Foundations of Chemical Kinetics

# Lecture 28: Diffusion-influenced reactions, Part I

Marc R. Roussel

Department of Chemistry and Biochemistry

University of Lethbridge

## The encounter pair

- One of the major differences between gas-phase and solution-phase kinetics is solvent caging.
  - In the gas phase, a collision is a single event with a very short lifetime.
  - In solution, once two molecules have come into direct contact with each other, they may stay in contact for a long time because the solvent molecules that surround them need to move in order to allow them to move away from each other.
- ► A pair of molecules that are in contact and surrounded by solvent is called an encounter pair.

## Diffusion-influenced reactions

- In solution, the rate of reaction will in general depend both on the intrinsic reaction rate and on the rate at which molecules diffuse into close proximity.
- We then say that the reaction is diffusion-influenced.
- In what follows, we will consider a reaction

$$A + B \xrightarrow{k} product(s)$$

with rate

$$v = k[A][B]$$

- For simplicity, assume spherical molecules.
- The molecules need to touch in order to react.
- They touch when their centres are a distance  $R_{AB} = R_A + R_B$  apart.
- We focus on one particular A molecule, and assume the solution is sufficiently dilute that the distribution of B molecules around one A molecule does not affect the distribution around the others.
- ▶ We assume that the A molecules are stationary. We can compensate for this by replacing D<sub>B</sub> by the relative diffusion coefficient D<sub>AB</sub> = D<sub>A</sub> + D<sub>B</sub>.
- ▶ If the concentration of encounter pairs is in a steady state, the rate at which B molecules reach a distance *R*<sub>AB</sub> from the centre of an A molecule balances the rate of reaction.

- ► The flux at r = R<sub>AB</sub> is the rate of arrival of molecules of B at a sphere of radius R<sub>AB</sub> centered on a given A molecule per unit area.
- Therefore, the steady-state condition becomes

$$v/[A] = k[B] = -4\pi R_{AB}^2 J_B(r = R_{AB})$$

In a steady state, because the concentration between any two shells of arbitrary radii r<sub>1</sub> and r<sub>2</sub> is constant, the flux through a shell of any r must be the same. Thus,

$$k[\mathsf{B}] = -4\pi r^2 J_B(r)$$

 The steady-state distribution of B around a given A should be (on average) spherical. Thus,

$$J_{B} = -\left[D_{AB}\frac{d[B]_{r}}{dr} + \frac{z_{B}e}{k_{B}T}D_{AB}[B]_{r}\frac{dV}{dr}\right]$$

where  $[B]_r$  is the concentration of B at distance r from a molecule of A.

• Using  $U(r) = z_B eV(r)$ , we get

$$J_{B} = -D_{AB} \left[ \frac{d[B]_{r}}{dr} + \frac{1}{k_{B}T} [B]_{r} \frac{dU}{dr} \right]$$

The steady state condition becomes

$$k[\mathsf{B}] = 4\pi r^2 D_{AB} \left[ \frac{d[\mathsf{B}]_r}{dr} + \frac{1}{k_B T} [\mathsf{B}]_r \frac{dU}{dr} \right]$$

- In this equation [B] is the average concentration of B in the solution, which is also the expected concentration of B far from any given A molecule.
- Note:

$$\frac{d}{dr}\left[\left[\mathsf{B}\right]_{r}\exp\left(\frac{U(r)}{k_{B}T}\right)\right] = \exp\left(\frac{U(r)}{k_{B}T}\right)\left[\frac{d[\mathsf{B}]_{r}}{dr} + \frac{1}{k_{B}T}[\mathsf{B}]_{r}\frac{dU}{dr}\right]$$

Therefore,

$$k[B] = 4\pi r^2 D_{AB} \exp\left(-\frac{U(r)}{k_B T}\right) \frac{d}{dr} \left[ [B]_r \exp\left(\frac{U(r)}{k_B T}\right) \right]$$

Rearranging, we get

$$\frac{d}{dr}\left[[\mathsf{B}]_r \exp\left(\frac{U(r)}{k_B T}\right)\right] = \frac{k[\mathsf{B}]}{4\pi r^2 D_{AB}} \exp\left(\frac{U(r)}{k_B T}\right)$$

We can solve this equation by separation of variables subject to the boundary conditions [B]<sub>r</sub> → [B] and U(r) → 0 as r → ∞:

$$\int_{r=R_{AB}}^{\infty} d\left[[B]_{r} \exp\left(\frac{U(r)}{k_{B}T}\right)\right] = \int_{R_{AB}}^{\infty} \frac{k[B]}{4\pi r^{2} D_{AB}} \exp\left(\frac{U(r)}{k_{B}T}\right) dr$$
$$\therefore \left[[B]_{r} \exp\left(\frac{U(r)}{k_{B}T}\right)\right]_{R_{AB}}^{\infty} = \frac{k[B]}{4\pi D_{AB}} \int_{R_{AB}}^{\infty} \frac{1}{r^{2}} \exp\left(\frac{U(r)}{k_{B}T}\right) dr$$

Applying the limits, we get

$$[\mathsf{B}] - [\mathsf{B}]_{R_{AB}} \exp\left(\frac{U(R_{AB})}{k_B T}\right) = \frac{k[\mathsf{B}]}{4\pi D_{AB}\beta} \tag{1}$$

where

$$\beta^{-1} = \int_{R_{AB}}^{\infty} \frac{1}{r^2} \exp\left(\frac{U(r)}{k_B T}\right) dr$$

There is some intrinsic rate constant for reaction when A and B are in contact, k<sub>R</sub>, such that

$$v = k_R[A][B]_{R_{AB}} = k[A][B]$$
$$\therefore [B]_{R_{AB}} = k[B]/k_R$$

The next step is to substitute for [B]<sub>RAB</sub> in equation (1) and solving for k.

The result is

$$k = \frac{4\pi D_{AB}\beta k_R}{k_R + 4\pi D_{AB}\beta \exp\left(\frac{U(R_{AB})}{k_B T}\right)}$$

In order to get a rate constant in molar units, we need to multiply this equation by L:

$$k = \frac{4\pi L D_{AB}\beta k_R}{k_R + 4\pi D_{AB}\beta \exp\left(\frac{U(R_{AB})}{k_B T}\right)}$$

• This equation allows us to calculate the rate constant for a diffusion-influenced reaction, provided we know the intrinsic rate constant  $k_R$ , the potential energy U(r), and the diffusion coefficients and radii of A and B.

### The diffusion-limited rate constant

- Suppose that k<sub>R</sub> is very large, i.e. that A and B react nearly every time they meet in solution.
  We then say that the reaction is diffusion-limited.
- For  $k_R$  very large, we get

$$k = \frac{4\pi L D_{AB}\beta k_R}{k_R + 4\pi D_{AB}\beta \exp\left(\frac{U(R_{AB})}{k_BT}\right)} \to 4\pi D_{AB}\beta$$

This quantity is the diffusion-limited rate constant:

$$k_D = 4\pi L D_{AB}\beta$$

 The diffusion-influenced rate constant can consequently be written

$$k = \frac{k_D k_R}{k_R + k_D \exp\left(\frac{U(R_{AB})}{k_B T}\right)}$$

#### The diffusion-limited rate constant Weak intermolecular forces

- If intermolecular forces between A and B are weak, then U(r) ≈ 0, except when A and B are very close.
- In this case,

$$\beta^{-1} = \int_{R_{AB}}^{\infty} \frac{1}{r^2} \exp\left(\frac{U(r)}{k_B T}\right) dr \approx \int_{R_{AB}}^{\infty} \frac{1}{r^2} dr = \frac{1}{R_{AB}}$$

or  $\beta = R_{AB}$ .

The diffusion-limited rate constant becomes

$$k_D = 4\pi L D_{AB} R_{AB}$$

and the diffusion-influenced rate constant is

$$k = \frac{k_D k_R}{k_R + k_D}$$

#### The diffusion-limited rate constant The Coulomb interaction

If we have a reaction between ions,

$$U(r) = \frac{z_A z_B e^2}{4\pi\epsilon r}$$

where  $\epsilon$  is the permittivity of the solvent.

Note: The textbook writes this formula in cgs units, which is why the factor of  $4\pi$  doesn't appear in their formula.

Moreover, in their formula,  $\epsilon$  is the dielectric constant, not the permittivity.

For this potential, we get

$$\beta = \frac{z_A z_B e^2}{4\pi\epsilon k_B T \left[ \exp\left(\frac{z_A z_B e^2}{4\pi\epsilon k_B T R_{AB}}\right) - 1 \right]}$$

The diffusion-limited rate constant The Coulomb interaction (continued)

$$\beta = \frac{z_A z_B e^2}{4\pi\epsilon k_B T \left[ \exp\left(\frac{z_A z_B e^2}{4\pi\epsilon k_B T R_{AB}}\right) - 1 \right]}$$

For z<sub>A</sub>z<sub>B</sub> > 0, putting in some typical numbers, you would find β ≪ R<sub>AB</sub>.

Thus, as you might expect,  $k_D$  is decreased due to Coulomb repulsion.

Conversely, for z<sub>A</sub>z<sub>B</sub> < 0, β > R<sub>AB</sub>, so k<sub>D</sub> is increased when the reactants are subject to an attractive potential.