\mathbf{z}) \right) \\ &\left(\int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}y_{1}} g_{0}(\mathbf{y}) e^{j\omega_{2}y_{2}} \ d^{2}\mathbf{y} - \mathbf{j} \int\limits_{\mathbb{R}^{2}} e^{i\omega_{1}y_{1}} g_{2}(\mathbf{y}) e^{j\omega_{2}y_{2}} \ d^{2}\mathbf{y} \right) e^{-j\omega_{2}z_{2}} d^{2}\mathbf{z} \\ &+ \int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}z_{1}} \left(jf_{2}(\mathbf{z}) + kf_{3}(\mathbf{z}) \right) \\ &\left(-i \int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}y_{1}} g_{1}(\mathbf{y}) e^{j\omega_{2}y_{2}} \ d^{2}\mathbf{y} - \mathbf{k} \int\limits_{\mathbb{R}^{2}} e^{i\omega_{1}y_{1}} g_{3}(\mathbf{y}) e^{j\omega_{2}y_{2}} \ d^{2}\mathbf{y} \right) e^{-j\omega_{2}z_{2}} d^{2}\mathbf{z} \\ &= \int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}z_{1}} \left(f_{0}(\mathbf{z}) + if_{1}(\mathbf{z}) \right) \left(\mathcal{F}_{q}\{g_{0}\}(-\boldsymbol{\omega}) - j\mathcal{F}_{q}\{g_{2}\}(\omega_{1}, -\omega_{2}) \right) e^{-j\omega_{2}z_{2}} d^{2}\mathbf{z} \\ &+ \int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}z_{1}} \left(jf_{2}(\mathbf{z}) + kf_{3}(\mathbf{z}) \right) \left(\mathcal{F}_{q}\{g_{0}\}(\omega_{1}, -\omega_{2}) - j\mathcal{F}_{q}\{g_{2}\}(-\boldsymbol{\omega}) \right) e^{-j\omega_{2}z_{2}} d^{2}\mathbf{z} \\ &+ \int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}z_{1}} \left(jf_{2}(\mathbf{z}) + kf_{3}(\mathbf{z}) \right) \left(-i\mathcal{F}_{q}\{g_{1}\}(\omega_{1}, -\omega_{2}) + k\mathcal{F}_{q}\{g_{1}\}(-\boldsymbol{\omega}) \right) e^{-j\omega_{2}z_{2}} d^{2}\mathbf{z} \\ &= \int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}z_{1}} \left(jf_{2}(\mathbf{z}) + kf_{3}(\mathbf{z}) \right) \left(-i\mathcal{F}_{q}\{g_{1}\}(\omega_{1}, -\omega_{2}) + k\mathcal{F}_{q}\{g_{1}\}(-\boldsymbol{\omega}) \right) e^{-j\omega_{2}z_{2}} d^{2}\mathbf{z} \\ &= \int\limits_{\mathbb{R}^{2}} e^{-i\omega_{1}z_{1}} \left(jf_{2}(\mathbf{z}) + kf_{3}(\mathbf{z}) \right) \left(-i\mathcal{F}_{q}\{g_{1}\}(\omega_{1}, -\omega_{2}) + k\mathcal{F}_{q}\{g_{1}\}(-\boldsymbol{\omega}) \right) e^{-j\omega_{2}z_{2}} d^{2}\mathbf{z} \end{aligned}$$

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$$\begin{split} &+\int\limits_{\mathbb{R}^2}e^{-i\omega_1z_1}\big(f_0(\mathbf{z})+\mathbf{i}f_1(\mathbf{z})\big)\ e^{\mathbf{j}\omega_2z_2}d^2\mathbf{z}\,\Big(-\mathbf{i}\mathcal{F}_q\{g_1\}(-\boldsymbol{\omega})-\mathbf{k}\mathcal{F}_q\{g_3\}(\omega_1,-\omega_2)\Big)\\ &+\int\limits_{\mathbb{R}^2}e^{-i\omega_1z_1}\left(\mathbf{j}f_2(\mathbf{z})+\mathbf{k}f_3(\mathbf{z})\right)e^{-j\omega_2z_2}d^2\mathbf{z}\,\Big(\mathcal{F}_q\{g_0\}(\omega_1,-\omega_2)-\mathbf{j}\mathcal{F}_q\{g_2\}(-\boldsymbol{\omega})\Big)\\ &+\int\limits_{\mathbb{R}^2}e^{-i\omega_1z_1}\big(\mathbf{j}f_2(\mathbf{z})+\mathbf{k}f_3(\mathbf{z})\big)\ e^{\mathbf{j}\omega_2z_2}d^2\mathbf{z}\,\Big(-\mathbf{i}\mathcal{F}_q\{g_1\}(\omega_1,-\omega_2)-\mathbf{k}\mathcal{F}_q\{g_3\}(-\boldsymbol{\omega})\Big)\\ &=\Big(\mathcal{F}_q\{f_0\}(\boldsymbol{\omega})+\mathbf{i}\mathcal{F}_q\{f_1\}(\boldsymbol{\omega})\Big)\Big(\mathcal{F}_q\{g_0\}(-\boldsymbol{\omega})-\mathbf{j}\mathcal{F}\{g_2\}(-\omega_1,\omega_2)\Big)\\ &+\Big(\mathcal{F}_q\{f_0\}(\omega_1,-\omega_2)+\mathbf{i}\mathcal{F}_q\{f_1\}(\omega_1,-\omega_2)\Big)\Big(-\mathbf{i}\mathcal{F}_q\{g_1\}(-\boldsymbol{\omega})-\mathbf{k}\mathcal{F}\{g_3\}(\omega_1,-\omega_2)\Big)\\ &+\Big(\mathbf{j}\mathcal{F}_q\{f_2\}(-\omega_1,\omega_2)+\mathbf{k}\mathcal{F}_q\{f_3\}(-\omega_1,\omega_2)\Big)\Big(\mathcal{F}_q\{g_0\}(\omega_1,-\omega_2)-\mathbf{j}\mathcal{F}\{g_2\}(-\boldsymbol{\omega})\Big)\\ &+\Big(\mathbf{j}\mathcal{F}_q\{f_2\}(-\boldsymbol{\omega})+\mathbf{k}\mathcal{F}_q\{f_3\}(-\boldsymbol{\omega})\Big)\Big(-\mathbf{i}\mathcal{F}_q\{g_1\}(\omega_1,-\omega_2)-\mathbf{k}\mathcal{F}\{g_3\}(-\boldsymbol{\omega})\Big). \end{split}$$

As a consequence of the above theorem, we immediately obtain

Lemma 2.3 Given any two quaternion-valued functions $f, g \in L^2(\mathbb{R}^2; \mathbb{H})$. If we assume that $\mathcal{F}_q\{f\}(\boldsymbol{\omega}) \in L^2(\mathbb{R}^2; \mathbb{R})$, then

$$\begin{split} \mathcal{F}_q\{f\otimes g\}(\pmb\omega) &= \mathcal{F}_q\left\{\,f\right\}(\omega_1,-\omega_2)\left(\mathcal{F}_q\left\{g_0\right\}(-\pmb\omega) - \pmb i\,\mathcal{F}_q\left\{g_1\right\}(-\pmb\omega)\right) \\ &- \mathcal{F}_q\left\{f\right\}(\omega_1,-\omega_2)\left(\pmb j\mathcal{F}_q\{g_2\}(\omega_1,-\omega_2) + \pmb k\mathcal{F}_q\{g_3\}(-\omega_1,\omega_2)\right). \end{split}$$

Proof. Straightforward computations show that $\mathcal{F}_q\{f \otimes g\}(\omega)$

$$\begin{split} &= \int\limits_{\mathbb{R}^2} e^{-i\omega_1 x_1} \left(\int\limits_{\mathbb{R}^2} f(\mathbf{x} + \mathbf{y}) \overline{g(\mathbf{y})} \ d^2 \mathbf{y} \right) e^{-j\omega_2 x_2} d^2 \mathbf{x} \\ &= \int\limits_{\mathbb{R}^2} \int\limits_{\mathbb{R}^2} e^{-i\omega_1 (z_1 - y_1)} \left(f(\mathbf{z}) \left(g_0(\mathbf{y}) - i g_1(\mathbf{y}) \right) - f(\mathbf{z}) \left(j g_2(\mathbf{y}) + k g_3(\mathbf{y}) \right) \right) e^{-j\omega_2 (z_2 - y_2)} d^2 \mathbf{y} \ d^2 \mathbf{z} \\ &= \int\limits_{\mathbb{R}^2} e^{-i\omega_1 (z_1 - y_1)} f(\mathbf{z}) \int\limits_{\mathbb{R}^2} \left(g_0(\mathbf{y}) - i g_1(\mathbf{y}) \right) e^{j\omega_2 y_2} d^2 \mathbf{y} \ e^{-j\omega_2 z_2} \ d^2 \mathbf{z} \\ &- \int\limits_{\mathbb{R}^2} e^{-i\omega_1 (z_1 - y_1)} f(\mathbf{z}) \int\limits_{\mathbb{R}^2} \left(j g_2(\mathbf{y}) + k g_3 \mathbf{y} \right) e^{j\omega_2 y_2} d^2 \mathbf{y} \ e^{-j\omega_2 z_2} \ d^2 \mathbf{z} \\ &= \mathcal{F}_q \left\{ f \right\} (\omega_1, -\omega_2) \int\limits_{\mathbb{R}^2} e^{-i\omega_1 y_1} \left(g_0(\mathbf{y}) - i g_1(\mathbf{y}) \right) e^{-j\omega_2 y_2} d^2 \mathbf{y} \end{split}$$

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$$-\mathcal{F}_q\left\{f\right\}(\omega_1,-\omega_2)\int\limits_{\mathbb{R}^2}e^{-i\omega_1y_1}\big(jg_2(y)+kg_3(y)\big)\;e^{-j\omega_2y_2}\;d^2y.$$

We further obtain

$$\mathcal{F}_a\{f\otimes g\}(\boldsymbol{\omega})$$

$$\begin{split} &= \mathcal{F}_{q} \left\{ f \right\} (\omega_{1}, -\omega_{2}) \int_{\mathbb{R}^{2}} e^{i\omega_{1}y_{1}} \left(g_{0}(\mathbf{y}) - ig_{1}(\mathbf{y}) \right) e^{j\omega_{2}y_{2}} d^{2}\mathbf{y} \\ &- \mathcal{F}_{q} \left\{ f \right\} (\omega_{1}, -\omega_{2}) \int_{\mathbb{R}^{2}} e^{i\omega_{1}y_{1}} \left(jg_{2}(\mathbf{y}) + \mathbf{k}g_{3}(\mathbf{y}) \right) e^{j\omega_{2}y_{2}} d^{2}\mathbf{y} \\ &= \mathcal{F}_{q} \left\{ f \right\} (\omega_{1}, -\omega_{2}) \left(\int_{\mathbb{R}^{2}} e^{i\omega_{1}y_{1}} g_{0}(\mathbf{y}) e^{j\omega_{2}y_{2}} d^{2}\mathbf{y} - \int_{\mathbb{R}^{2}} e^{i\omega_{1}y_{1}} ig_{1}(\mathbf{y}) e^{j\omega_{2}y_{2}} d^{2}\mathbf{y} \right) \\ &- \mathcal{F}_{q} \left\{ f \right\} (\omega_{1}, -\omega_{2}) \left(\int_{\mathbb{R}^{2}} e^{i\omega_{1}y_{1}} jg_{2}(\mathbf{y}) e^{j\omega_{2}y_{2}} d^{2}\mathbf{y} + \int_{\mathbb{R}^{2}} e^{i\omega_{1}y_{1}} \mathbf{k} g_{3}(\mathbf{y}) e^{j\omega_{2}y_{2}} d^{2}\mathbf{y} \right) \\ &= \mathcal{F}_{q} \left\{ f \right\} (\omega_{1}, -\omega_{2}) \left(\mathcal{F}_{q} \left\{ g_{0} \right\} (-\boldsymbol{\omega}) - i \mathcal{F}_{q} \left\{ g_{1} \right\} (-\boldsymbol{\omega}) \right) \\ &- \mathcal{F}_{q} \left\{ f \right\} (\omega_{1}, -\omega_{2}) \left(j\mathcal{F}_{q} \left\{ g_{2} \right\} (\omega_{1}, -\omega_{2}) + \mathbf{k}\mathcal{F}_{q} \left\{ g_{3} \right\} (-\omega_{1}, \omega_{2}) \right). \end{split}$$

which was to be proved.

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