# OSCILLATION CRITERIA FOR SECOND ORDER DIFFERENTIAL EQUATIONS WITH DAMPING 

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#### Abstract

New oscillation criteria are given for second order nonlinear ordinary differential equations with alternating coefficients. The results involve a condition obtained by Kamenev for linear differential equations. The obtained criterion for superlinear differential equations is a complement of the work established by Kwong and Wong, and Wong and Philos, for sublinear differential equations and by Yan for linear differential equations.


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## 1. Introduction

This paper deals with the problem of oscillation of second order ordinary differential equations of the form

$$
\begin{equation*}
\left(a(t) x^{\bullet}(t)\right)^{\bullet}+p(t) x^{\bullet}(t)+q(t)|x(t)|^{\lambda} \operatorname{sgn} x(t)=0, \quad \lambda>0, \quad\left(^{\bullet}=\frac{d}{d t}\right) \tag{1}
\end{equation*}
$$

where $a, p, q:\left[t_{0}, \infty\right) \rightarrow R=(-\infty, \infty)$ are continuous and $a(t)>0$ for $t \geq t_{0}>0$.

Only such solutions, $x$, of equation (1) which exist on some interval $\left[t_{x}, \infty\right), t_{x} \geq t_{0}>0$, are considered. A solution of equation (1) is said to be oscillatory if it has arbitrarily large zeros; otherwise, it is said to be
nonoscillatory. Equation (1) is called oscillatory if all of its solutions are oscillatory.

The oscillation problem for second order nonlinear ordinary differential equations with alternating coefficients has been investigated by many authors. Some of the more important and useful tests involve the average behavior of the integral of the alternating coefficients. These tests have been motivated by the averaging criterion of Kamenev [5] and its generalizations. For such averaging techniques in second order nonlinear oscillation, we choose to refer to the paper of Butler [1], Grace and Lalli [2]-[4], Kanemev [5], Kwong and Wong [6], Philos [7], [8], Wong [9], [10], Yan [11] and Yen [12].

Several years ago, Kwong and Wong [6, Theorem 1] obtained an interesting criterion for the oscillation of equation (1) with $a(t)=1, p(t)=0$ and $0<\lambda<1$. This criterion has been extended by Philos [8] to more general equations of the form

$$
x^{\bullet \bullet}(t)+q(t) f(x(t))=0,
$$

where $q:\left[t_{0}, \infty\right) \rightarrow R$ and $f: R \rightarrow R$ are continuous, $x f(x)>0$ for $x \neq 0, f$ is nondecreasing and $\int_{ \pm 0} d u / f(u)<\infty$.

Recently, Yan [11, Theorem 2] proved that Kwong and Wong's theorem in [6] with averaging condition of Kamenev's type [5] remains valid for equation (1) when $\lambda=1$. The results in [6], [8], and [11] are not applicable to equation (1) with $\lambda>1$.

Therefore, the main purpose of this paper is to establish new criteria for the oscillation of equation (1) by using an averaging condition of the type introduced by Kamenev [5]. The result obtained for the superlinear case, that is, for $\lambda>1$, is a continuation of the work done by Kwong and Wong [6], Philos [8] and Yan [11]. We also mention that the results of this paper for the sublinear case, that is, for $0<\lambda<1$, are independent of those in [6], [8] and [11], but that they are similar to that of Yan [11] when $\lambda=1$.

## 2. Main results

Theorem 1. Suppose that there exist a differentiable function $\rho:\left[t_{0}, \infty\right) \rightarrow$ $(0, \infty)$ and a continuous function $\phi:\left[t_{0}, \infty\right) \rightarrow R$ such that

$$
\begin{align*}
& \lim _{t \rightarrow \infty} \operatorname{\operatorname {nf}} \int_{t_{0}}^{t} \rho(s) q(s) d s>-\infty  \tag{2}\\
& a(t) \rho^{\bullet}(t)-p(t) \rho(t)=\gamma(t) \geq 0 \text { and } \gamma^{\bullet}(t) \leq 0 \quad \text { for } t \geq t_{0}  \tag{3}\\
& \int^{\infty} \eta(s) d s=\infty, \quad \text { where } \eta(t)=\frac{1}{a(t) \rho(t)} \tag{4}
\end{align*}
$$

and for every constant $\theta>0$ and $\alpha \in(1, \infty)$

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \frac{1}{t^{\alpha}} \int_{t_{0}}^{t}(t-u)^{\alpha} \rho(u) q(u) d u<\infty \tag{5}
\end{equation*}
$$

$$
\begin{align*}
& \lim _{t \rightarrow \infty} \inf \frac{1}{t^{\alpha}} \int_{s}^{t}(t-u)^{\alpha-2}\left[(t-u)^{2} \rho(u) q(u)\right.  \tag{6}\\
& \quad-\frac{\theta}{4 \lambda} \frac{1}{v(u)}[\gamma(u) \eta(u)(t-u)-\alpha]^{2} d u>\phi(s)
\end{align*}
$$

and

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \int_{t_{0}}^{t} v(s) \phi_{+}^{2}(s) d s=\infty \tag{7}
\end{equation*}
$$

where $v(t)=\eta(t) / \int_{t_{0}}^{t} \eta(s) d s, \phi_{+}(t)=\max \{\phi(t), 0\}$.
Then equation (1) is oscillatory for all $\lambda>1$.
Proof. Let $x(t)$ be a nonoscillatory solution of equation (1). Without loss of generality, we assume that $x(t) \neq 0$ for $t \geq t_{0}$. Furthermore, we suppose that $x(t)>0$ for $t \geq t_{0}$, since the substitution $u=-x$ transforms equation (1) into an equation of the same form subject to the assumptions of the theorem. Now, we define

$$
W(t)=p(t) \frac{a(t) x^{\bullet}(t)}{x^{\lambda}(t)} \quad \text { for } t \geq t_{0} .
$$

Then it follows from equation (1) that

$$
\begin{equation*}
W^{\bullet}(t)=-\rho(t) q(t)+\gamma(t) \frac{x^{\bullet}(t)}{x^{\lambda}(t)}-\lambda a(t) \rho(t) \frac{\dot{x}^{2}(t)}{x^{\lambda+1}(t)}, \quad t \geq t_{0}, \tag{8}
\end{equation*}
$$

and consequently

$$
\begin{align*}
a(t) \rho(t) \frac{x^{\bullet}(t)}{x^{\lambda}(t)}= & c_{1}-\int_{t_{0}}^{t} \rho(s) q(s) d s+\int_{t_{0}}^{t} \gamma(s) \frac{x^{\bullet}(s)}{x^{\lambda}(s)} d s  \tag{9}\\
& -\int_{t_{0}}^{t} \lambda a(s) \rho(s)\left(\frac{x^{\bullet}(s)}{x^{\beta}(s)}\right)^{2} d s,
\end{align*}
$$

where $\beta=(\lambda+1) / 2$ and $c_{1}=a\left(t_{0}\right) \rho\left(t_{0}\right) x^{\bullet}\left(t_{0}\right) / x^{\lambda}\left(t_{0}\right)$. By the Bonnet theorem, for any $t \geq t_{0}$ there exists $\xi \in\left[t_{0}, t\right]$ such that

$$
\begin{aligned}
\int_{t_{0}}^{t} \gamma(s) \frac{x^{\bullet}(s)}{x^{\lambda}(s)} d s & =\gamma\left(t_{0}\right) \int_{t_{0}}^{\xi} \frac{x^{\bullet}(s)}{x^{\lambda}(s)} d s=\gamma\left(t_{0}\right) \int_{t_{0}}^{x(\xi)} u^{-\lambda} d u \\
& =\frac{\gamma\left(t_{0}\right)}{\lambda-1}\left[x^{1-\lambda}\left(t_{0}\right)-x^{1-\lambda}(\xi)\right] \leq \frac{\gamma\left(t_{0}\right)}{\lambda-1} x^{1-\lambda}\left(t_{0}\right)=M_{1} .
\end{aligned}
$$

Thus, for $t \geq t_{0}$, we get

$$
\begin{equation*}
a(t) \rho(t) \frac{x^{\bullet}(t)}{x^{\lambda}(t)}+\lambda \int_{t_{0}}^{t} a(s) \rho(s)\left(\frac{x^{\bullet}(s)}{x^{\beta}(s)}\right)^{2} d s+\int_{t_{0}}^{t} \rho(s) q(s) d s \leq L \tag{10}
\end{equation*}
$$

or

$$
\begin{equation*}
W(t)+\lambda \int_{t_{0}}^{t} \frac{1}{a(s) \rho(s)} x^{2(\beta-1)}(s) W^{2}(s) d s+\int_{t_{0}}^{t} \rho(s) q(s) d s \leq L \tag{11}
\end{equation*}
$$

where $L=c_{1}+M_{1}$.
Next, we consider the following three cases for the behavior of $x^{\bullet}$.
Case 1. $x^{\bullet}$ is oscillatory. Then there exists a sequence $\left\{t_{m}\right\}_{m=1,2, \ldots}$ in $\left[t_{0}, \infty\right)$ with $\lim _{m \rightarrow \infty} t_{m}=\infty$ and such that $x^{\bullet}\left(t_{m}\right)=0(m=1,2, \ldots)$. Thus (8) gives

$$
\int_{t_{0}}^{t_{m}} \lambda a(s) \rho(s)\left(\frac{x^{\bullet}(s)}{x^{\beta}(s)}\right)^{2} d s \leq L-\int_{t_{0}}^{t_{m}} \rho(s) q(s) d s \quad(m=1,2, \ldots),
$$

and hence, by taking into account (2), we conclude that

$$
\begin{equation*}
\int_{t_{0}}^{\infty} a(s) \rho(s)\left(\frac{x^{\bullet}(s)}{x^{\beta}(s)}\right)^{2} d s<\infty \tag{12}
\end{equation*}
$$

So, for some positive constant $N$, we have

$$
\int_{t_{0}}^{t} a(s) \rho(s)\left(\frac{x^{\bullet}(s)}{x^{\beta}(s)}\right)^{2} d s \leq N \quad \text { for } t \geq t_{0}
$$

By the Schwarz inequality

$$
\begin{aligned}
\left|\int_{t_{0}}^{t}\left(\frac{x^{\bullet}(s)}{x^{\beta}(s)}\right) d s\right|^{2} & \leq\left(\int_{t_{0}}^{t} \frac{d s}{a(s) \rho(s)}\right)\left(\int_{t_{0}}^{t} a(s) \rho(s)\left(\frac{x^{\bullet}(s)}{x^{\beta}(s)}\right)^{2} d s\right) \\
& \leq N \int_{t_{0}}^{t} \frac{d s}{a(s) \rho(s)}=N \int_{t_{0}}^{t} \eta(s) d s,
\end{aligned}
$$

or

$$
\left|x^{1-\beta}(t)-x^{1-\beta}\left(t_{0}\right)\right| \leq|1-\beta|\left(N \int_{t_{0}}^{t} \eta(s) d s\right)^{1 / 2}
$$

There exists $t_{1} \geq t_{0}$ and a constant $M>0$ so that

$$
\left|x^{1-\beta}(t)\right| \leq M\left(\int_{t_{0}}^{t} \eta(s) d s\right)^{1 / 2} \quad \text { for all } t \geq t_{1}
$$

Using (13) in (8) we get

$$
\begin{equation*}
W^{\bullet}(t) \leq-\rho(t) q(t)+\frac{\gamma(t)}{a(t) \rho(t)} W(t)-\frac{\lambda}{M^{2}} v(t) W^{2}(t), \quad t \geq t_{1} \tag{14}
\end{equation*}
$$

and consequently, for all $t>s \geq t_{1}$

$$
\begin{aligned}
\int_{s}^{t}(t-u)^{\alpha} W^{\bullet}(u) d u & \leq-\int_{s}^{t}(t-u)^{\alpha} \rho(u) q(u) d u \\
& -\int_{s}^{t}(t-u)^{\alpha}\left[\frac{\lambda}{M^{2}} v(u) W^{2}(u)-\gamma(u) \eta(u) W(u)\right] d u
\end{aligned}
$$

Since

$$
\int_{s}^{t}(t-u)^{\alpha} W^{\bullet}(u) d u=-(t-s)^{\alpha} W(s)+\alpha \int_{s}^{t}(t-u)^{\alpha-1} W(u) d u
$$

we obtain that

$$
\begin{align*}
& \int_{s}^{t}(t-u)^{\alpha} \rho(u) q(u) d u  \tag{15}\\
& \leq(t-s)^{\alpha} W(s)-\int_{s}^{t} {\left[(t-u)^{\alpha} \frac{\lambda}{M^{2}} v(u) W^{2}(u)\right.} \\
&\left.-(t-u)^{\alpha-1}[\gamma(u) \eta(u)(t-u)-\alpha] W(u)\right] d u
\end{align*}
$$

and hence

$$
\begin{align*}
& \int_{s}^{t}\left[(t-u)^{\alpha} \rho(u) q(u)-\frac{m^{2}}{4 \lambda} \frac{(t-u)^{\alpha-2}}{v(u)}(\gamma(u) \eta(u)(t-u)-\alpha)^{2}\right] d u  \tag{16}\\
& \leq(t-s)^{\alpha} W(s)-\int_{s}^{t}\left[\frac{1}{M}(t-u)^{\alpha / 2} \sqrt{\lambda v(u)} W(u)\right. \\
& \left.-\frac{M(t-u)^{\alpha / 2-1}(\gamma(u) \eta(u)(t-u)-\alpha)}{2 \sqrt{\lambda v(u)}}\right]^{2} d u \\
& \quad \leq(t-s)^{\alpha} W(s), \quad s \geq t_{1}
\end{align*}
$$

Now we proceed in a way similar to that in the proof of [11, Theorem 2]. Dividing (16) by $t^{\alpha}$ and taking the lower limit as $t \rightarrow \infty$, we obtain $\phi(s) \leq$ $W(s), s \geq t_{1}$, which implies that

$$
\begin{equation*}
\phi_{+}^{2}(s) \leq W^{2}(s) \tag{17}
\end{equation*}
$$

We define functions

$$
y(t)=t^{-\alpha} \int_{t_{1}}^{t}[\gamma(u) \eta(u)(t-u)-\alpha](t-u)^{\alpha-1} W(u) d u
$$

and

$$
z(t)=t^{-\alpha} \int_{t_{1}}^{t} \frac{\lambda}{M^{2}}(t-u)^{\alpha} v(u) W^{2}(u) d u, \quad t \geq t_{1}
$$

From (15), we get

$$
\begin{equation*}
y(t)=z(t) \leq t^{-\alpha}\left(t-t_{1}\right)^{\alpha} W\left(t_{1}\right)-\int_{t_{1}}^{t}(t-u)^{\alpha} \rho(u) q(u) d u \tag{18}
\end{equation*}
$$

and we see that (6) implies that

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \inf t^{-\alpha} \int_{s}^{t}(t-u)^{\alpha} \rho(u) q(u) d u \geq \phi(s), \quad s \geq t_{1} \tag{19}
\end{equation*}
$$

and
(20) $\lim _{t \rightarrow \infty} \sup t^{-\alpha} \int_{t_{1}}^{t}(t-u)^{\alpha} \rho(u) q(u) d u$

$$
-\lim _{t \rightarrow \infty} \inf \frac{M^{2}}{4 \lambda} t^{-\alpha} \int_{t_{1}}^{t}[\gamma(u) \eta(u)(t-u)-\alpha]^{2} \frac{1}{v(u)}(t-u)^{\alpha-2} d u \geq \phi\left(t_{1}\right)
$$

Together with (5), (20) shows that there exists a sequence

$$
\begin{equation*}
\left\{T_{n}\right\}_{n=1,2, \ldots}, \quad T_{n}>t_{1}, n=1,2, \ldots, \lim _{n \rightarrow \infty} T_{n}=\infty \tag{21}
\end{equation*}
$$

such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \frac{M^{2}}{4 \lambda} t^{-\alpha} \int_{t_{1}}^{T_{n}}[\gamma(u) \eta(u)(t-u)-\alpha]^{2} \frac{1}{v(u)}(t-u)^{\alpha-2} d u<\infty \tag{22}
\end{equation*}
$$

Next, taking the upper limit as $t \rightarrow \infty$ in (18) and using (19), we have
(23) $\lim _{t \rightarrow \infty} \sup [y(t)+z(t)] \leq W\left(t_{1}\right)-\lim _{t \rightarrow \infty} \inf t^{-\alpha} \int_{t_{1}}^{t}(t-u)^{\alpha} \rho(u) q(u) d u$

$$
\leq W\left(t_{1}\right)-\phi\left(t_{1}\right)=K
$$

Hence for all sufficiently large $n$

$$
\begin{equation*}
y\left(T_{n}\right)+z\left(T_{n}\right)<K \tag{24}
\end{equation*}
$$

Since

$$
z(t)=\frac{\lambda}{M^{2}} \int_{t_{1}}^{t}\left(1-\frac{u}{t}\right)^{\alpha} v(u) W^{2}(u) d u>0
$$

is increasing in $t>t_{1}$, we see that $\lim _{t \rightarrow \infty} z(t)=b$, where $b=\infty$ or is a positive constant. Suppose that $b=\infty$. Then $\lim _{n \rightarrow \infty} z\left(T_{n}\right)=\infty$ and by (24) we have

$$
\begin{equation*}
\lim _{n \rightarrow \infty} y\left(T_{n}\right)=-\infty \tag{25}
\end{equation*}
$$

Now (24) and (25) lead to

$$
\frac{y\left(T_{n}\right)}{z\left(T_{n}\right)}+1<\sigma, \quad \text { where } 0<\sigma<1 \text { is a constant }
$$

that is,

$$
\begin{equation*}
\frac{y\left(T_{n}\right)}{z\left(T_{n}\right)}<\sigma-1<0 \quad \text { for all large } T_{n} \tag{26}
\end{equation*}
$$

On the other hand, by the Schwarz inequality we have

$$
\begin{aligned}
0 \leq & T_{n}^{-2 \alpha}\left(\int_{t_{1}}^{T_{n}}\left[\gamma(u) \eta(u)\left(T_{n}-u\right)-\alpha\right]\left(T_{n}-u\right)^{\alpha-1} W(u) d u\right)^{2} \\
\leq & \left(T_{n}^{-\alpha} \int_{t_{1}}^{T_{n}}\left[\gamma(u) \eta(u)\left(T_{n}-u\right)-\alpha\right]^{2}\left(T_{n}-u\right)^{\alpha-2} \frac{1}{v(u)} d u\right)^{2} \\
& \cdot\left(T_{n}^{-\alpha} \int_{t_{1}}^{T_{n}}\left(t_{n}-u\right)^{\alpha} v(u) W^{2}(u) d u\right), \quad \text { for all large } T_{n}
\end{aligned}
$$

Thus,

$$
0 \leq \frac{y^{2}\left(T_{n}\right)}{z\left(T_{n}\right)} \leq \frac{M^{2}}{\lambda} T_{n}^{-\alpha} \int_{t_{1}}^{T_{n}}\left[\gamma(u) \eta(u)\left(T_{n}-u\right)-\alpha\right]^{2} \frac{1}{v(u)}\left(T_{n}-u\right)^{\alpha-2} d u
$$

By (22), we get

$$
\begin{aligned}
0 & \leq \lim _{n \rightarrow \infty} \frac{\left.y^{( } T_{n}\right)}{z\left(T_{n}\right)} \\
& \leq \frac{M^{2}}{\lambda} \lim _{n \rightarrow \infty} \int_{t_{1}}^{T_{n}}\left[\gamma(u) \eta(u)\left(T_{n}-u\right)-\alpha\right]^{2} \frac{1}{v(u)}\left(T_{n}-u\right)^{\alpha-2} d u<\infty
\end{aligned}
$$

which contradicts (25) and (26). Hence $\lim _{t \rightarrow \infty} z(t)=b<\infty$. Using (17), we obtain

$$
\begin{aligned}
\lim _{t \rightarrow \infty} t^{-\alpha} \int_{t_{1}}^{t}(t-u)^{\alpha} v(u) \phi_{+}^{2}(u) d u & \leq \lim _{t \rightarrow \infty} t^{-\alpha} \int_{t_{1}}^{t}(t-u)^{\alpha} v(u) W^{2}(u) d u \\
& =\frac{M^{2} b}{\lambda}<\infty
\end{aligned}
$$

which contradicts condition (7).
Case 2. $x^{\bullet}>0$ on $[T, \infty)$ for some $T \geq t_{0}$. From (2) and (10) it follows that (12) holds, and hence we can complete the proof by the procedure of Case 1.

Case 3. $x^{\bullet}<0$ on $[T, \infty)$ for some $T \geq t_{0}$. If (12) holds, then we can arrive at a contradiction by the procedure of Case 1 . So, we suppose that the
integral in (12) diverges. Using (2) in (10) we have

$$
\begin{equation*}
-a(t) \rho(t) \frac{x^{\bullet}(t)}{x^{\lambda}(t)} \geq-C+\lambda \int_{T}^{t} a(s) \rho(s)\left(\frac{\left(x^{\bullet}(s)\right)^{2}}{x^{\lambda+1}(s)}\right) d s \tag{27}
\end{equation*}
$$

where $C$ is a constant. By the assumptions, we can choose $T_{1} \geq T$ so that

$$
\lambda \int_{T}^{T_{1}} a(s) \rho(s) \frac{\left(x^{\bullet}(s)\right)^{2}}{x^{\lambda+1}(s)} d s=1+C
$$

and then for any $t \geq T_{1}$ we get

$$
\frac{-a(t) \rho(t) \frac{x^{\bullet}(t)}{x^{\lambda}(t)}\left(-\lambda \frac{x^{\bullet}(t)}{x(t)}\right)}{-C+\lambda \int_{T}^{t} a(s) \rho(s) \frac{\left(x^{\star}(s)\right)^{2}}{x^{\alpha+1}(s)} d s} \geq-\lambda \frac{x^{\bullet}(t)}{x(t)} .
$$

Integrating the above inequality from $T_{1}$ to $t$ we obtain

$$
\ln \left[-C+\lambda \int_{T}^{t} a(s) \rho(s) \frac{\left(x^{\bullet}(s)\right)^{2}}{x^{\lambda+1}(s)} d s\right] \geq \lambda \int_{T_{1}}^{t}\left(\frac{-x^{\bullet}(s)}{x(s)}\right) d s=\ln \left(\frac{x\left(T_{1}\right)}{x(t)}\right)^{\lambda}
$$

which together with (27) yields

$$
-a(t) \rho(t) \frac{x^{\bullet}(t)}{x^{\lambda}(t)} \geq\left(\frac{x\left(T_{1}\right)^{\lambda}}{x(t)}\right)
$$

from which it follows that

$$
x^{\bullet}(t) \leq-\left(-x\left(T_{1}\right)\right)^{\lambda} \frac{1}{a(t) \rho(t)}<0 \quad \text { for } t \geq T_{1}
$$

or

$$
x(t) \leq x\left(T_{1}\right)-\left(x\left(T_{1}\right)\right)^{\lambda} \int_{T_{1}}^{t} \frac{1}{a(s) \rho(s)} d s \rightarrow-\infty \quad \text { as } t \rightarrow \infty
$$

contradicting the fact that $x(t)>0$ for $t \geq t_{0}$. This completes the proof.
The follow result is concerned with the oscillatory behavior of equation (1) for all $\lambda>0$.

Corollary 1. Let the differentiable function $\rho$ assumed in Theorem 1 be defined by

$$
\begin{equation*}
\rho(t)=\exp \left(\int_{t_{0}}^{t} \frac{p(s)}{a(s)} d s\right) \quad \text { for } t \geq t_{0} \tag{28}
\end{equation*}
$$

and let conditions (2), (4)-(7) hold. Then equation (1) is oscillatory for all $\lambda>0$.

Proof. Let $x(t)$ be a nonoscillatory solution of equation (1), say $x(t)>0$ for $t \geq t_{0}$. From (28), we see that $\gamma(t)=0$ for all $t \geq t_{0}$. Now, if $W$ is
defined as in the proof of Theorem 1, then we obtain (10) or (11). The rest of the proof is similar to that of Theorem 1 and hence is omitted.

In the following corollary we study the oscillatory behavior of the undamped equation

$$
\begin{equation*}
\left(a(t) x^{\bullet}(t)\right)^{\bullet}+q(t)|x(t)|^{\lambda} \operatorname{sgn} x(t)=0, \quad \lambda>0, \tag{29}
\end{equation*}
$$

where the functions $a$ and $q$ are defined as in equation (1).

Corollary 2. Suppose that there exists a continuous function $\phi:\left[t_{0}, \infty\right)$
$\rightarrow R$ and $\alpha \in(1, \infty)$ such that

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \inf \int_{t_{0}}^{t} q(s) d s>-\infty \tag{30}
\end{equation*}
$$

$$
\begin{align*}
& \int^{\infty} \frac{1}{a(s)} d s=\infty  \tag{31}\\
& \lim _{t \rightarrow \infty} \sup \frac{1}{t^{\alpha}} \int_{t_{0}}^{t}(t-u)^{\alpha} q(u) d u<\infty \tag{32}
\end{align*}
$$

and for every $\theta>0$

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \inf \frac{1}{t^{\alpha}} \int_{s}^{t}(t-u)^{\alpha-2}\left[(t-u)^{2} q(u)-\frac{\theta \alpha^{2}}{4 \lambda v(u)}\right] d u>\phi(s) \tag{33}
\end{equation*}
$$

and

$$
\begin{equation*}
\int^{\infty} v(s) \phi_{+}^{2}(s) d s=\infty \tag{34}
\end{equation*}
$$

where $\phi_{+}(t)=\max \{\phi(t), 0\}$ and $v(t)=\frac{1}{a(t)}\left(\int_{t_{0}}^{t} \frac{1}{a(s)} d s\right)^{-1}$.
Then equation (29) is oscillatory for all $\lambda>0$.

Proof. This follows from Corollary 1 if we let $p(t)=0$ and $\rho(t)=1$ for $t \geq t_{0}$.

For illustration we consider the following example.
Example 1. Consider the differential equation

$$
\begin{align*}
&\left(\frac{1}{t} x^{\bullet}(t)\right)^{\bullet}+\frac{1}{t^{2}} x^{\bullet}(t)+\frac{\cos t}{t}|x(t)|^{\lambda} \operatorname{sgn} x(t)=0,  \tag{35}\\
& \lambda>0, t \geq t_{0}>0 .
\end{align*}
$$

Taking $\rho(t)=t$ and $\alpha=2$, we have

$$
\begin{aligned}
& \gamma(t)=0, \quad v(t)=\frac{1}{t} \quad \text { for } t \geq t_{0}>0 \\
& \lim _{t \rightarrow \infty} \inf \int_{t_{0}}^{t} \cos s d s>-\infty \\
& \lim _{t \rightarrow \infty} \sup \frac{1}{t^{2}} \int_{t_{0}}^{t}(t-u)^{2} \cos u d u=-\sin t_{0}<\infty \\
& \lim _{t \rightarrow \infty} \inf \frac{1}{t^{2}} \int_{s}^{t}\left[(t-u)^{2} \cos u-\frac{\theta}{\lambda u}\right] d u \geq-\sin s-k
\end{aligned}
$$

where $k$ is a positive constant. Set $\phi(s)=-\sin s-k$.
Next, we consider an integer $N$ such that $(2 N+1) \pi+\pi / 4 \geq t_{0}$. Then for all integer $n \geq N$ and $(2 n+1) \pi+\pi / 4 \leq s \leq(2 n+1) \pi-\pi / 4$,

$$
\phi(s)=-\sin s-k \geq \delta s
$$

where $\delta$ is a small constant. Thus,

$$
\lim _{t \rightarrow \infty} \int_{t_{0}}^{t} v(s) \phi_{+}^{2}(s) d s \geq \sum_{N=n}^{\infty} \delta^{2} \int_{(2 n+1) \pi+\pi / 4}^{2(n+1) \pi-\pi / 4} s d s=\infty
$$

Hence equation (35) is oscillatory for all $\lambda>0$ by Corollary 1 , whereas none of the known criteria can cover this result.

Remark 1. It is easy to check that the conditions (2), (3) and (4) of Corollary 1 are superfluous if $0<\lambda \leq 1$.

Remark 2. Corollary 1 , when $0<\lambda<1$, is new and is independent of the results of Kwong and Wong [6] and Philos [8], while Corollary 1 for the linear case $(\lambda=1)$ is of the type obtained by Tan [11].

The following result is concerned with the oscillatory solution of equation (1) when condition (5) fails.

Theorem 2. Suppose that there exist a differentiable function $\rho:\left[t_{0}, \infty\right) \rightarrow$ $(0, \infty)$ and a constant $\alpha \in(1, \infty)$ such that conditions (2)-(4) hold and for every $\theta>0$

$$
\begin{align*}
\lim _{t \rightarrow \infty} \sup \frac{1}{t^{\alpha}} \int_{t_{0}}^{t}(t-u)^{\alpha-2}[ & (t-u)^{2} \rho(u) q(u)  \tag{36}\\
& \left.-\frac{\theta}{4 \lambda} \frac{1}{v(u)}[\gamma(u) \eta(u)(t-u)-\alpha]^{2}\right] d u=\infty
\end{align*}
$$

where $\eta$ and $v$ are defined as in Theorem 1.
Then equation (1) is oscillatory for all $\lambda>1$.
Proof. Let $x(t)$ be a nonoscillatory solution of equation (1), say $x(t)>0$ for $t \geq t_{0}$. As in the proof of Theorem 1 (Case 1), we obtain (16). Divide
(16) by $t^{\alpha}$ and take the upper limit as $t \rightarrow \infty$. Using (36) we obtain a contradiction. The proofs in the cases when $x^{\bullet}$ is either positive or negative on $[T, \infty), T \geq t_{0}$, are similar to the proofs in Cases 2 and 3 of Theorem 1 , and hence will be omitted.

Corollary 3. Let condition (36) in Theorem 1 be replaced by

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \frac{1}{t^{\alpha}} \int_{t_{0}}^{t}(t-u)^{\alpha} \rho(u) q(u) d u=\infty \tag{37}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \int_{t_{0}}^{t} \frac{(t-u)^{\alpha-2}}{v(u)}[\gamma(u) \eta(u)(t-u)-\alpha]^{2} d u<\infty . \tag{38}
\end{equation*}
$$

Then the conclusions of Theorem 2 holds.

The following corollaries extend Wong's criterion in [10] to more general equations of the form of (1) and (29). The proofs are immediate consequences of Theorem 2 and will be omitted.

Corollary 4. Let the function $\rho$ in Theorem 2 be defined by (28) and suppose that conditions (2), (4) and (36) hold. Then equation (1) is oscillatory for all $\lambda>0$.

Corollary 5. Let conditions (30) and (31) hold for every $\theta>0$ and $\alpha \in(1, \infty)$ and suppose

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \sup \frac{1}{t^{\alpha}} \int_{t_{0}}^{t}(t-u)^{\alpha-2}\left[(t-u)^{2} q(u)-\frac{\theta \alpha^{2}}{4 v(u)}\right] d u=\infty \tag{39}
\end{equation*}
$$

where

$$
v(t)=\frac{1}{a(t)}\left(\int_{t_{0}}^{t} \frac{1}{a(s)} d s\right)^{-1} .
$$

Then equation (29) is oscillatory for all $\lambda>0$.

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